

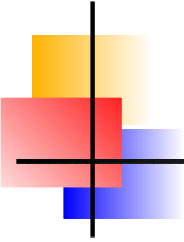
Particle Theory of $0\nu\beta\beta$ decay

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<http://www.astroparticles.es/>



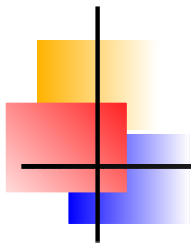
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II. $0\nu\beta\beta$ decay, LNV operators and m_ν

III. $0\nu\beta\beta$ decay, LNV and LHC

IV. Conclusions

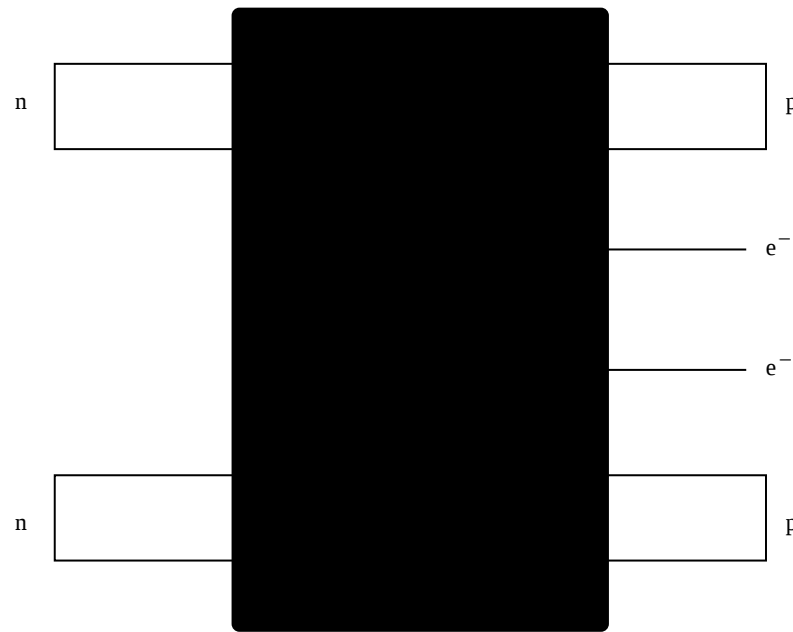


I.

Introduction

Black box!

Experimentalist observes:



Neutrinoless double
beta decay is:

$$nn \rightarrow pp + e^- e^-$$

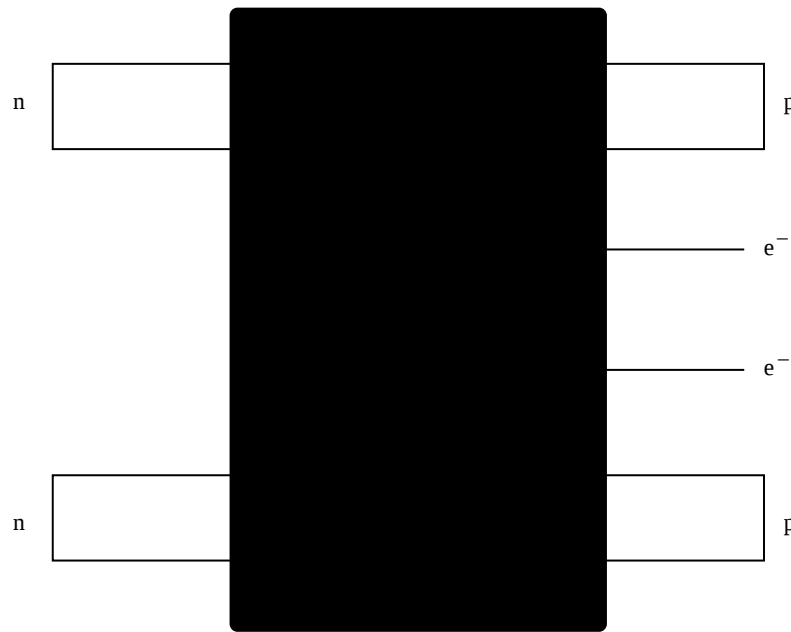
(no missing energy)

Lepton number is violated:

$$\Delta L = 2 !$$

Black box!

Experimentalist observes:



Express half-life as:

$$\left[T_{1/2}^{0\nu\beta\beta} \right]^{-1} = G_{0\nu} |\epsilon_i|^2 |M_i^{0\nu\beta\beta}|^2$$

$G_{0\nu}$ - phase space

$M_i^{0\nu\beta\beta}$ - nuclear matrix element

ϵ_i - particle physics

Many, many
possible
contributions:

RPV SUSY

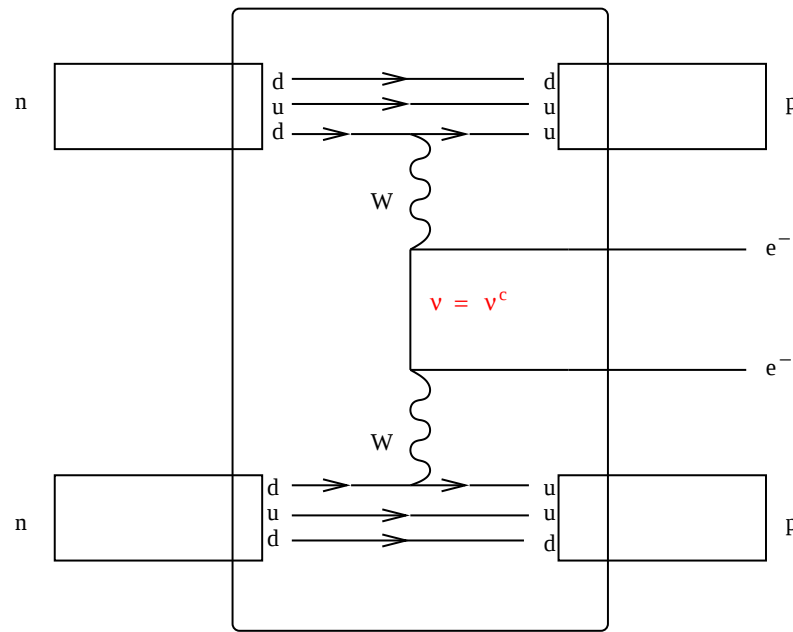
Leptoquarks

Left-right symmetry

... etc ... etc ...

Mass mechanism

Experimentalist observes:



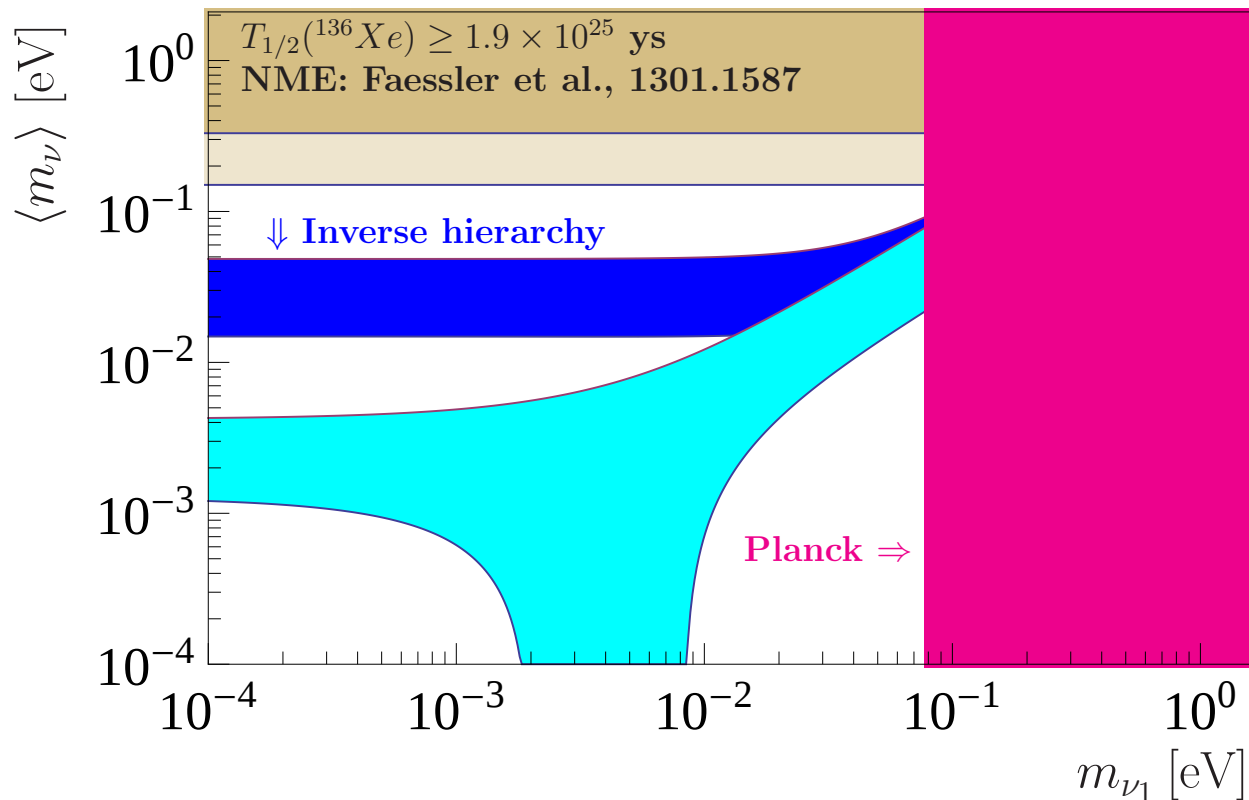
Neutrino
propagator:

$$\int \frac{d^4 p}{(2\pi)^4} \frac{m_\nu + \not{p}}{p^2 - m_\nu^2}$$

“Mass mechanism” because weak interaction is left-handed:

$$P_L(m_\nu + \not{p})P_L = m_\nu P_L$$

$\langle m_\nu \rangle$ versus m_{ν_1} - status 2015



Global fit data from:

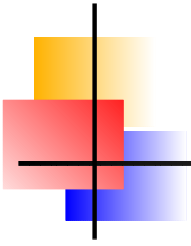
Forero, Tortola
& Valle;

arXiv:1405.7540

all ranges at 1σ c.l.

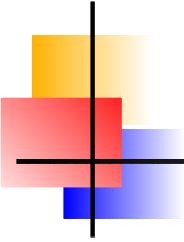
$\Rightarrow$ Planck - limits from cosmological data

$\Rightarrow T_{1/2}(^{136}\text{Xe})$ - limit from KamLAND-Zen



III.

$0\nu\beta\beta$ decay,
LNV operators and m_ν



Theoretical expectation?

Majorana Neutrino mass

$$m_\nu \simeq \frac{(Y\nu)^2}{\Lambda}.$$

Weinberg, 1979

Smallness of neutrino mass
can be “explained” by:

Minkowski, 1977

⇒ High scale: Large Λ
“classical” seesaw

Yanagida, 1979

Gell-Mann, Ramond, Slansky, 1979

Mohapatra, Senjanovic, 1980

Schechter, Valle, 1980

..., ..., ...

Foot et al., 1988

Theoretical expectation?

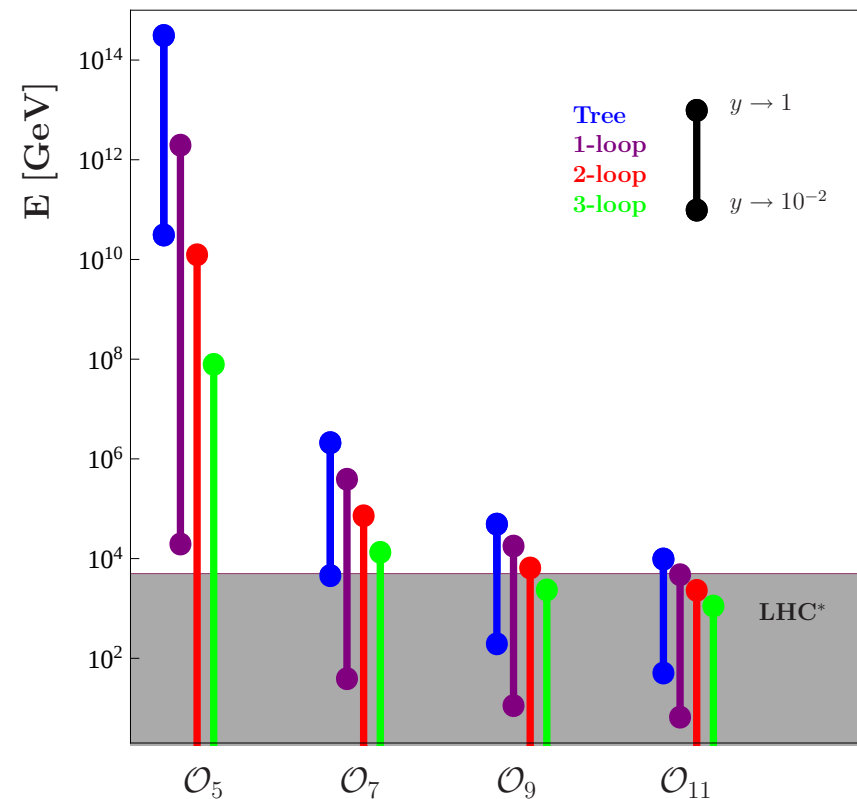
Majorana Neutrino mass generated from an n -loop dimension d diagram:

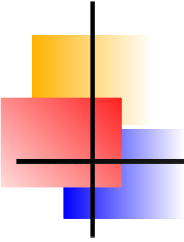
$$m_\nu \simeq \frac{(Yv)^2}{\Lambda} \cdot \epsilon \cdot \left(\frac{Y^2}{16\pi^2} \right)^n \cdot \left(\frac{Yv}{\Lambda} \right)^{d-5}$$

Smallness of neutrino mass
can be “explained” by:

- ⇒ High scale: Large Λ
“classical” seesaw
- ⇒ Loop factor: $n \geq 1$
+ “smallish” $Y \sim \mathcal{O}(10^{-3} - 10^{-1})$
- ⇒ Higher order: $d = 7, 9, 11$
- ⇒ Nearly conserved L ,
i.e. small ϵ (“inverse seesaw”)

... or combination thereof





$\Delta L = 2$ operators

$d = 5$:

