

Neutrinoless double-beta decay and nuclear theory

Achim Schwenk



TECHNISCHE
UNIVERSITÄT
DARMSTADT



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Bundesministerium
für Bildung
und Forschung



Outline

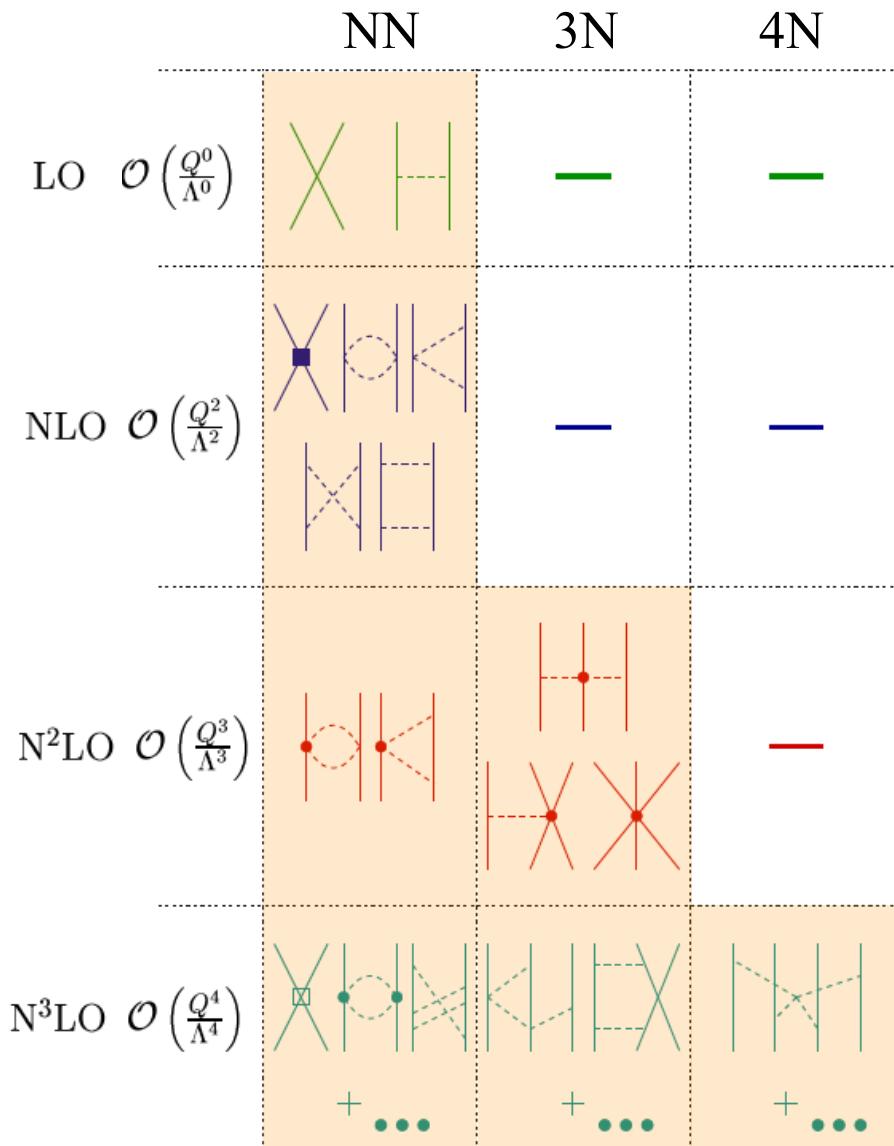
Advances in nuclear forces and nuclear structure

Progress in nuclear matrix elements for
neutrinoless double beta decay

Perspectives

Chiral effective field theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



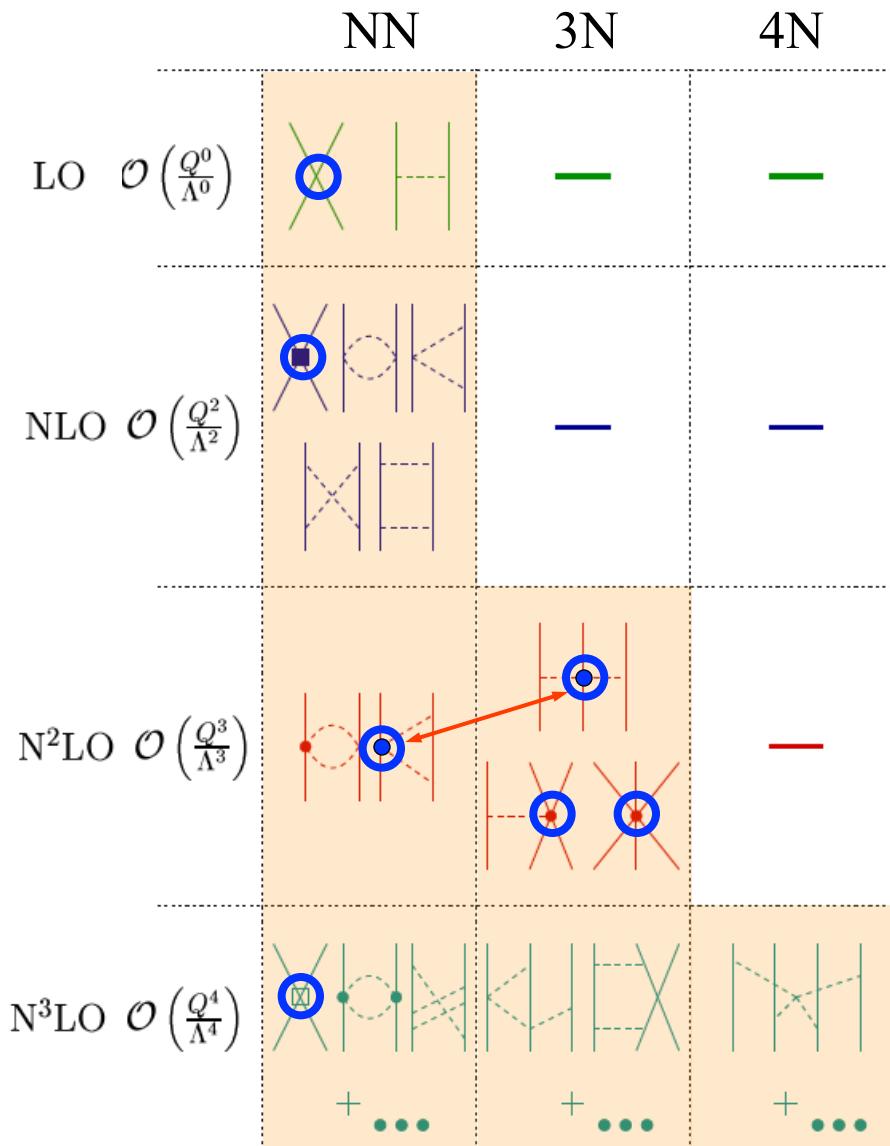
limited resolution at low energies,
can expand in powers $(Q/\Lambda_b)^n$