Weinberg, 1979

$$\mathcal{O}_W \propto \frac{c_{ij}}{\Lambda}(L_i H)(L_j H)$$

One $d=5$

$\Delta L = 2$ operators

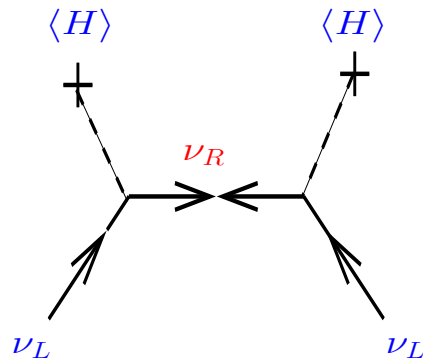
$d = 5$:

Weinberg, 1979

$$\mathcal{O}_W \propto \frac{c_{ij}}{\Lambda} (L_i H)(L_j H)$$

One $d=5$

Example realization, seesaw type-I:



$$\Lambda \simeq M_{\nu R_k}$$

$$c_{ij} \propto Y_{ik}^\nu Y_{jk}^\nu$$

$\Delta L = 2$ operators

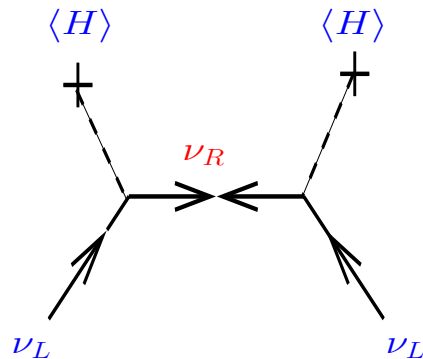
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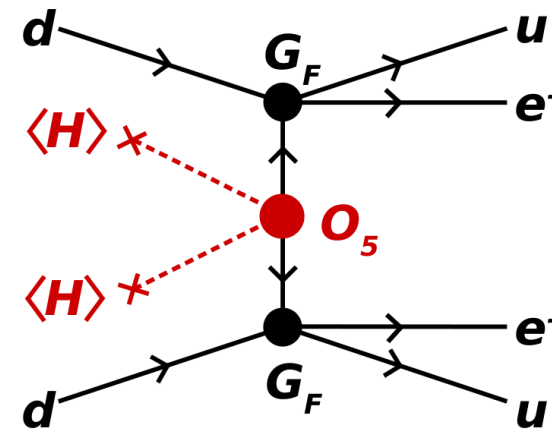
Example realization, seesaw type-I:



$$\Lambda \simeq M_{\nu_R k}$$

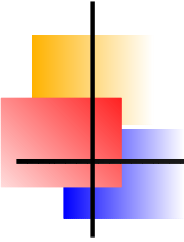
$$c_{ij} \propto Y_{ik}^\nu Y_{jk}^\nu$$

$0\nu\beta\beta$ decay:



(a)

Mass mechanism!



$\Delta L = 2$ operators

$d = 5$:

Weinberg, 1979

$$\mathcal{O}_W \propto \frac{c_{ij}}{\Lambda}(L_i H)(L_j H)$$

One $d=5$

$d = 7$:

Babu & Leung, 2001

de Gouvea & Jenkins, 2007

$$\mathcal{O}_2 \propto L L L e^c H$$

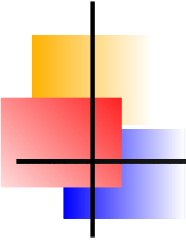
$$\mathcal{O}_3 \propto L L Q d^c H$$

$$\mathcal{O}_4 \propto L L \bar{Q} \bar{u}^c H$$

$$\mathcal{O}_8 \propto L \bar{e}^c \bar{u}^c d^c H$$

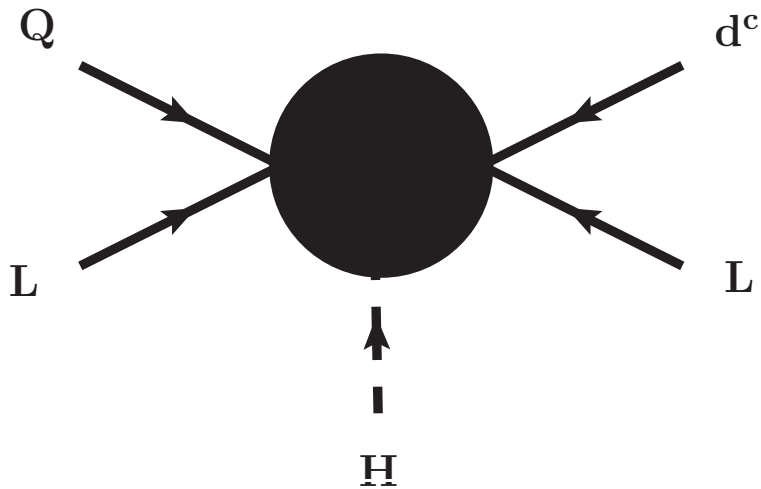
$$\mathcal{O} \propto (LH)(LH)(H_u H_d)$$

4 (+1) $d = 7$



Example $d = 7$: $LLQd^cH$

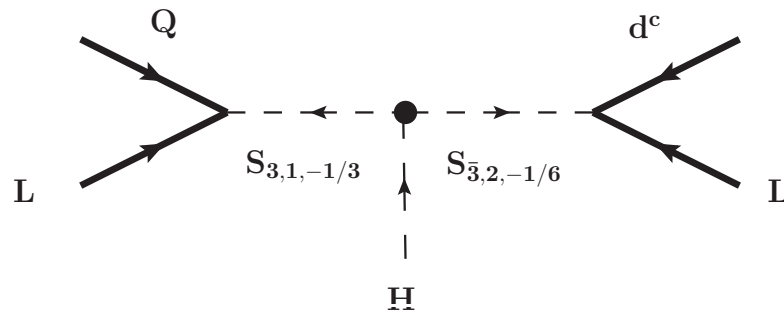
Graphically:



Example $d = 7$: $LLQd^cH$

Again, more than one realization.

Example:



$S_{3,1,-1/3}$ - singlet leptoquark

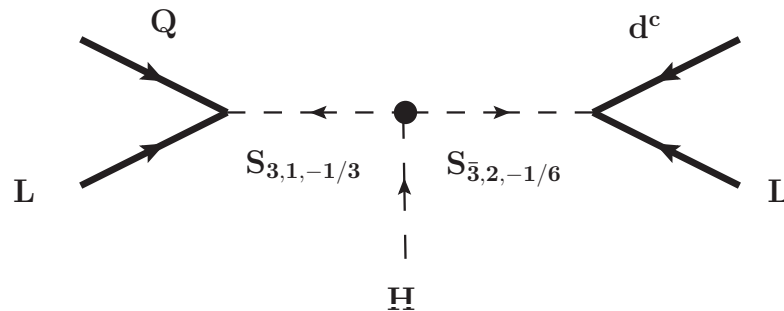
$S_{\bar{3},2,-1/6}$ - doublet leptoquark

$\Delta L = 2$, so ...

Example $d = 7$: $LLQd^cH$

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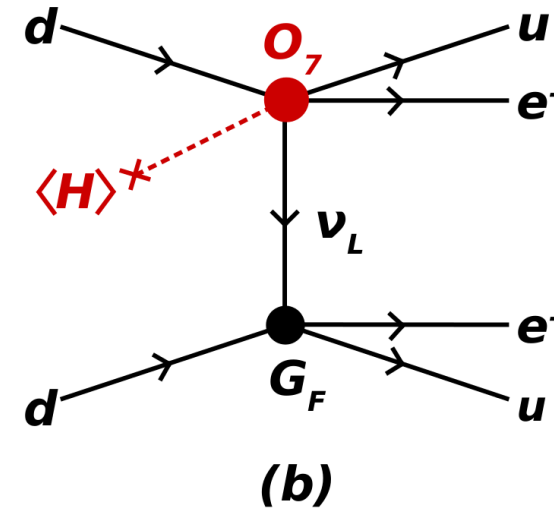


$S_{3,1,-1/3}$ - singlet leptoquark

$S_{3,2,1/6}$ - doublet leptoquark

$\Delta L = 2$, so ...

$0\nu\beta\beta$ decay:



Long range contribution!

$$A \propto \frac{\mu \times \langle H^0 \rangle}{m_{3,1,1/3}^2 m_{3,2,1/6}^2}$$

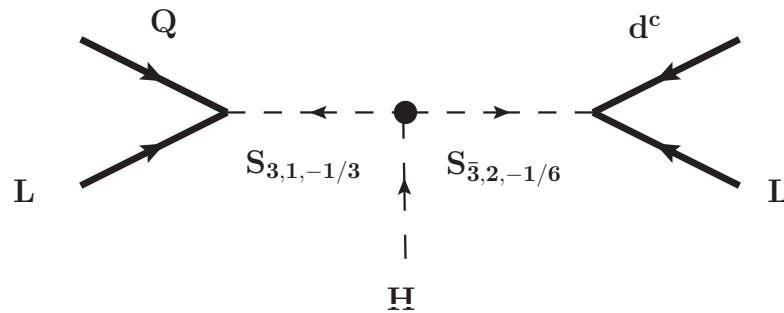
$$\propto \frac{v}{\Lambda^3}$$

No helicity suppression!

Example $d = 7$: $LLQd^cH$

Again, more than one realization.

Example:

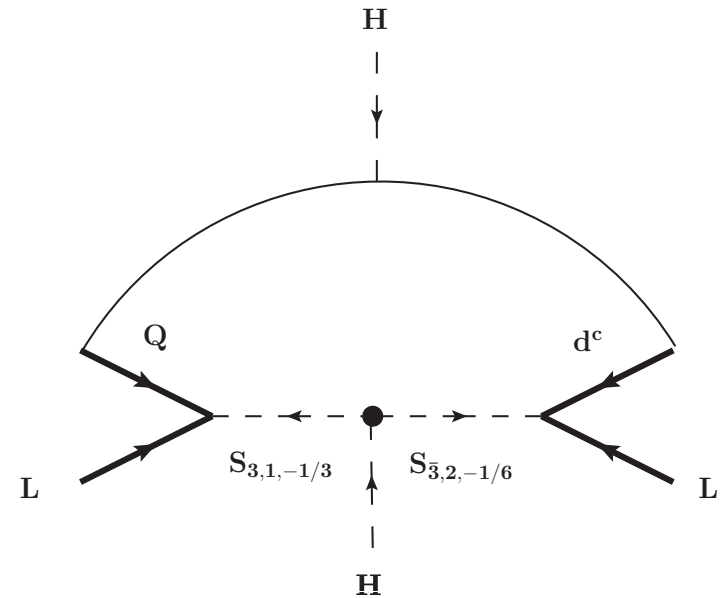


$S_{3,1,-1/3}$ - singlet leptoquark

$S_{\bar{3},2,-1/6}$ - doublet leptoquark

$\Delta L = 2$, so ...

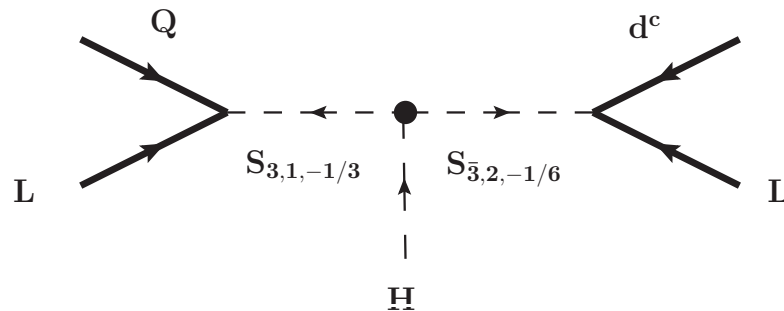
1-loop neutrino mass:



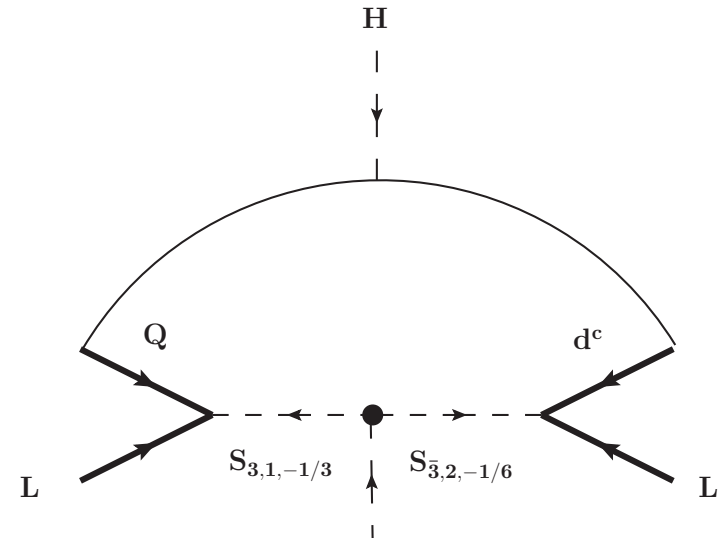
Example $d = 7$: $LLQd^cH$

Again, more than one realization.

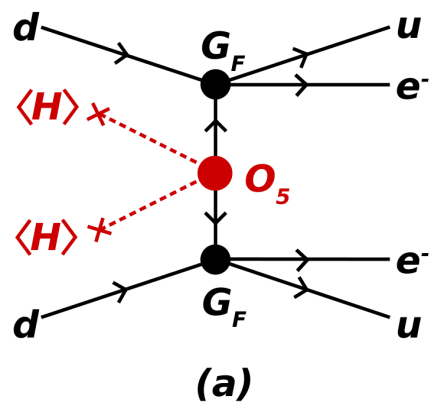
Example:



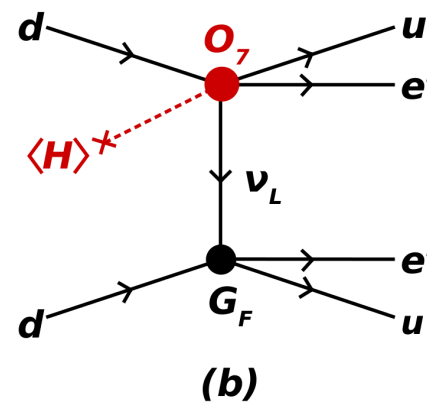
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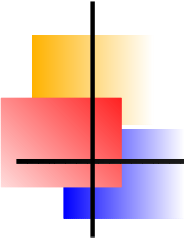


$0\nu\beta\beta$ decay has both contributions:



+





$\Delta L = 2$ operators

$d = 5$:

Weinberg, 1979

$$\mathcal{O}_W \propto \frac{c_{ij}}{\Lambda} (L_i H)(L_j H)$$

One $d=5$

$d = 7$:

Babu & Leung, 2001

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$$\mathcal{O}_3 \propto LLQd^c H$$

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$$\mathcal{O}_8 \propto L\bar{e}^c \bar{u}^c d^c H$$

$$\mathcal{O} \propto (LH)(LH)(H_u H_d)$$

4 (+1) $d = 7$

$d = 9$:

many $d = 9$ and $d = 11$ ops

$$\mathcal{O}_5 \propto LLQd^c HHH^\dagger$$

$$\mathcal{O}_6 \propto LL\bar{Q}\bar{u}^c HHH^\dagger H$$

$$\mathcal{O}_7 \propto LQ\bar{e}^c \bar{Q}HHH^\dagger$$

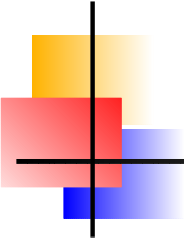
.....

$$\mathcal{O}_9 \propto LLLe^c Le^c$$

$$\mathcal{O}_{10} \propto LLLe^c Qd^c$$

$$\mathcal{O}_{11} \propto LLQd^c Qd^c$$

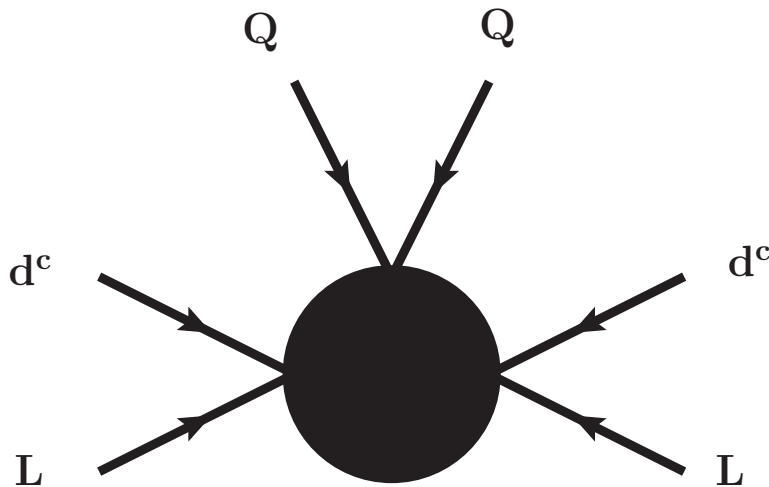
.....



Example $d = 9$: $LLQd^cQd^c$

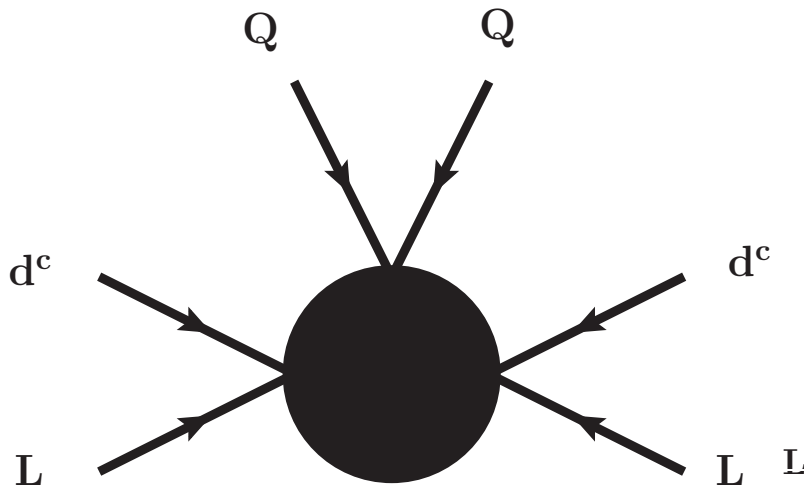
True $d = 9$ operator:

Many, many realizations ...

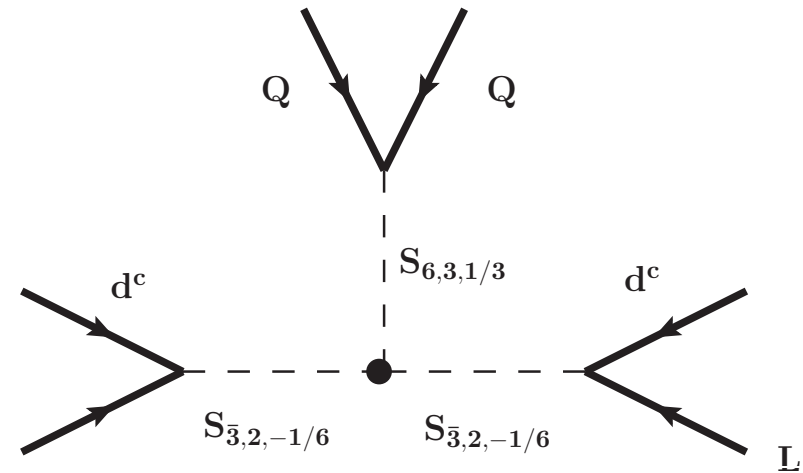


Example $d = 9$: $LLQd^cQd^c$

True $d = 9$ operator:



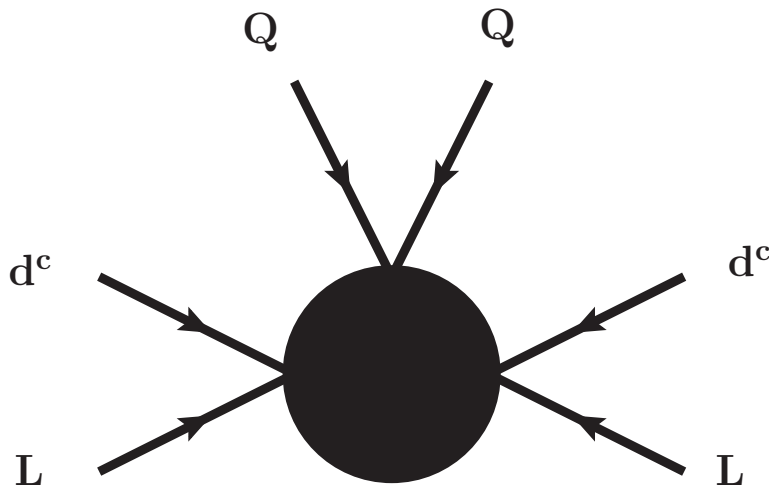
Many, many realizations ...
One example:



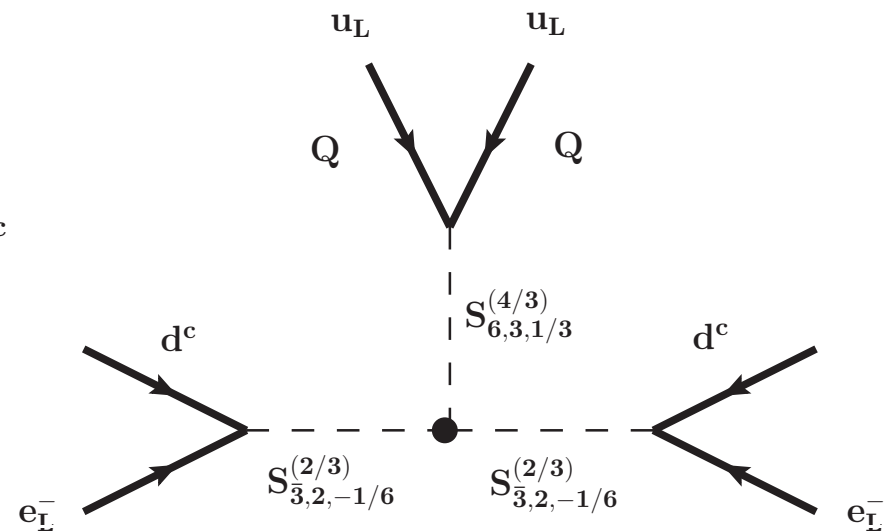
$S_{6,3,1/3}$ - triplet diquark
 $S_{\bar{3},2,-1/6}$ - doublet leptoquark

Example $d = 9$: $LLQd^cQd^c$

True $d = 9$ operator:



Many, many realizations ...
One example:

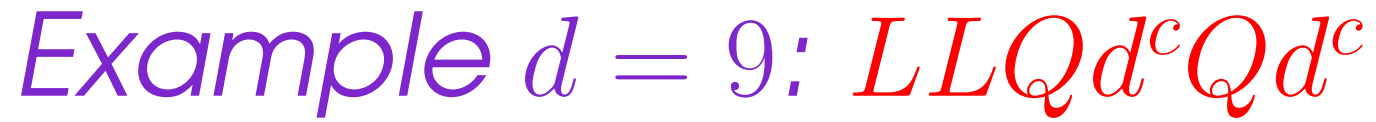


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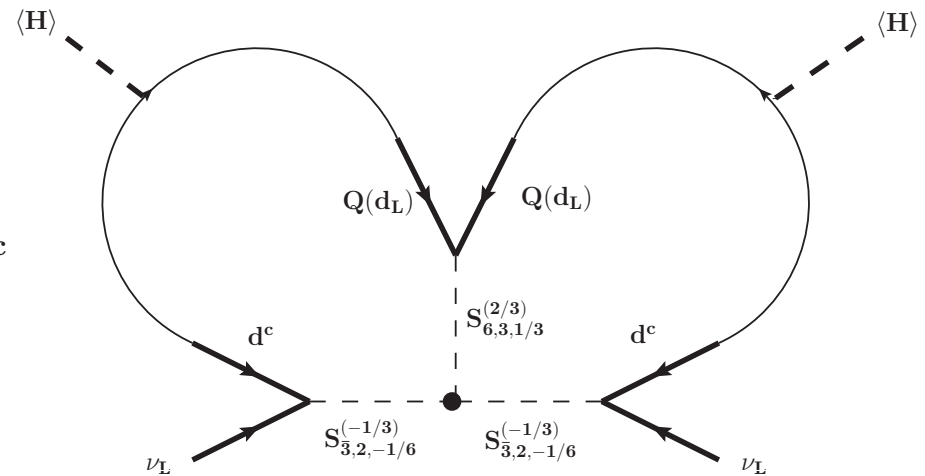
$S_{3,2,1/6}$ - doublet leptoquark

$0\nu\beta\beta$ decay without neutrino!

$\Delta L = 2$, so ...



Many, many realizations ...
One example:

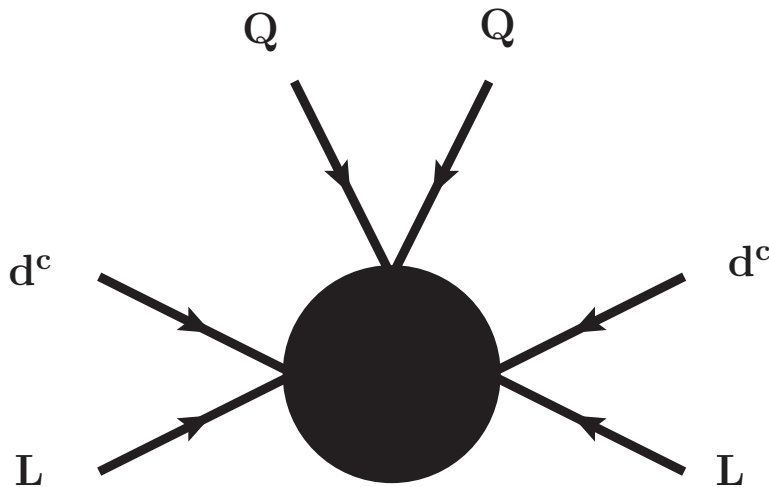


$S_{6,3,1/3}$ - triplet diquark
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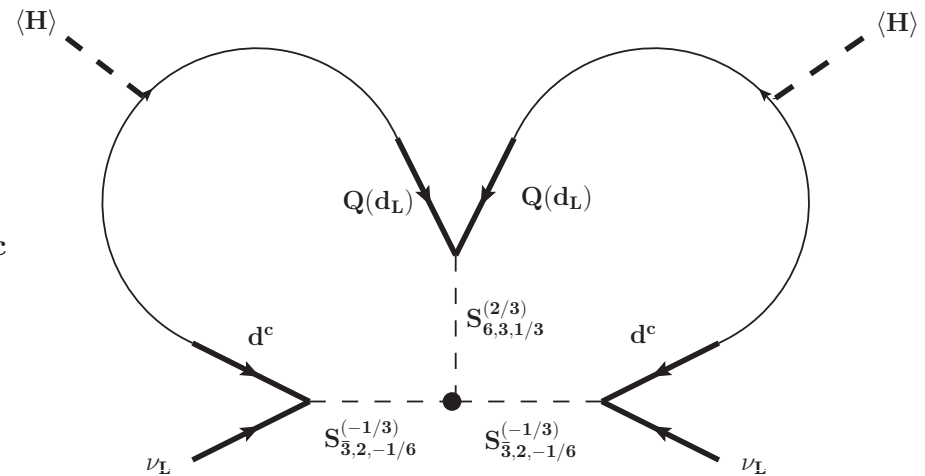
2-loop neutrino mass!