LO, n=0 - leading order,
NLO, n=2 - next-to-leading order,...
expansion parameter $\sim 1/3$



Chiral effective field theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



include long-range pion physics

few short-range couplings,
fit to experiment once

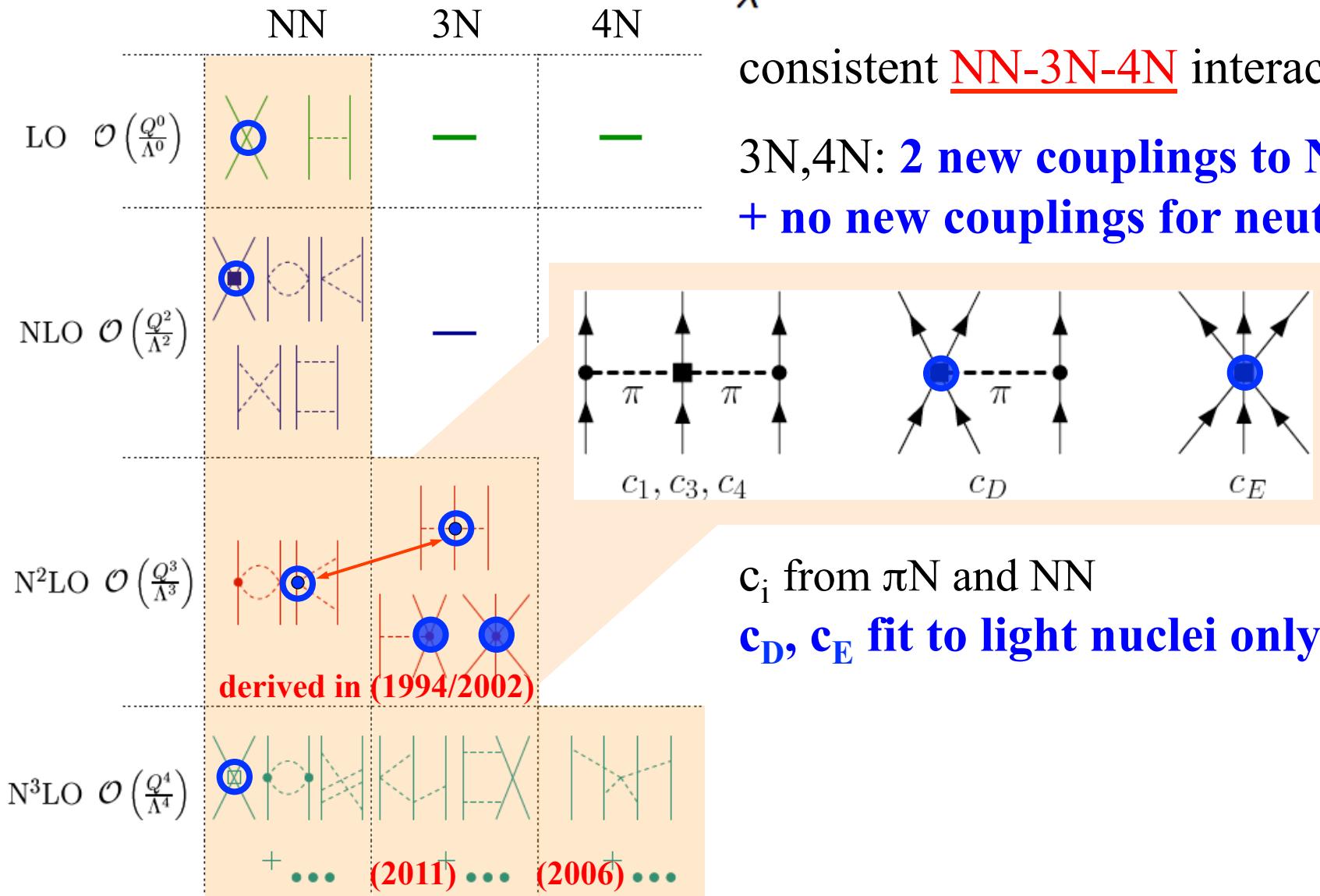
systematic: can work to desired
accuracy and obtain **error estimates**

consistent **electroweak interactions**
and **matching to lattice QCD**

new developments in power counting,
uncertainty quantification,
optimization,... **Ekström et al., Furnstahl et al.**