Example $d = 9$: $LLQd^cQd^c$

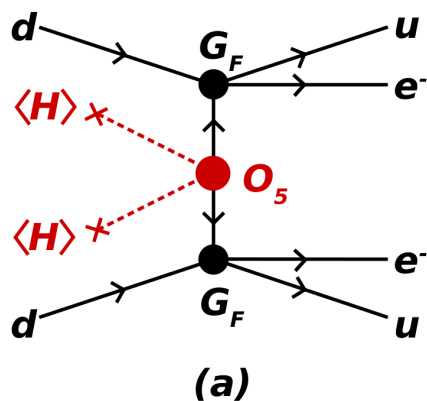
True $d = 9$ operator:



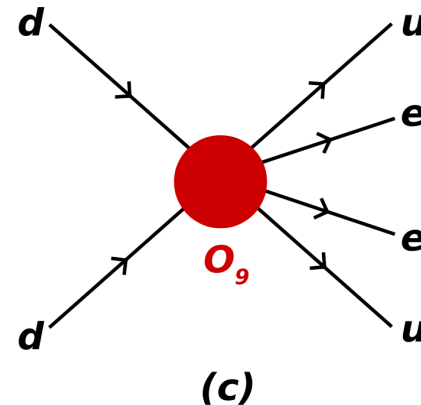
Many, many realizations ...
One example:



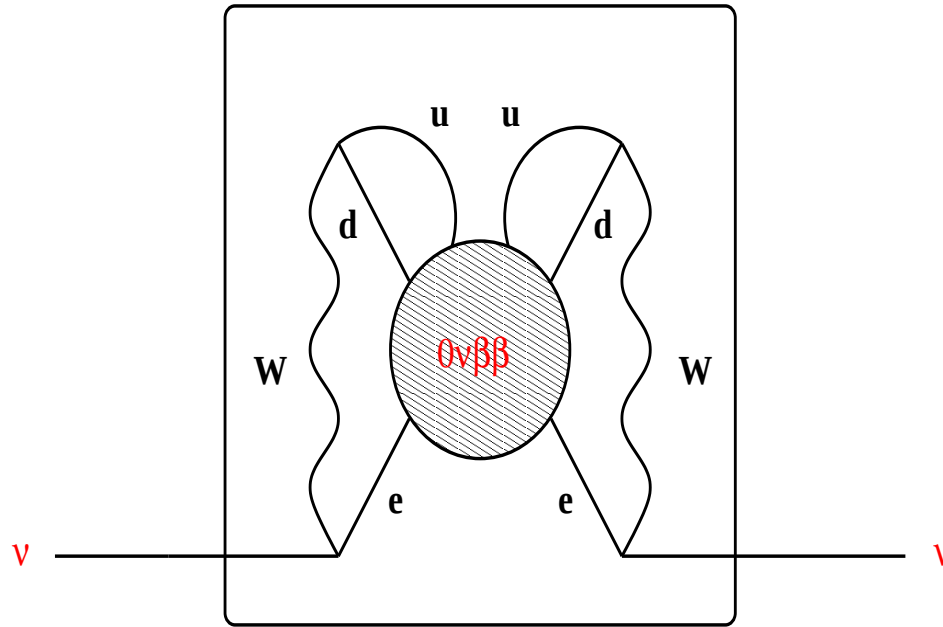
Again, $0\nu\beta\beta$ decay has two contributions:



+



Black Box Theorem



Schechter & Valle, PRD 1982

Takasugi, PLB 1984

If $0\nu\beta\beta$
is observed
the neutrino is a
Majorana particle!

⇒ 4-loop “butterfly” diagram: $m_\nu \sim 10^{-24}$ eV

Duerr et al 2011

⇒ Tree-level, 1-loop, ... 4-loop possible

⇒ Rule of thumb:

Helo et al., 2015

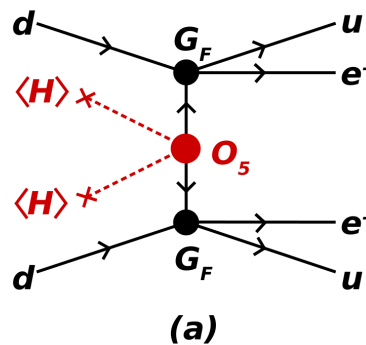
→ Models with tree-level, 1-loop m_ν - mass mechanism dominates

→ Models with 2-loop, 3-loop m_ν - mass mechanism $\sim$ SR

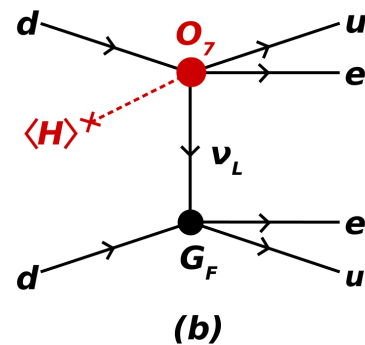
→ Models with 4-loop m_ν - SR dominates

Distinguish mechanisms?

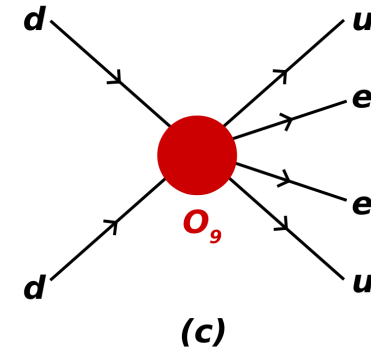
Amplitude for $(Z, A) \rightarrow (Z \pm 2, A) + e^\mp e^\mp$ can be divided into:



Mass mechanism



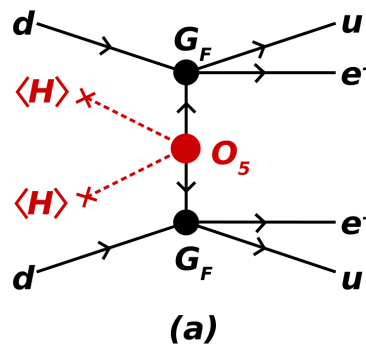
"long-range"



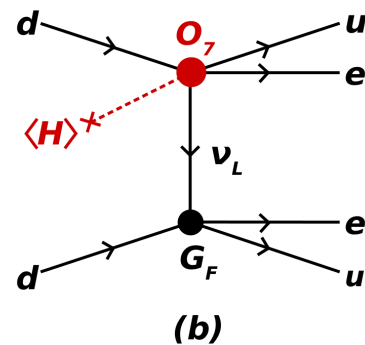
"short-range"

Distinguish mechanisms?

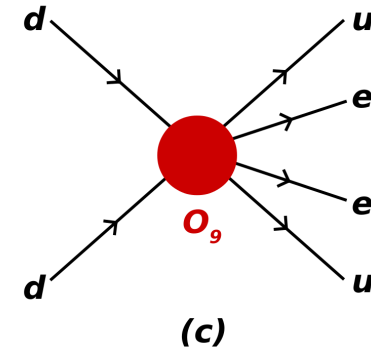
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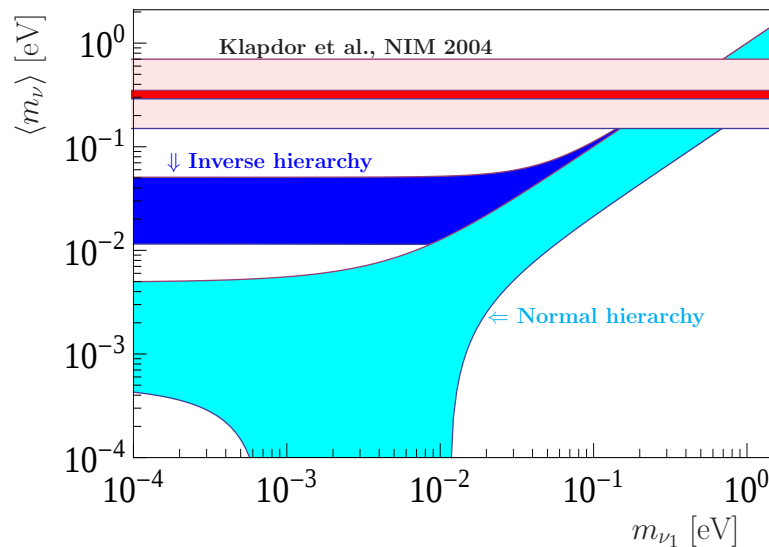


"long-range"



"short-range"

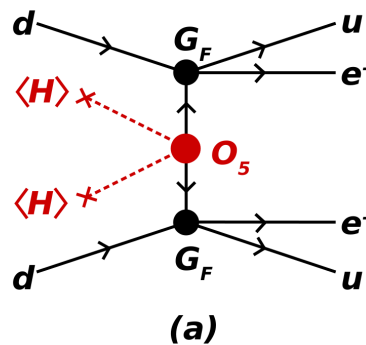
Compare with
other experiments:
Cosmology
KATRIN?



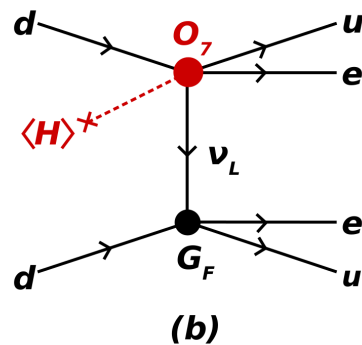
Red line:
Claim for
 $T_{1/2}^{0\nu\beta\beta}(^{76}\text{Ge}) =$
 $(1.19^{+0.37}_{-0.23}) \times 10^{25} \text{ ys}$

Distinguish mechanisms?

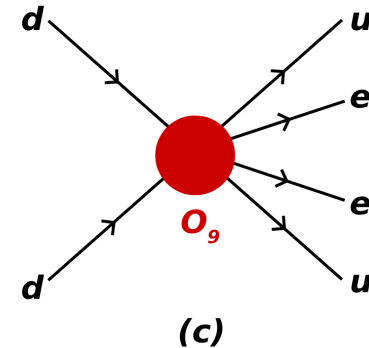
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Mass mechanism

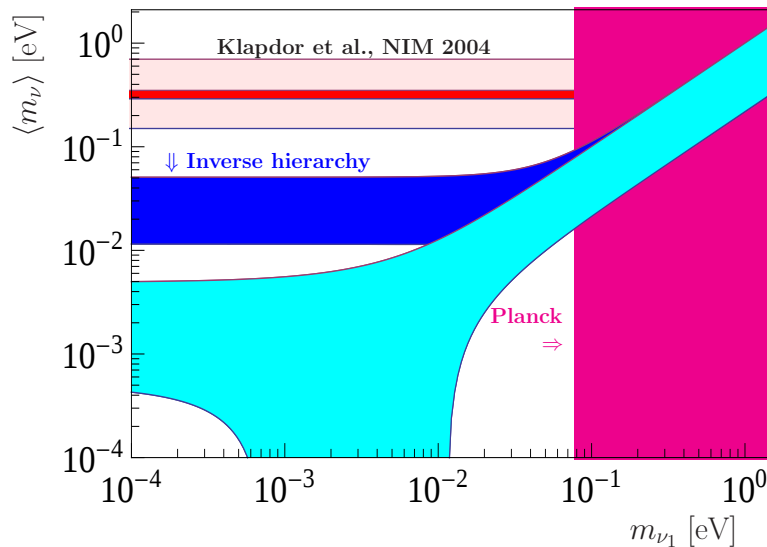


"long-range"



"short-range"

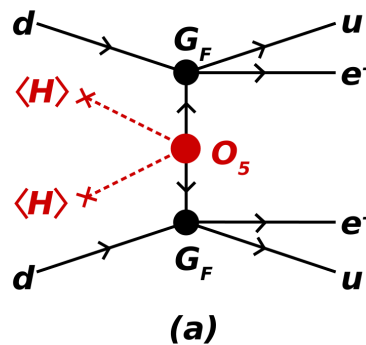
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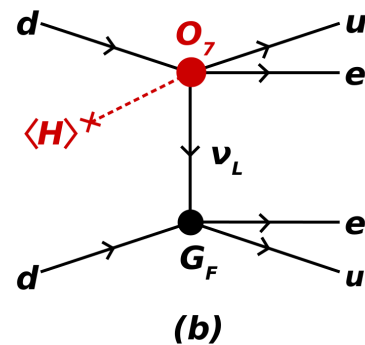
Planck + BAO:
NO overlap
with allowed
region of $\langle m_\nu \rangle$!

Distinguish mechanisms?

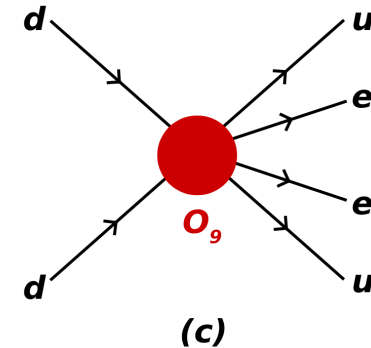
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Mass mechanism

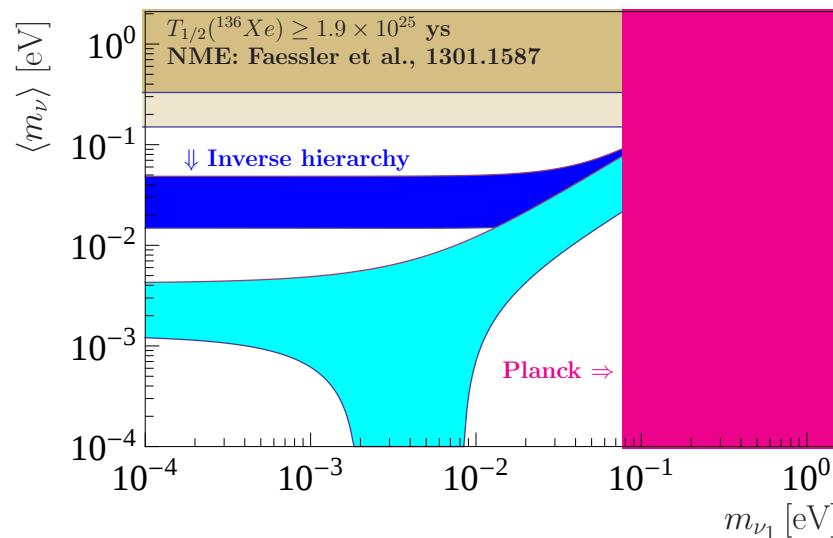


"long-range"



"short-range"

Compare with
other experiments:
Cosmology
KATRIN?

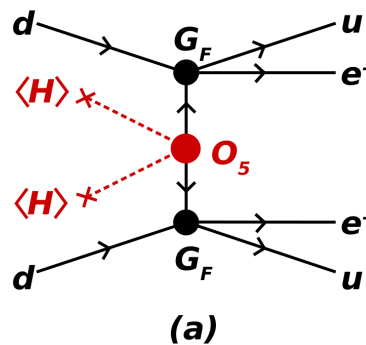


Claim now
ruled out by:

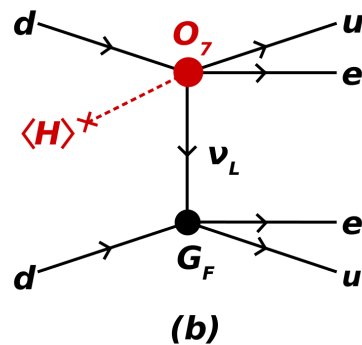
GERDA
EXO-200
KamLAND-Zen

Distinguish mechanisms?