Chiral effective field theory and **many-body forces**

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



consistent NN-3N-4N interactions

3N,4N: **2 new couplings to N³LO**
+ **no new couplings for neutrons**

c_i from πN and NN

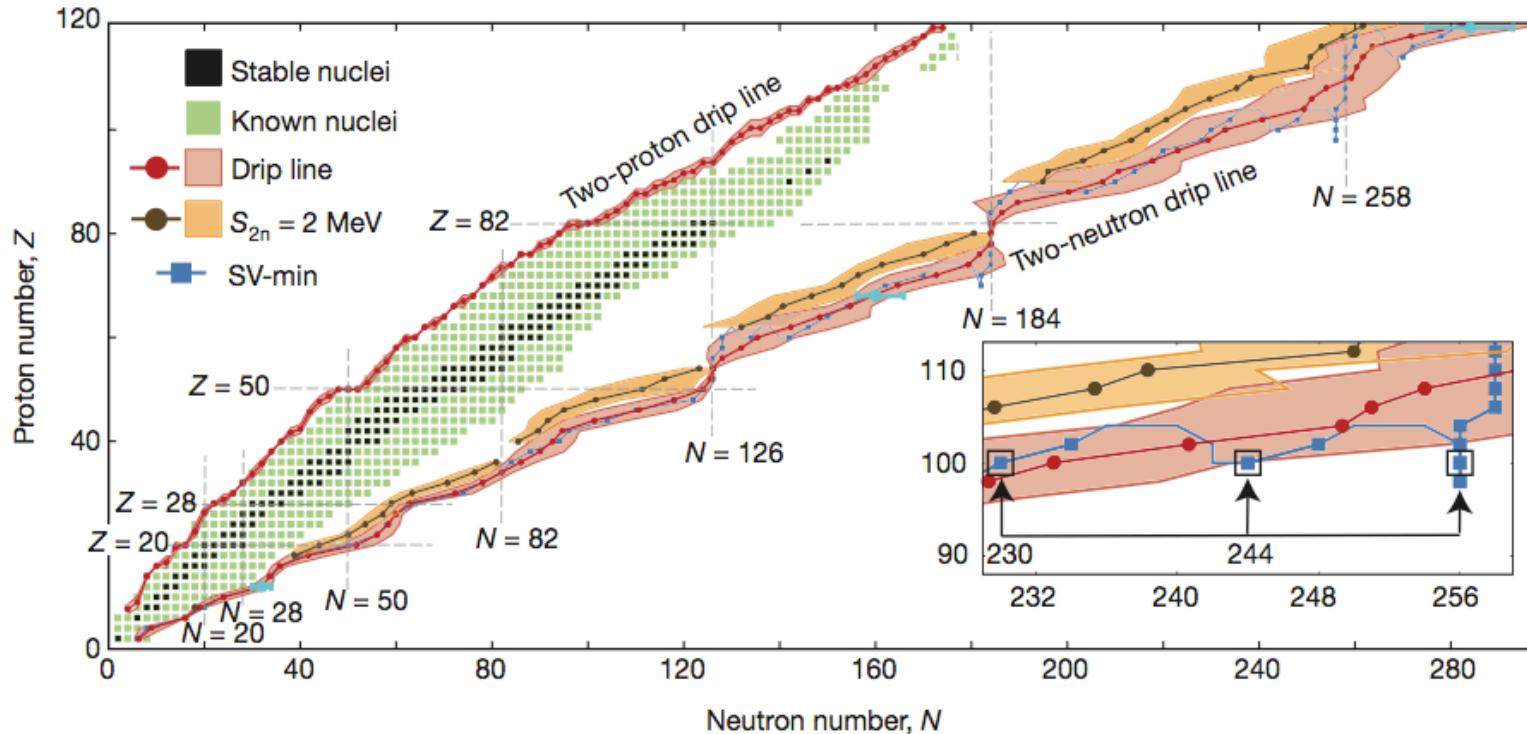
c_D, c_E fit to light nuclei only

Nuclei bound by strong interactions

doi:10.1038/nature11188

The limits of the nuclear landscape

Jochen Erler^{1,2}, Noah Birge¹, Markus Kortelainen^{1,2,3}, Witold Nazarewicz^{1,2,4}, Erik Olsen^{1,2}, Alexander M. Perhac¹ & Mario Stoitsov^{1,2,†}

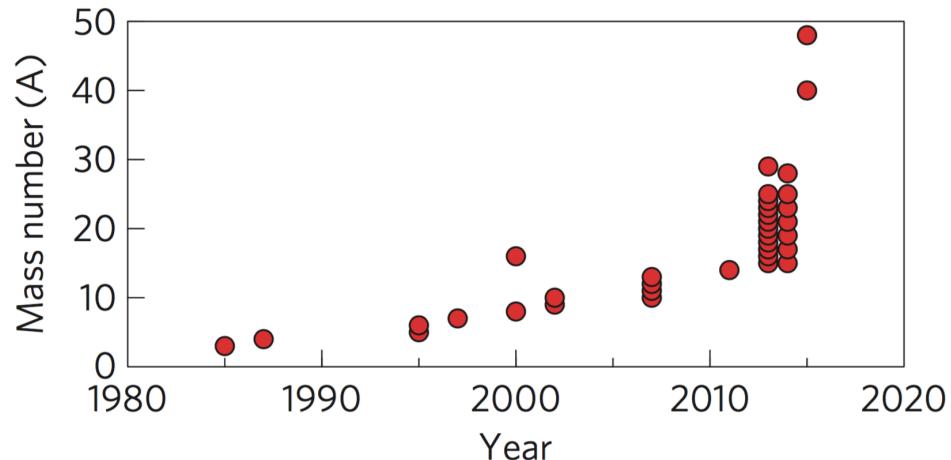


How does the nuclear chart emerge from quantum chromodynamics?

Lattice QCD and effective field theories of QCD

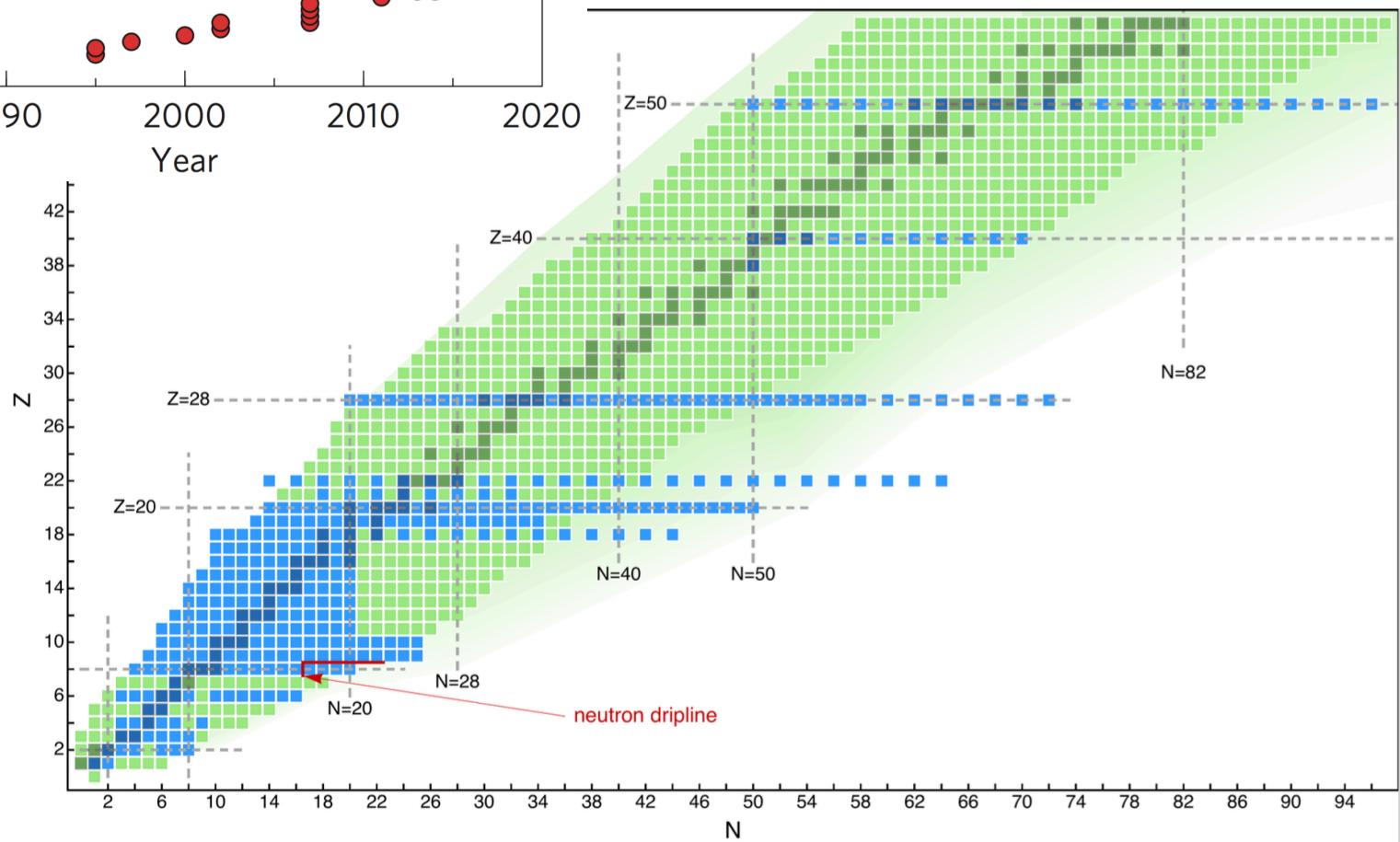
Progress in ab initio calculations of nuclei

dramatic progress in last 5 years to access nuclei up to $A \sim 50$



from Hagen et al., Nature Phys. (2016)

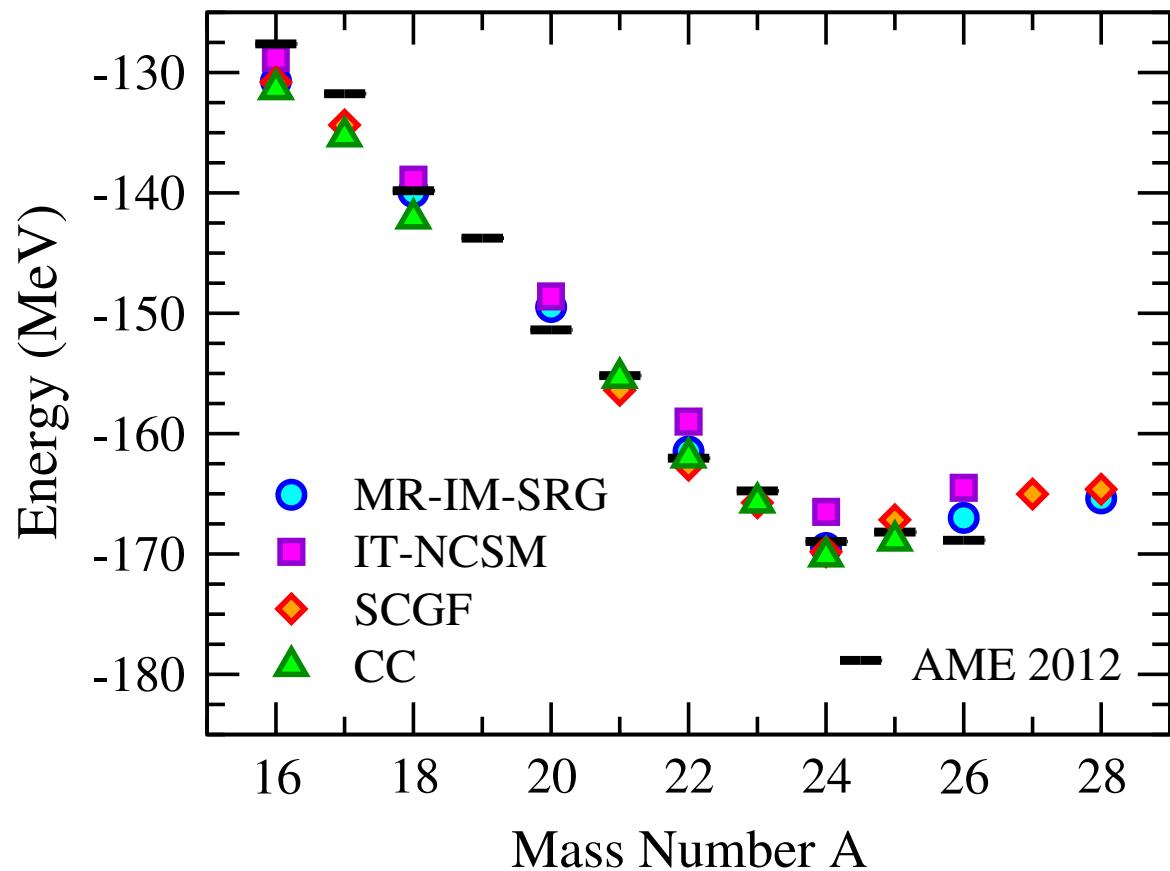
from Hergert et al., Phys. Rep. (2016)