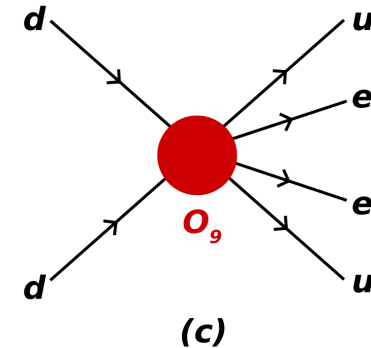
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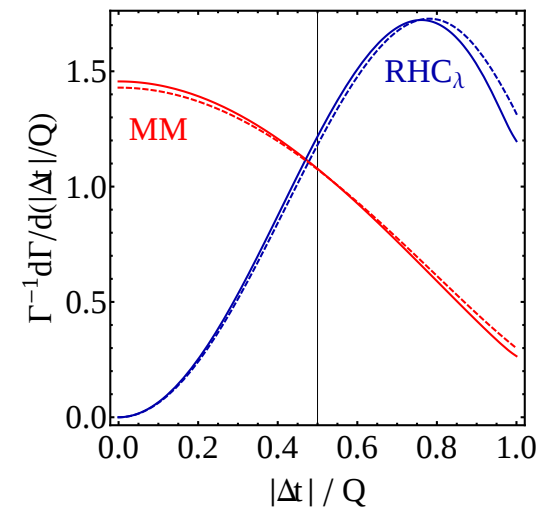
"long-range"



"short-range"

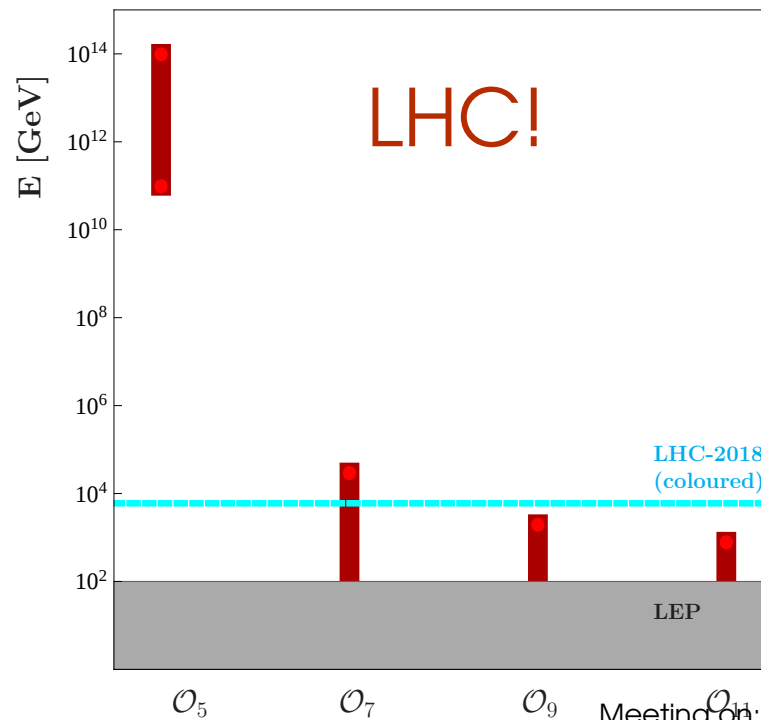
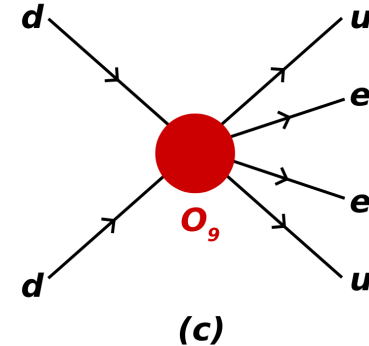
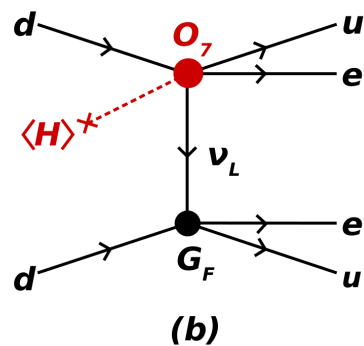
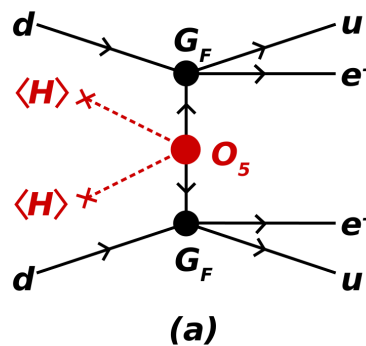
Angular correlations
 $0\nu\beta^+ / EC$ decays
 LHC?

SuperNEMO
 Arnold et al., 2010



Distinguish mechanisms?

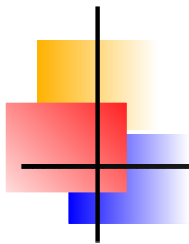
Amplitude for $(Z, A) \rightarrow (Z \pm 2, A) + e^\mp e^\mp$ can be divided into:



Estimate $0\nu\beta\beta$
sensitivity for
different operators:

LHC tests
 O_9 and O_{11}

and
partially O_7

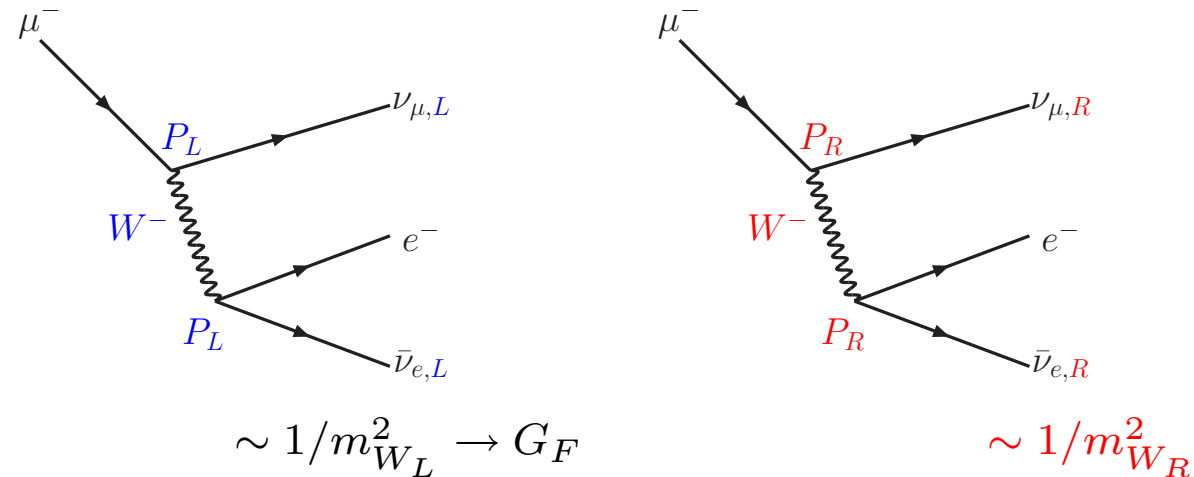


III.

$0\nu\beta\beta$ decay, LNV and LHC

Left-right symmetry

Motivation:



Extend standard model gauge group to:

$$SU(3)_c \times SU(2)_L \times U(1)_Y \rightarrow SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

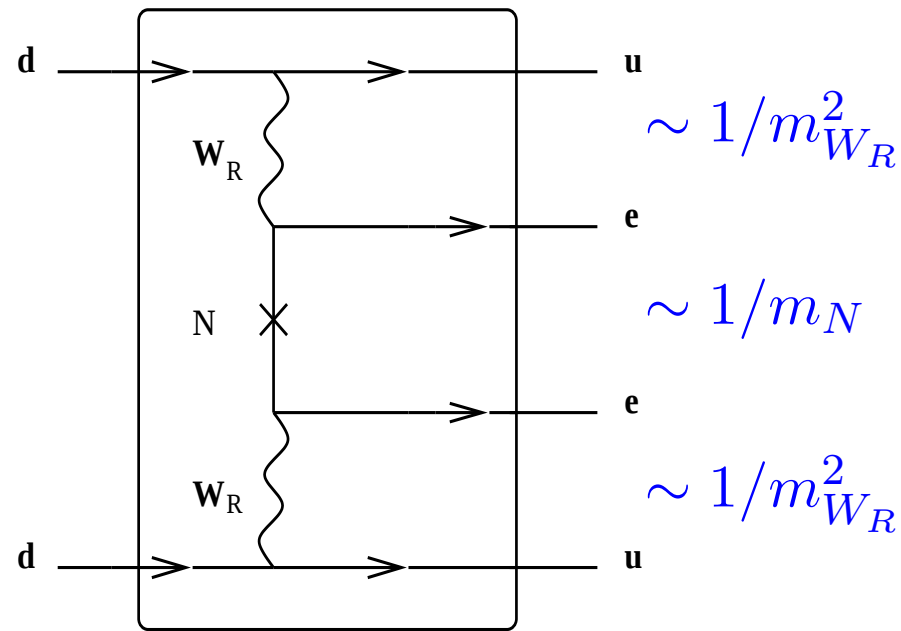
If $m_{W_R} \gg m_{W_L}$ - interactions mostly left-handed

$\Rightarrow$ LR symmetry implies: $L \leftrightarrow L^c$ - ν_R is part of theory!

$\Rightarrow$ Seesaw mechanism included in theory

W_R and $0\nu\beta\beta$ decay

The experimentalist sees:



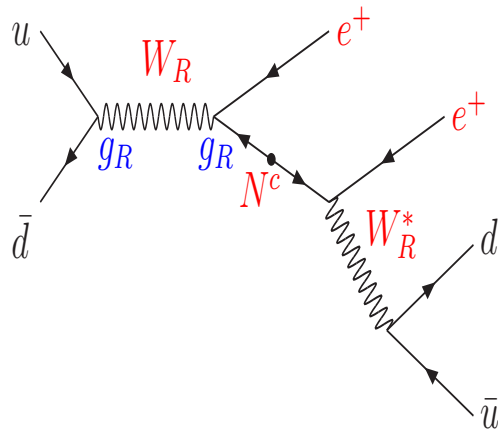
If $m_N \rightarrow \infty$
limit on

$m_{W_R} \rightarrow 0$

With $T_{1/2}^{0\nu\beta\beta}(^{136}\text{Xe}) \geq 1.6 \times 10^{25}$ ys:

$$m_{W_R} \gtrsim 1.3 \left(\frac{\langle m_N \rangle}{[1\text{TeV}]} \right)^{-1/4} \text{TeV}$$

Example: W_R @ LHC

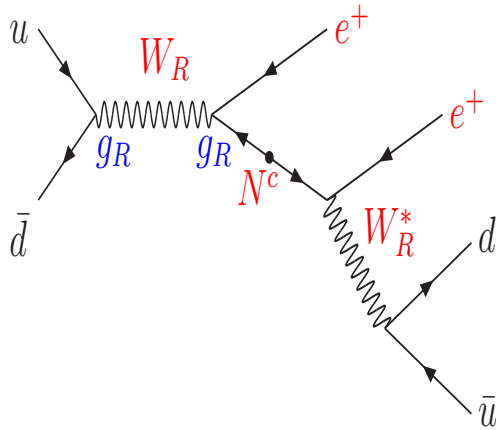


Keung & Senjanovic, 1983

Signal:

Same-sign and opposite-sign
di-lepton + jets, **no** $\cancel{E}_T$

Example: W_R @ LHC



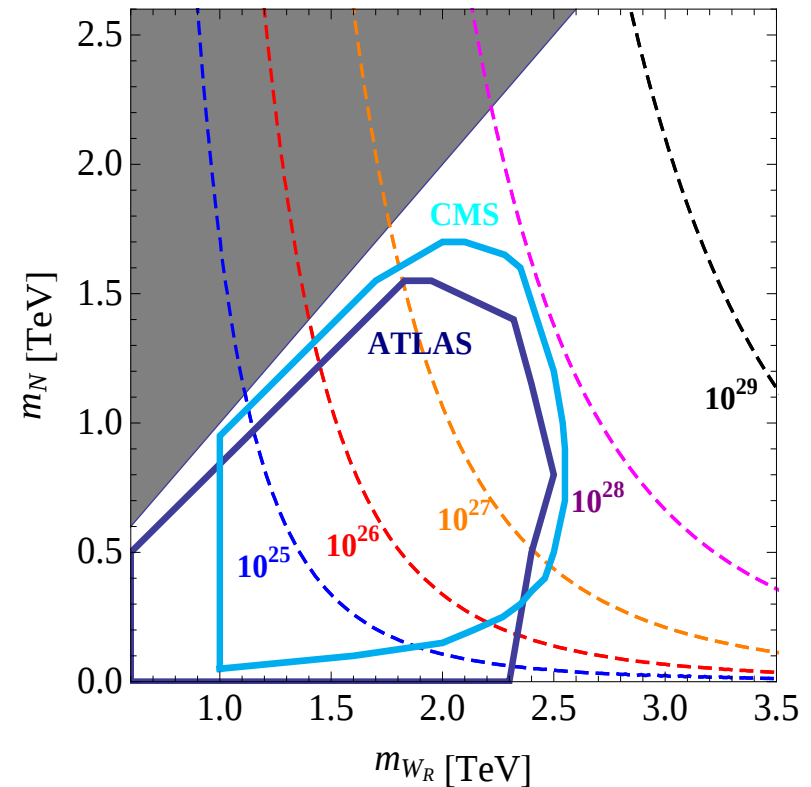
CMS (and ATLAS) with $\sqrt{s} = 8$ TeV:

Non-observation gives
stringent limits on
short-range W_R diagrams
for $0\nu\beta\beta$ decay.

Assumes: $g_R = g_L$!