Ab initio calculations of neutron-rich oxygen isotopes

impact of **3N forces key for neutron dripline** Otsuka et al., PRL (2010)

based on **same**
NN+3N interactions



using different many-body methods:

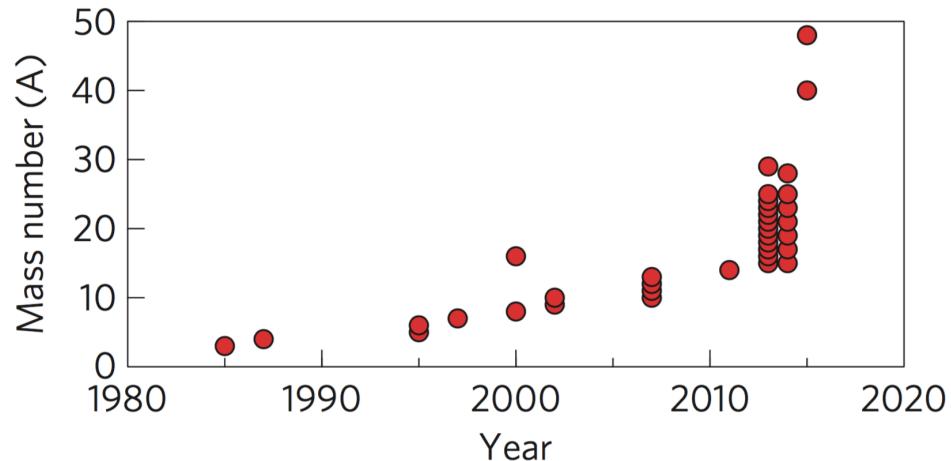
Coupled Cluster theory/CCEI Hagen et al., PRL (2012), Jansen et al., PRL (2014)

Multi-Reference In-Medium SRG and IT-NCSM Hergert et al., PRL (2013)

Self-Consistent Green's Function methods Cipollone et al., PRL (2013)

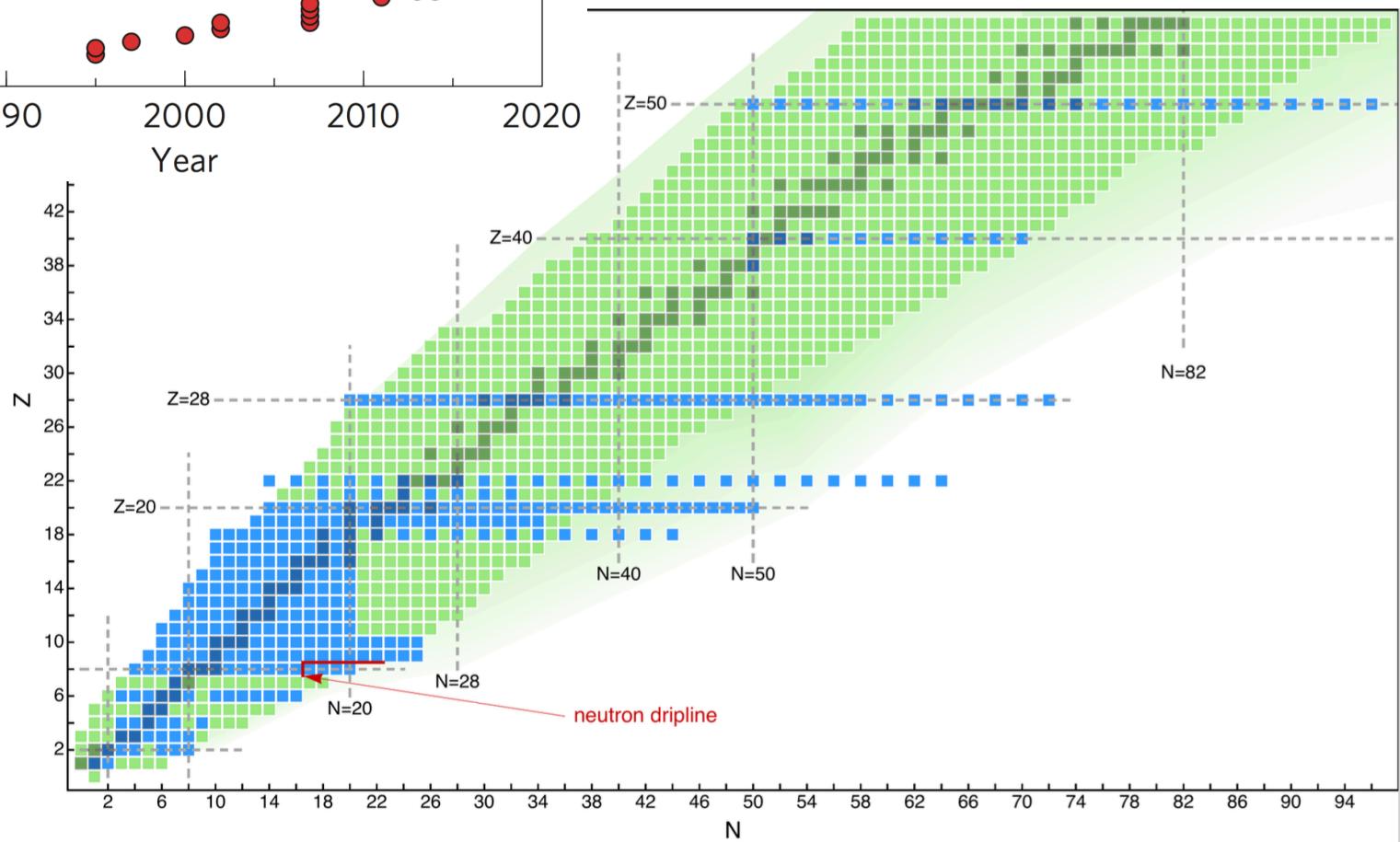
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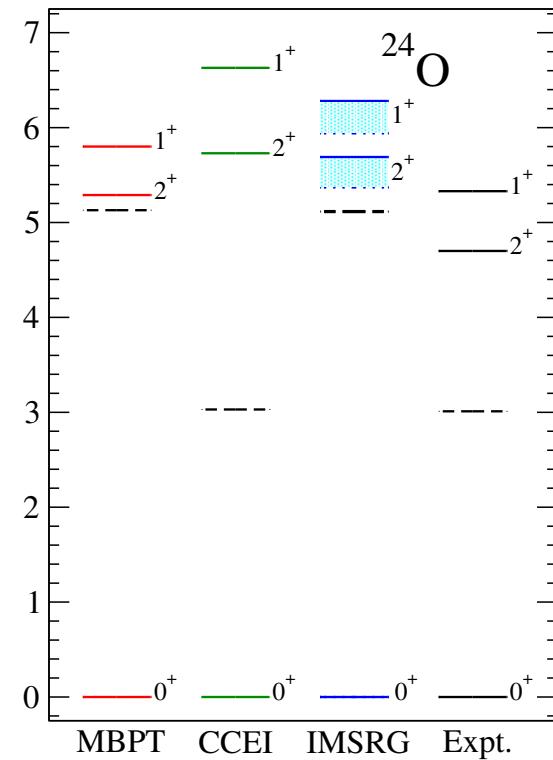
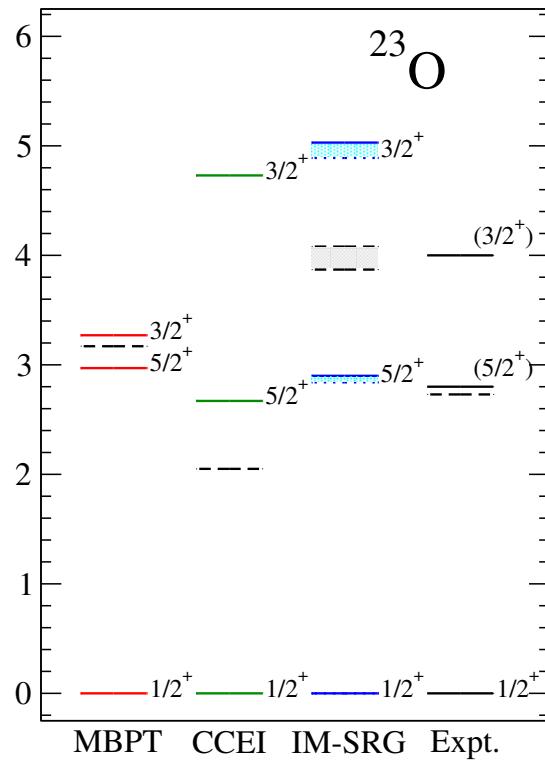
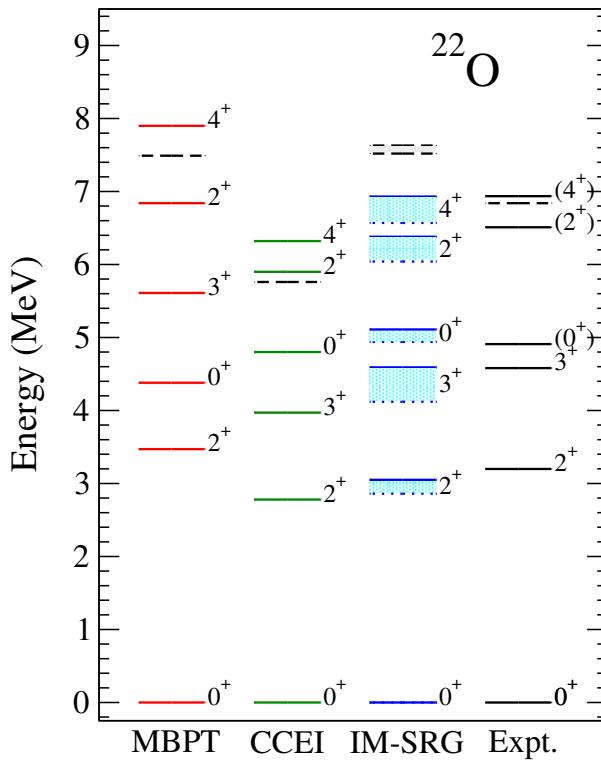


Ab initio calculations going open shell

In-Medium SRG to derive nonperturbative shell-model interactions

Tsukiyama, Bogner, AS, PRC (2012), Bogner et al., PRL (2014)

Coupled Cluster for effective interactions (CCEI) Jansen et al., PRL (2014)



Experiments at GANIL, GSI, NSCL, RIBF: ^{22}O and ^{24}O doubly magic

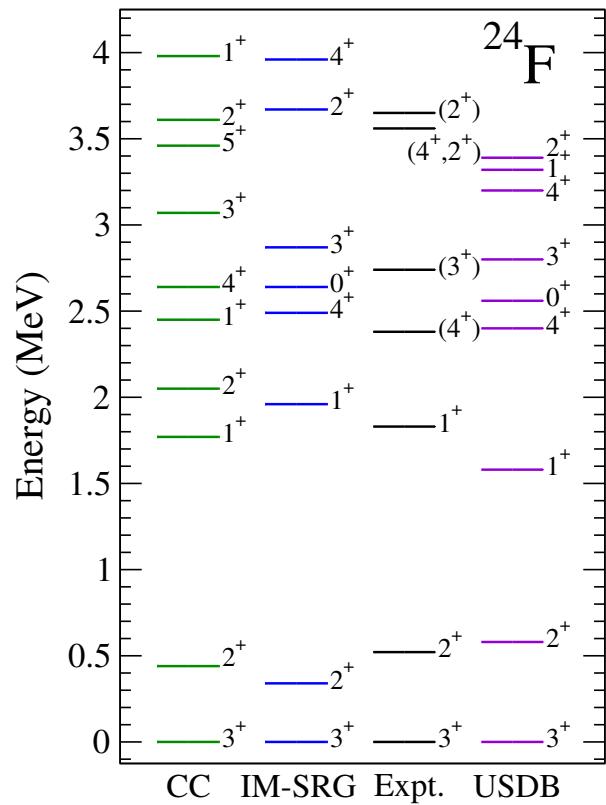
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Spectrum of ^{24}F Cáceres et al., PRC (2015)