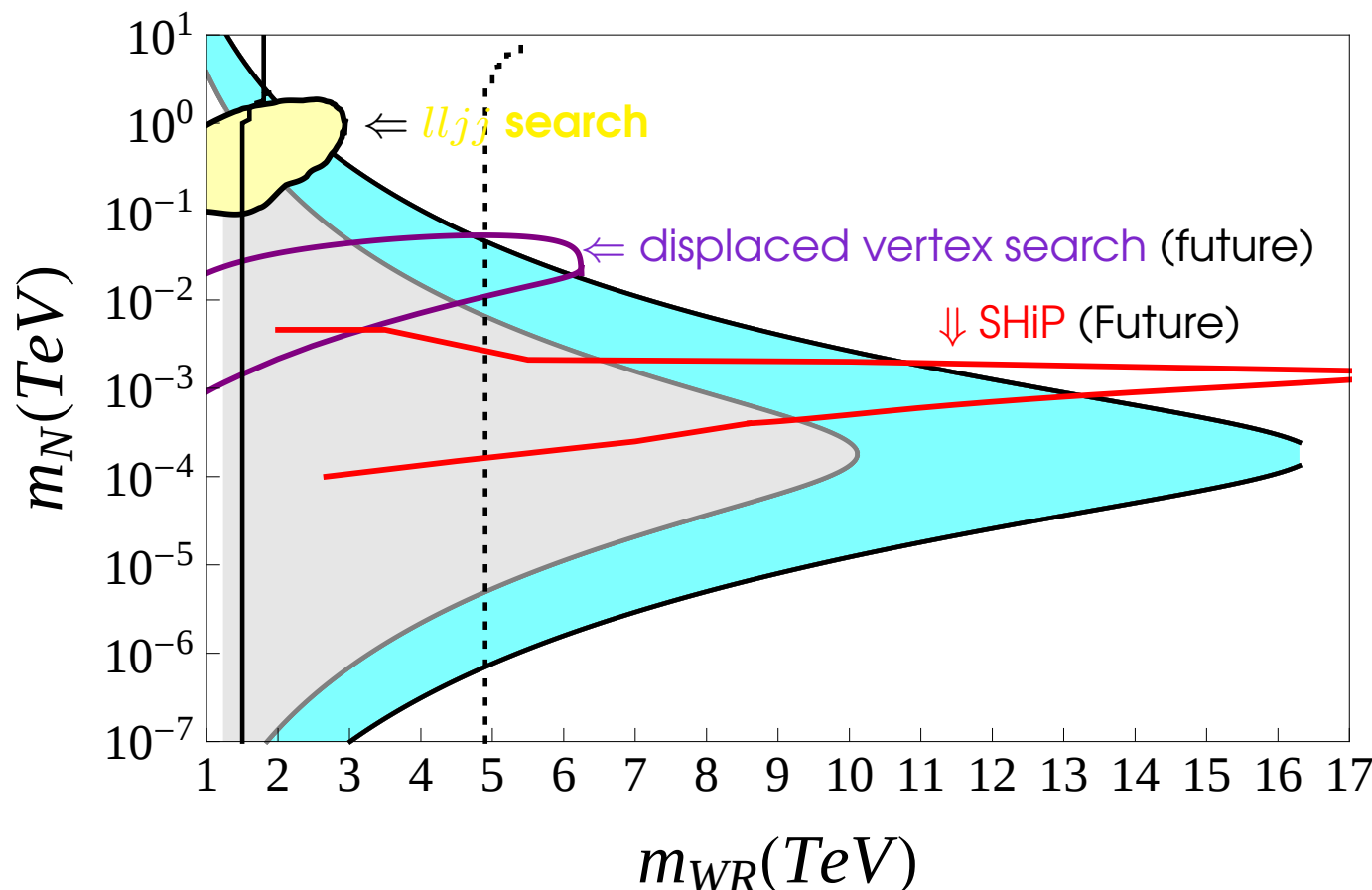
Signal:

Same-sign and opposite-sign
di-lepton + jets, **no** $\cancel{E}_T$



Dijet and $0\nu\beta\beta$ decay

Absence of $pp \rightarrow W_R \rightarrow jj$ gives limit on $0\nu\beta\beta$:



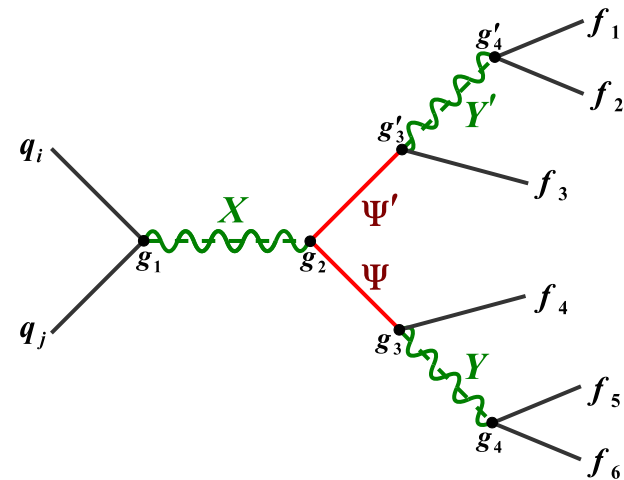
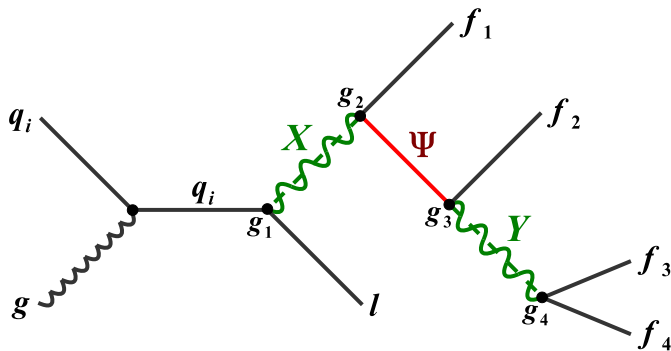
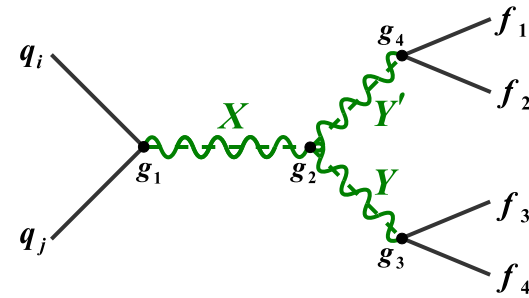
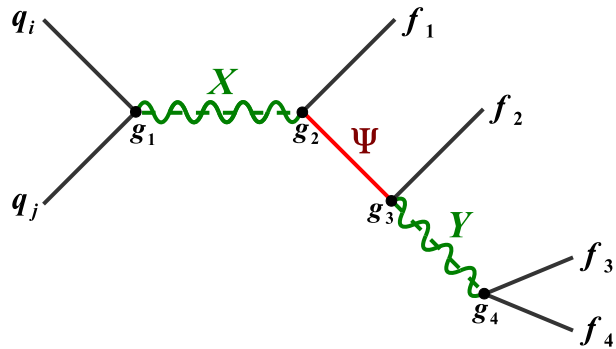
Helo & Hirsch
1509.00423

gray:
 $T_{1/2}^{0\nu\beta\beta} \gtrsim 10^{25} \text{ ys}$

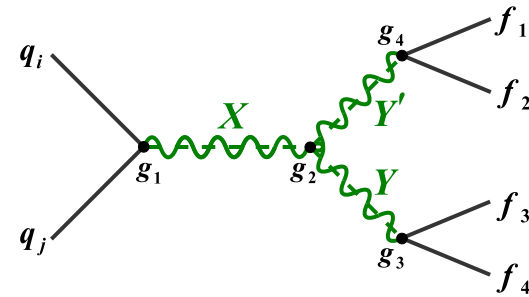
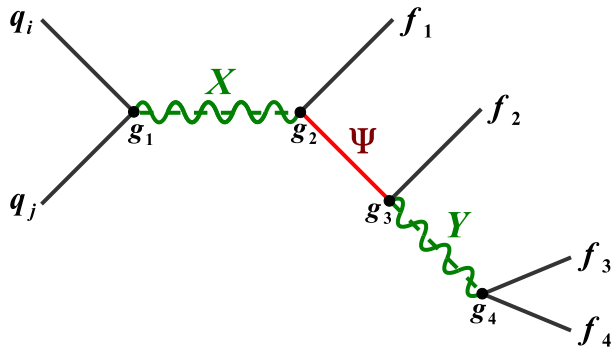
cyan:
 $T_{1/2}^{0\nu\beta\beta} \gtrsim 10^{27} \text{ ys}$

⇒ full (dashed) current limits
(future sensitivity) for dijet data

LNV @ LHC

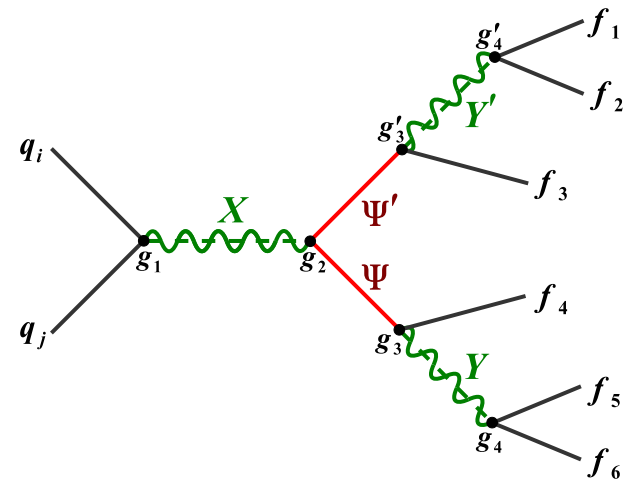
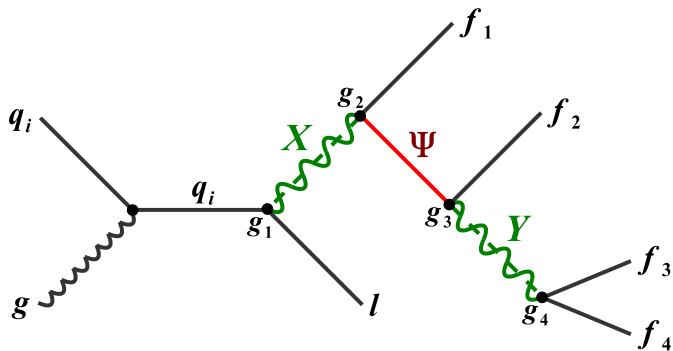


LNV @ LHC

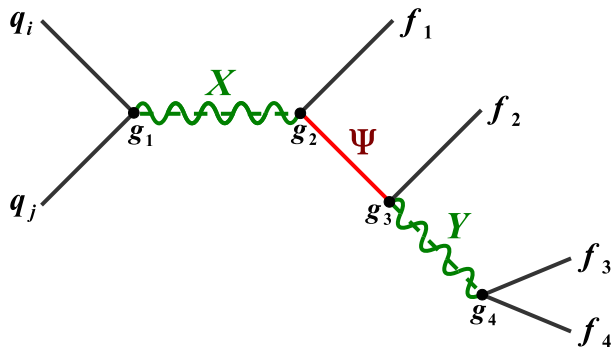


Example:

$$u\bar{d} \rightarrow W_R^+ \rightarrow l^+ N \rightarrow l^+ l^+ jj$$

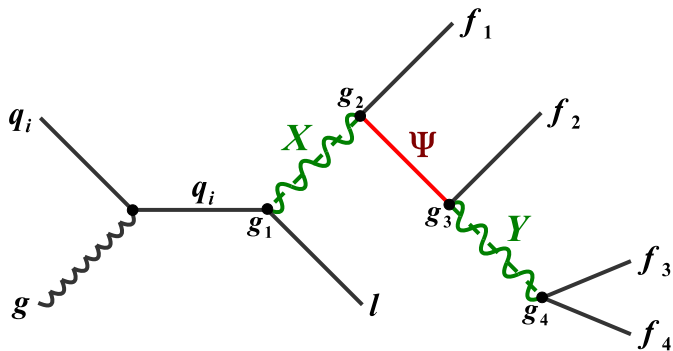


LNV @ LHC

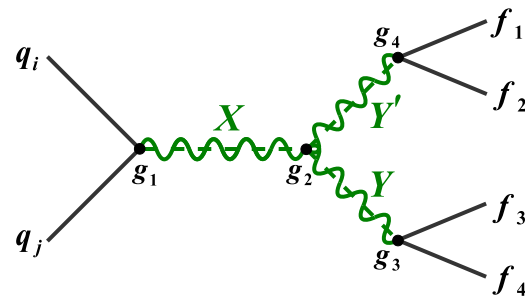


Example:

$$u\bar{d} \rightarrow W_R^+ \rightarrow l^+ N \rightarrow l^+ l^+ jj$$

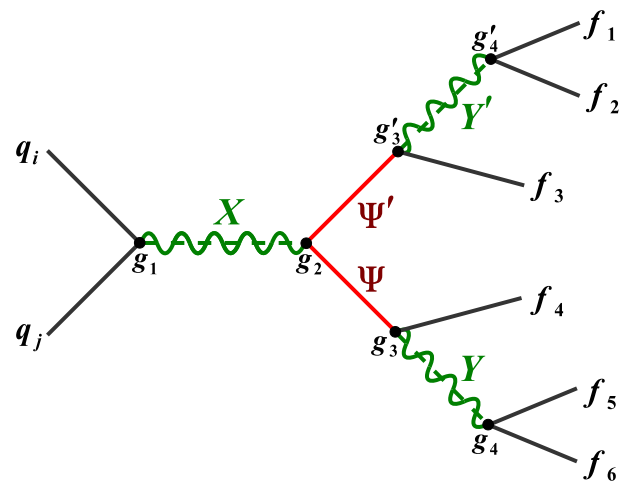


$$ug \rightarrow S_{3,1,1/3} + l^+ \rightarrow l^+ l^+ jjj$$



Example:

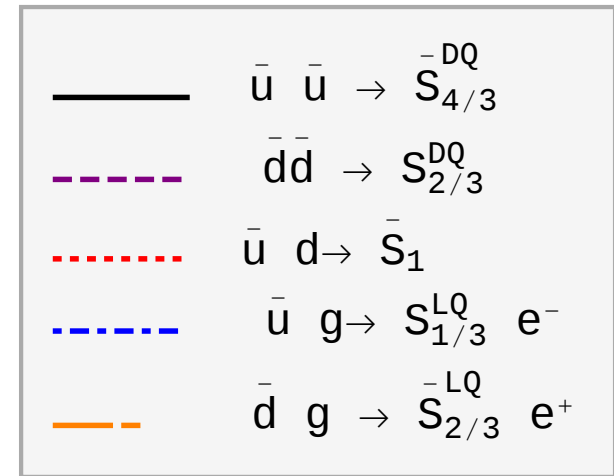
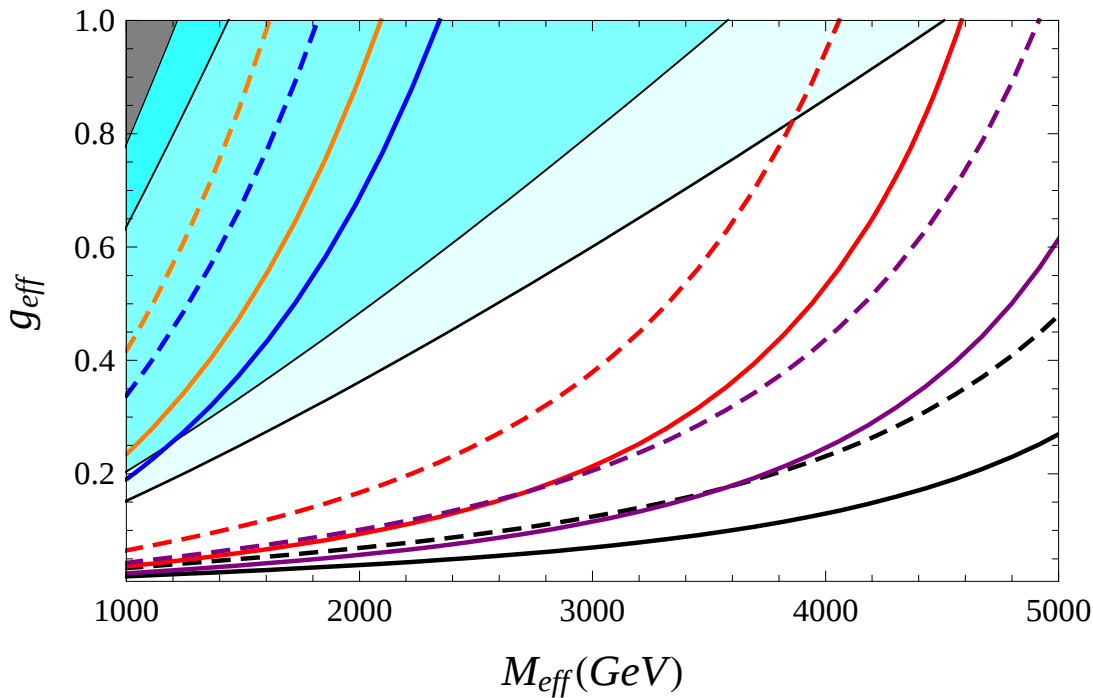
$$uu \rightarrow S_{6,3,1/3} \rightarrow 2S_{3,2,1/6} \rightarrow l^+ l^+ jj$$



$$q\bar{q} \rightarrow g \rightarrow \psi_{6,2,1/6} + \bar{\psi}_{6,2,1/6} \rightarrow l^+ l^+ jjjj$$