vs. new GANIL experiment

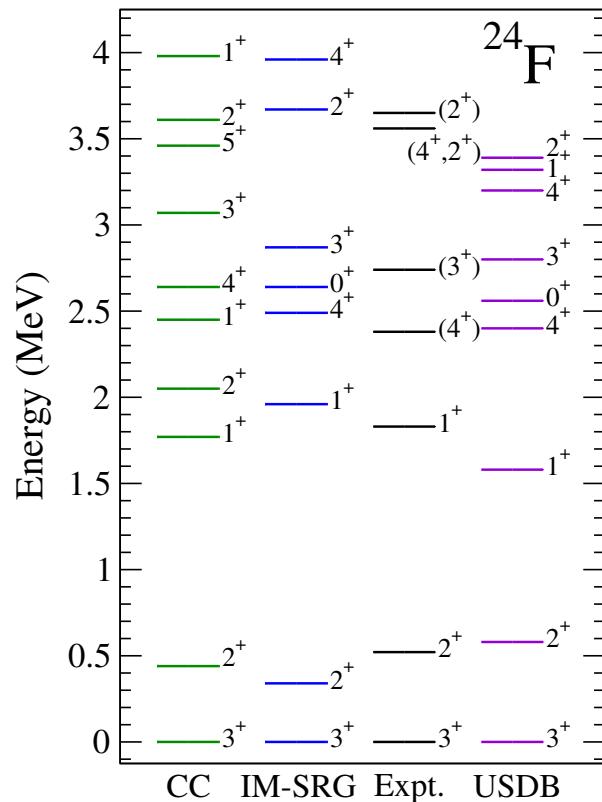


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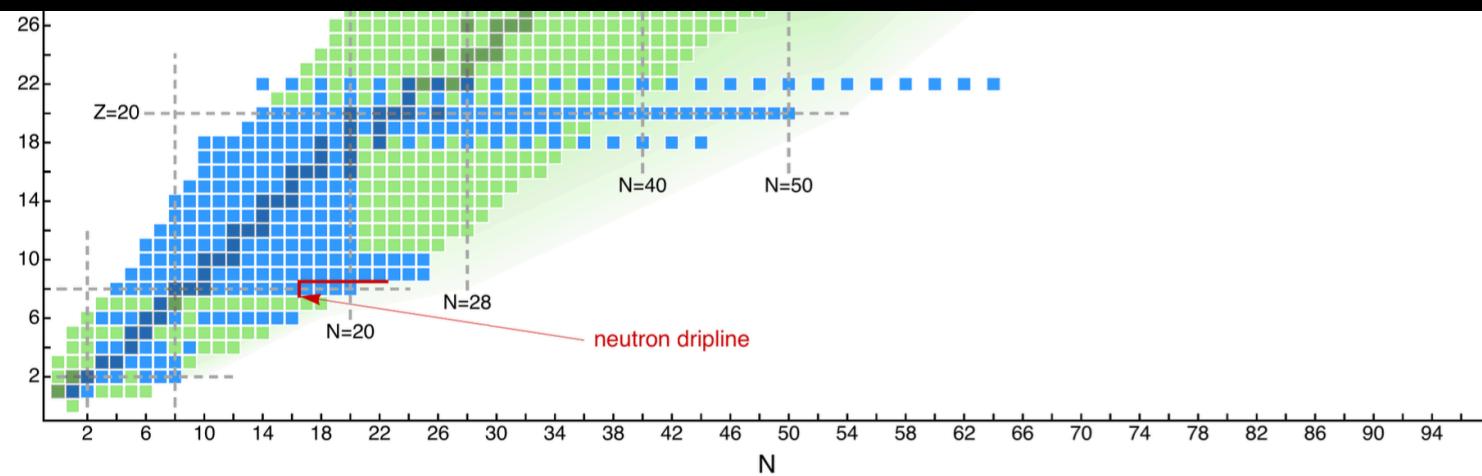
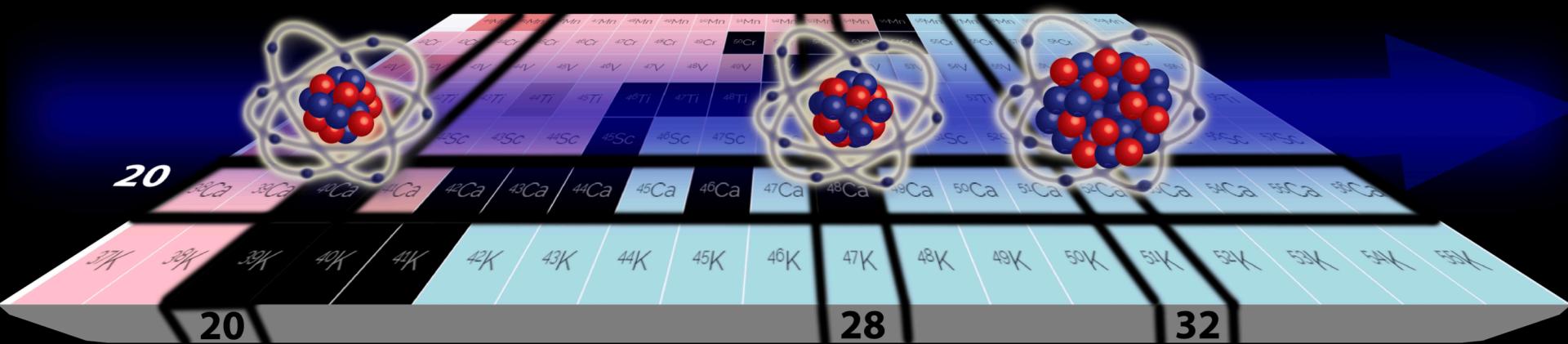
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Neutron-rich calcium isotopes