$0\nu\beta\beta$ and LHC ($\sqrt{s} = 14 \text{ TeV}$)

J.C. Helo et al,
PRD88 (2013)



g_{eff} - mean coupling
 M_{eff} - mean mass

⇒ Assumed upper limit on $\sigma(pp \rightarrow X)$: 10^{-2} fb

⇒ $m_F = 1000 \text{ GeV}$ (realistic (?) case)

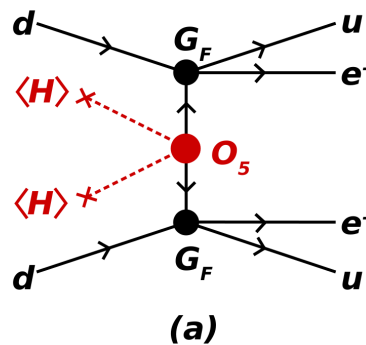
⇒ Full lines: $\text{Br} = 10^{-1}$, dashed lines $\text{Br} = 10^{-2}$

Conclusions

LVN & $0\nu\beta\beta$ decay:

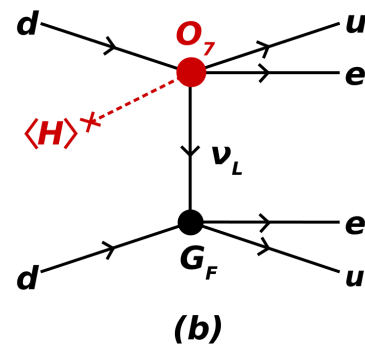
⇒ Majorana neutrino mass and $0\nu\beta\beta$ decay always related

⇒ What is the **scale** of LVN?



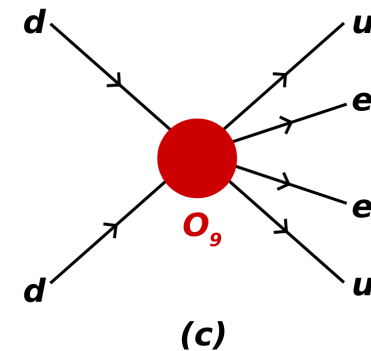
Mass mechanism

near GUT scale ?



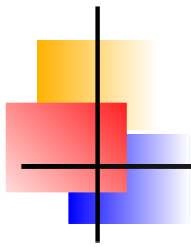
"long-range"

$(10^3 - 10^6)$ GeV?

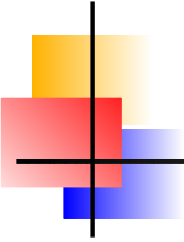


"short-range"

"few" TeV?



Backup slides

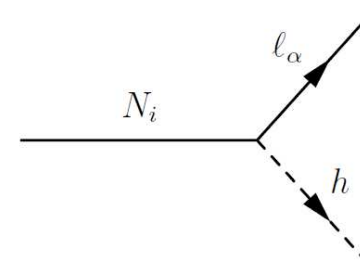


Leptogenesis, $0\nu\beta\beta$ and LHC

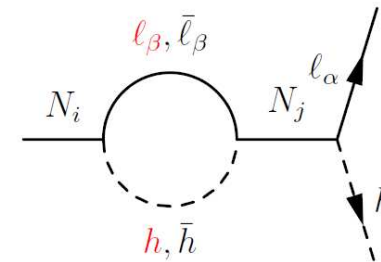
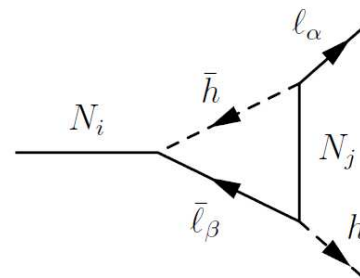
Leptogenesis

Sakharov's conditions:

- (i) Baryon number violation
- (ii) C and CP violation
- (iii) departure from thermal equilibrium



(e) Tree

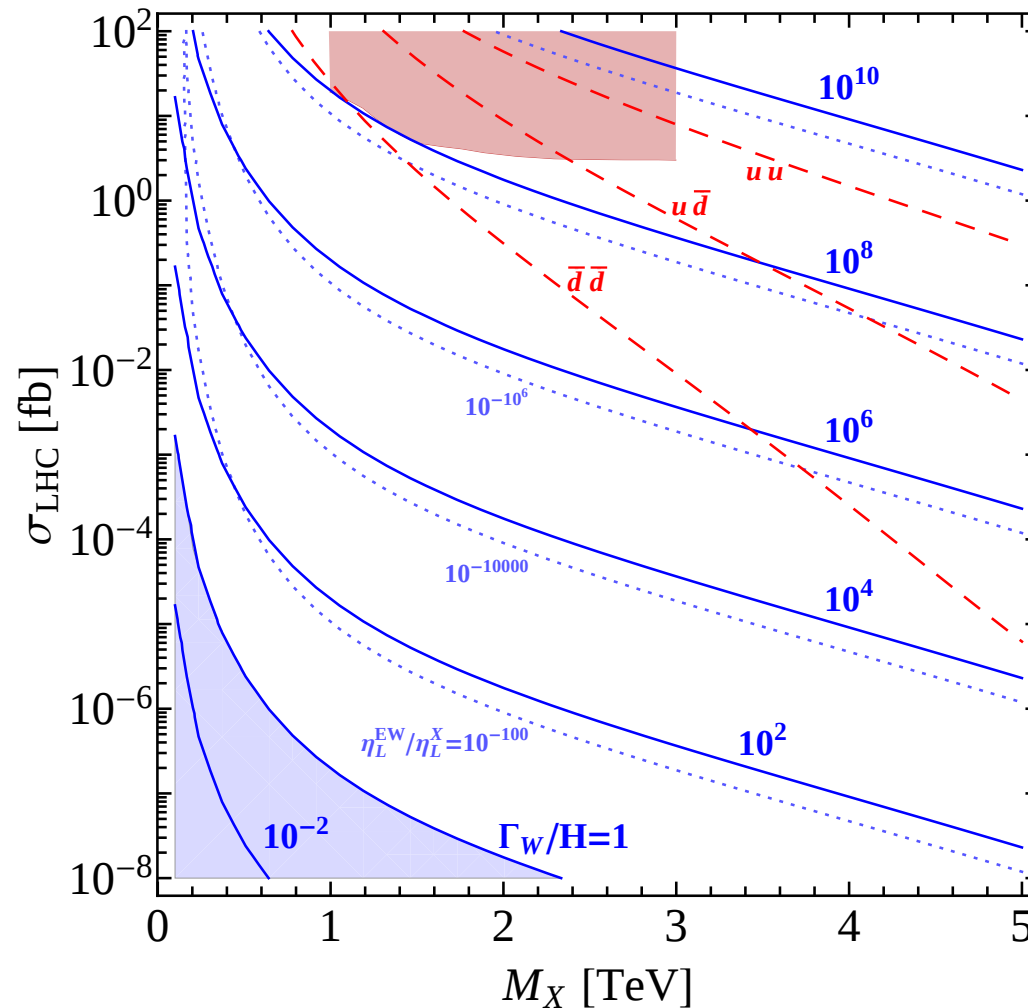


In Leptogenesis:

- (i) Convert L to B through SM sphalerons
- (ii) CP violation through interference tree $\leftrightarrow$ 1-loop
- (iii) L out of equilibrium via right-handed neutrino decay

Leptogenesis and LHC

Deppisch, Hartz & Hirsch
arXiv:1312.4447
PRL 112 (2014)



blue lines
washout factor Γ_W
- Suppression of $L \propto 10^{-\Gamma_W}$

Observation of
LNV @ LHC implies:
(High-scale) Leptogenesis
is ruled out!

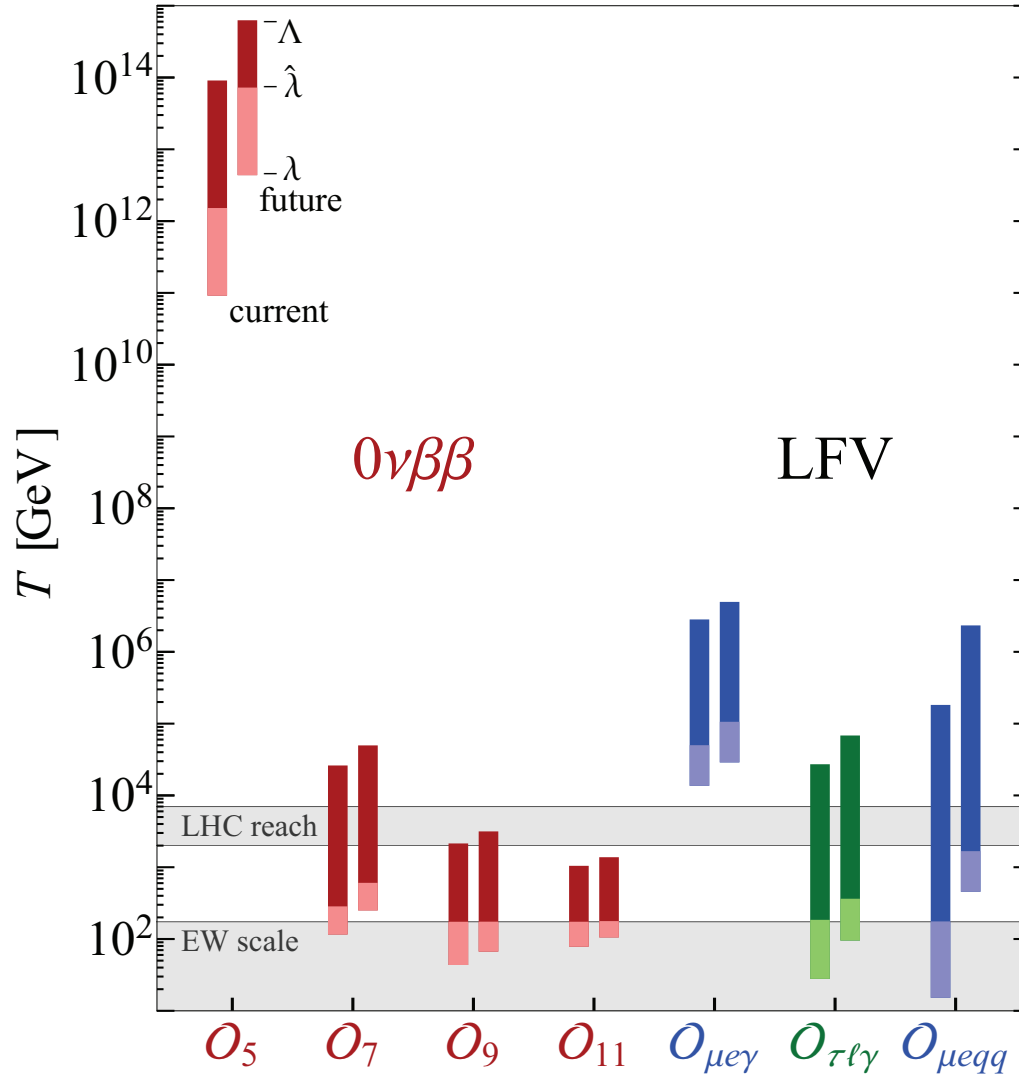
Loopholes???

(i) Resonant LG
with $m_N \ll m_X$?

(ii) Hide LG in τ 's?

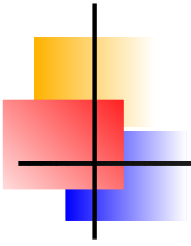
$$\sigma_{\text{LHC}} = \sigma_{pp \rightarrow l^\pm l^\pm + jj}$$

LG and $0\nu\beta\beta$ decay



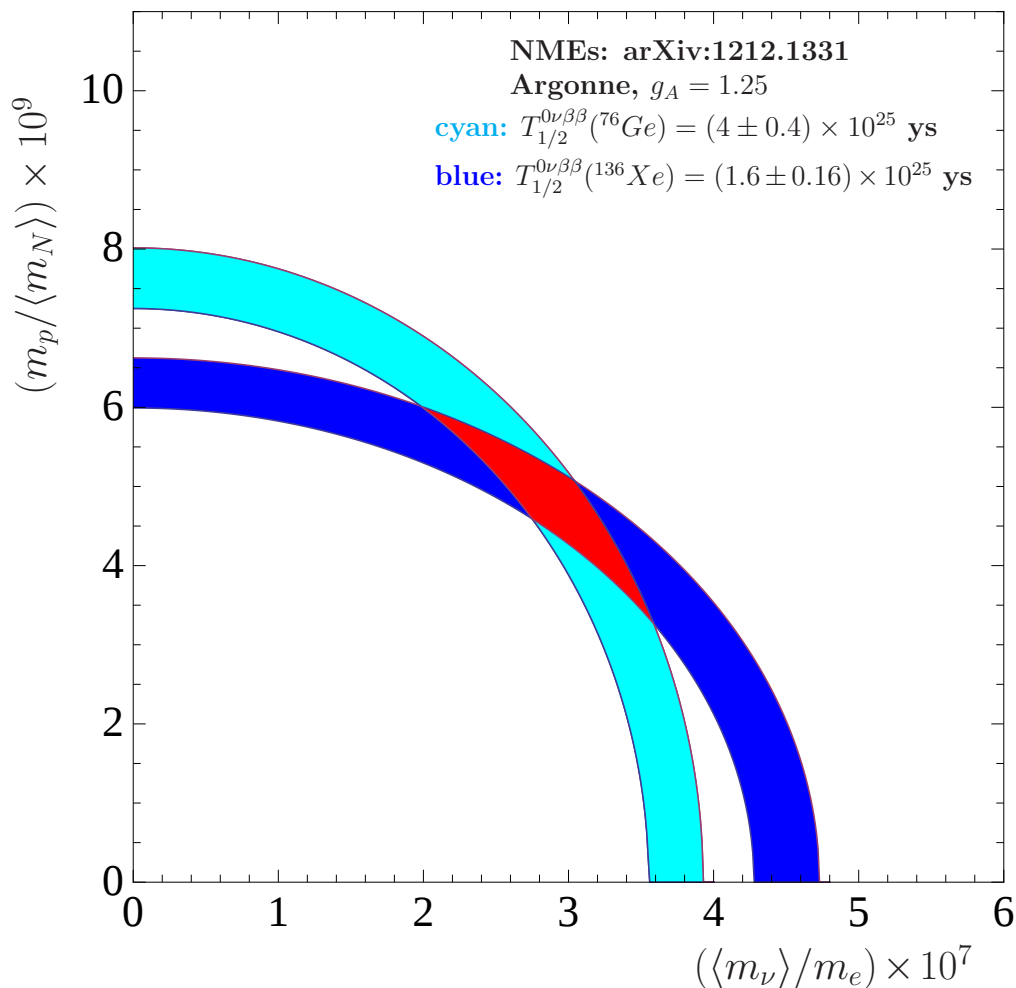
Deppisch et al.,
2015

If $0\nu\beta\beta$ is found
and demonstrated to be
not due to $\langle m_\nu \rangle$
LG ruled out above
scale λ



Comments on using different isotopes

More than one isotope?



Simple idea:

$$T_{1/2} = G \left\{ (\epsilon_1 \mathcal{M}_1)^2 + (\epsilon_2 \mathcal{M}_2)^2 \right\}$$

Only one example:

$$\epsilon_1 = (\langle m_\nu \rangle / m_e)$$

$$\epsilon_2 = (m_p / \langle m_N \rangle)$$

$\mathcal{M}_i$ - matrix elements

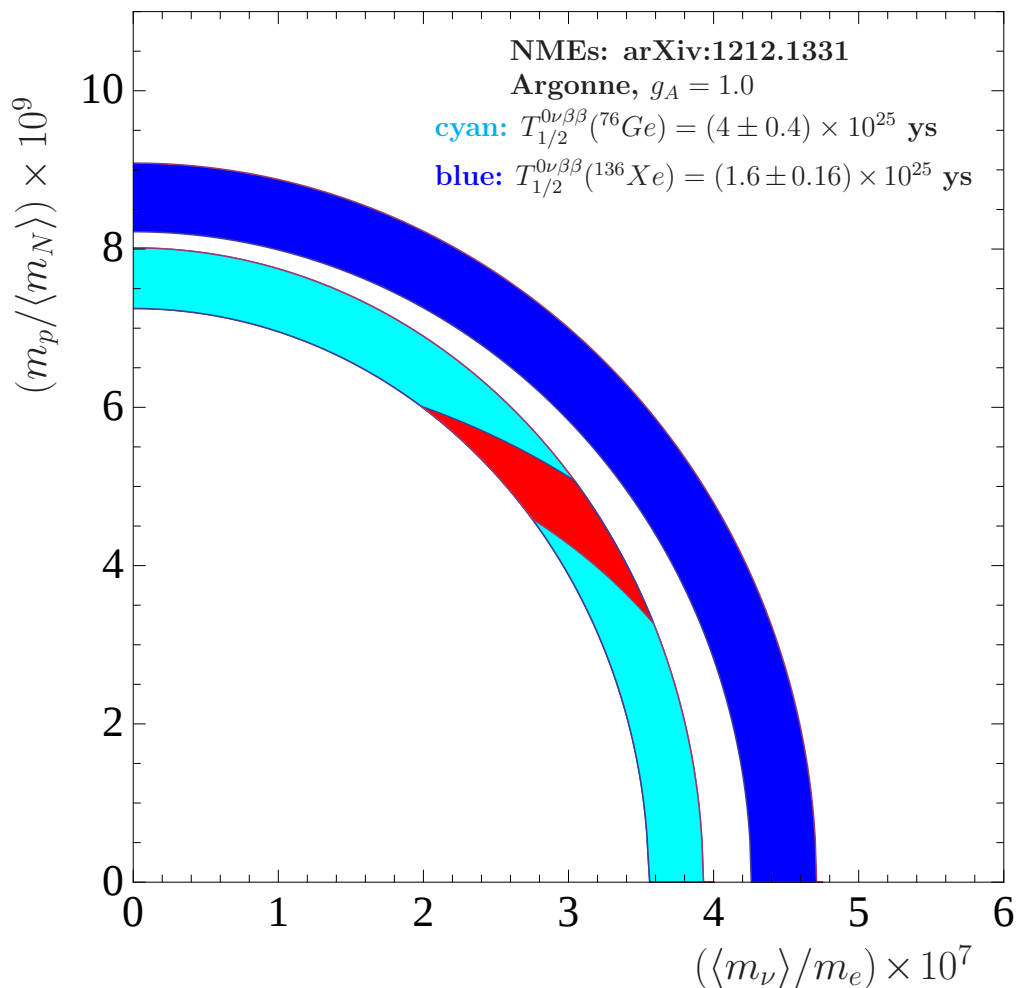
Use two (or more) isotopes
solve for two unknowns ...

Consistency region
shown in red

Note:

Plot is not based on experimental data! For illustration only.

More than one isotope?



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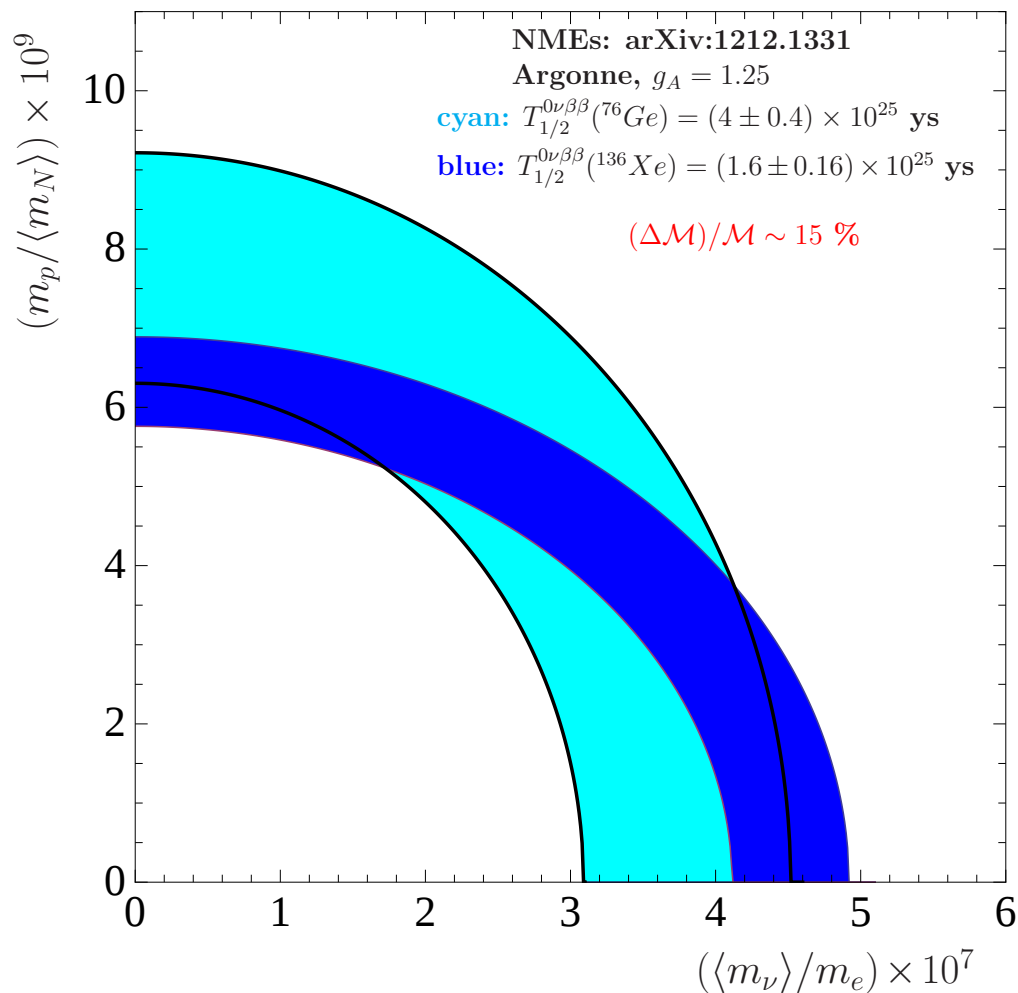
Use two (or more) isotopes
solve for two unknowns ...

No overlap for exactly
same numbers, except:
 $\mathcal{M}_i!$

Note:

Plot is not based on experimental data! For illustration only.

More than one isotope?



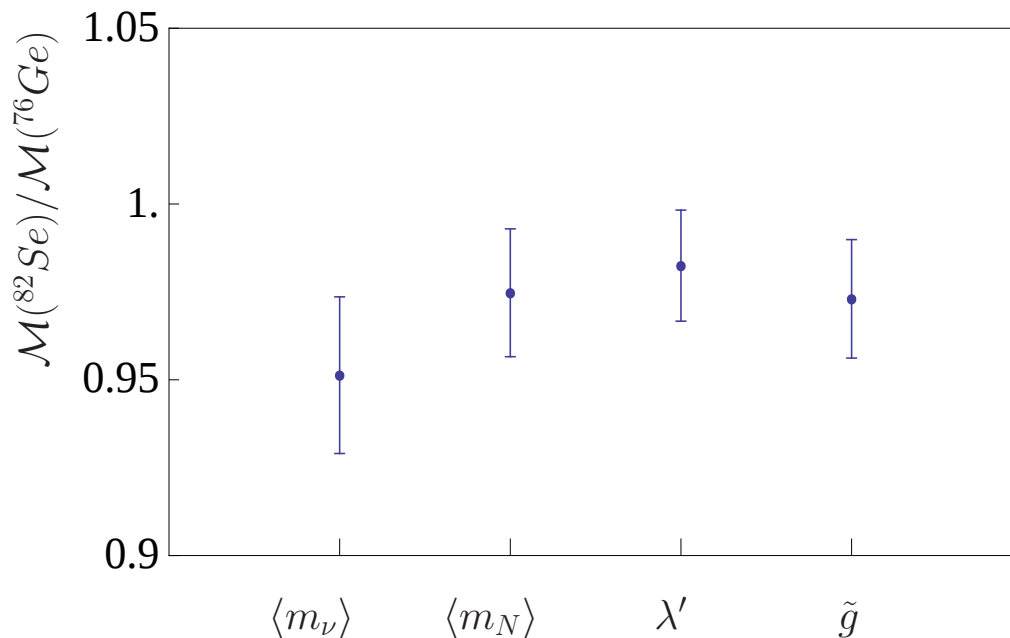
Same numbers except:

Arbitrary error for:
 $(\Delta\mathcal{M})/\mathcal{M} \sim 20 \%$ for ${}^{76}\text{Ge}$

Different mechanisms in
different nuclei lead to
(5-25) % differences,
much smaller than
typical error ($\Delta\mathcal{M}$)

Conclusion:
Better $\mathcal{M}$ needed!

Ratios of matrix elements



To distinguish:
Need different mechanisms
to have different NME ratios.

“Error bar” due to variation
of NMEs in Table 1 of 1212.1331

Plot shows ratio of matrix elements for different “mechanisms”, see 1212.1331:

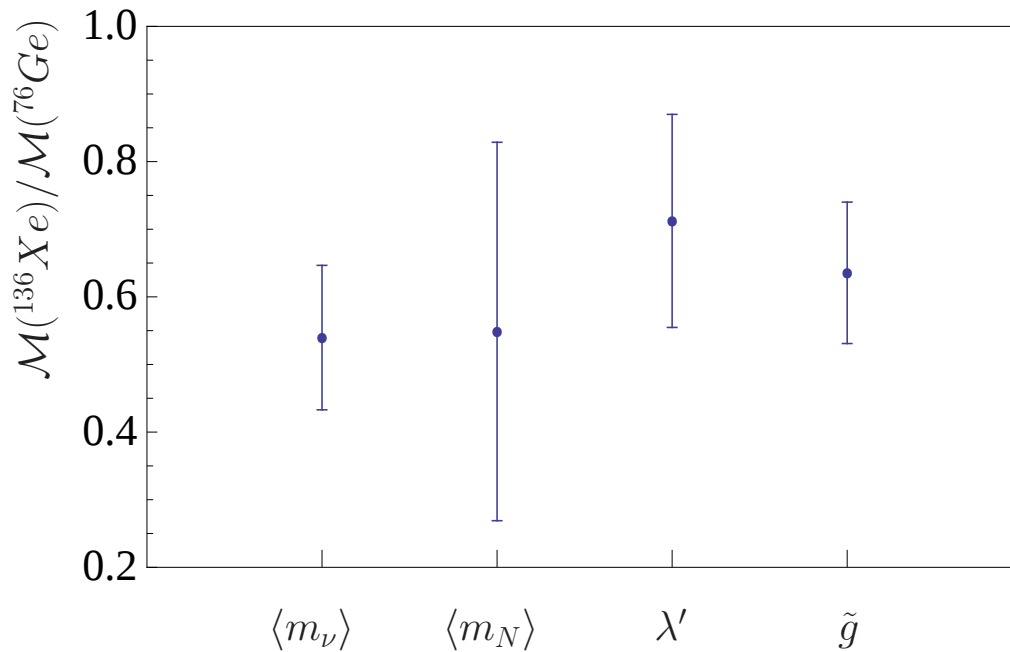
$\langle m_\nu \rangle$ - Mass mechanism

$\langle m_N \rangle$ - heavy neutrino exchange

λ' - “trilinear RPV”

$\tilde{g}$ - gluino RPV

Ratios of matrix elements



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