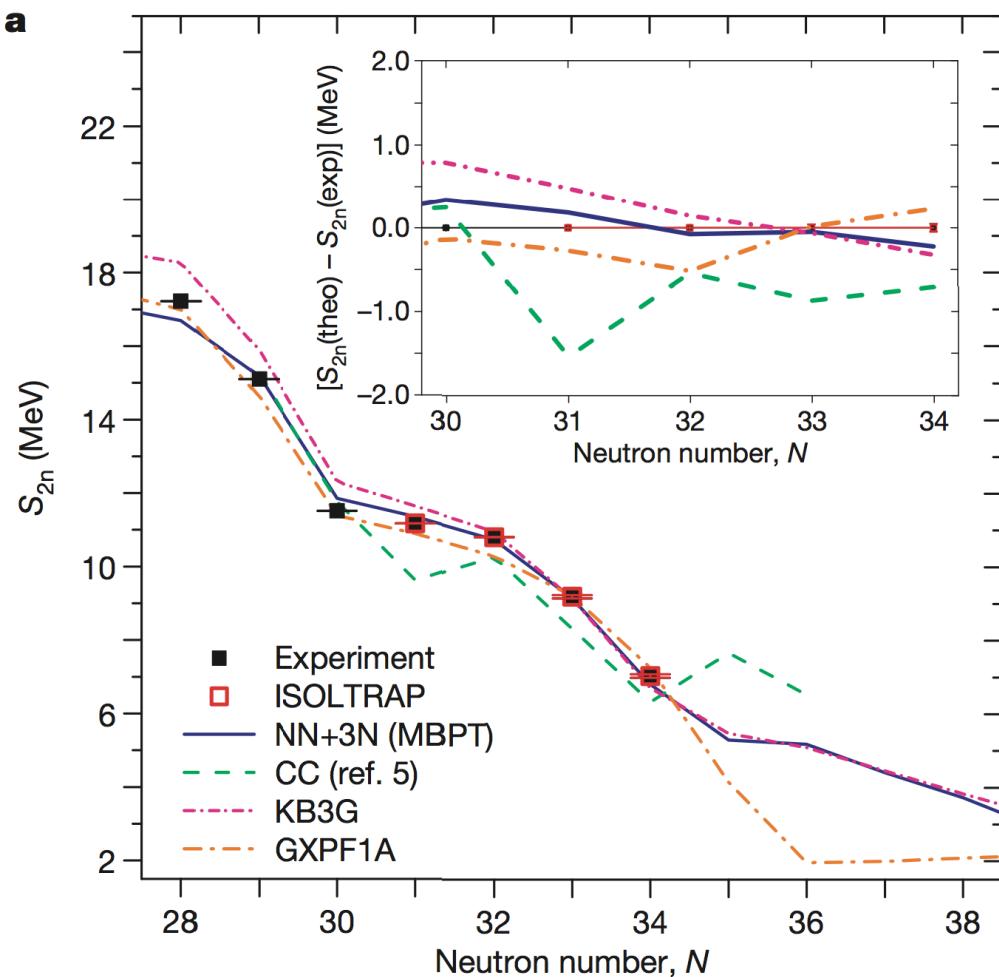
Masses of exotic calcium isotopes pin down nuclear forces

F. Wienholtz¹, D. Beck², K. Blaum³, Ch. Borgmann³, M. Breitenfeldt⁴, R. B. Cakirli^{3,5}, S. George¹, F. Herfurth², J. D. Holt^{6,7}, M. Kowalska⁸, S. Kreim^{3,8}, D. Lunney⁹, V. Manea⁹, J. Menéndez^{6,7}, D. Neidherr², M. Rosenbusch¹, L. Schweikhard¹, A. Schwenk^{7,6}, J. Simonis^{6,7}, J. Stanja¹⁰, R. N. Wolf¹ & K. Zuber¹⁰

$^{53,54}\text{Ca}$ masses measured at ISOLTRAP/CERN using new MR-TOF mass spectrometer

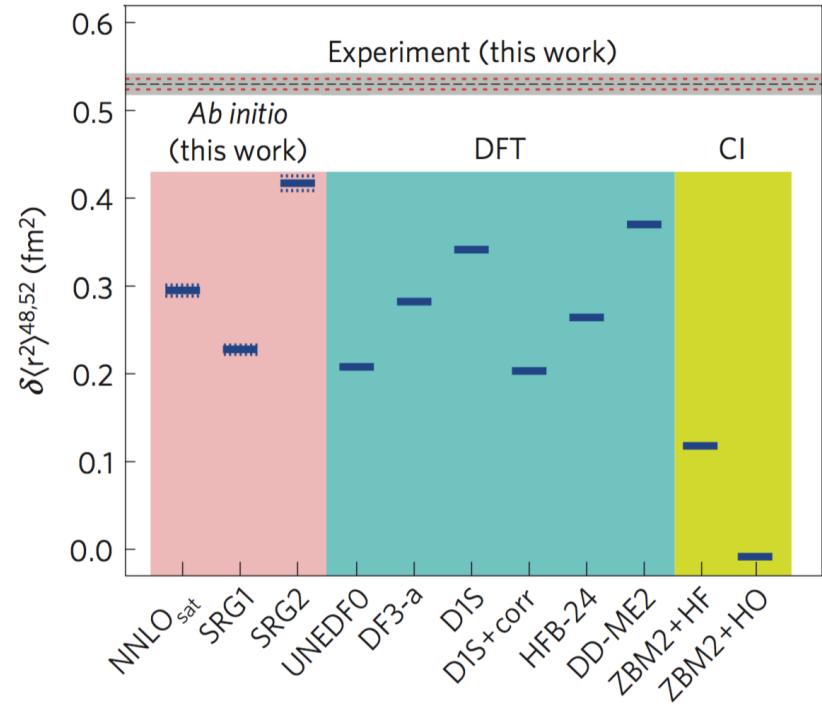
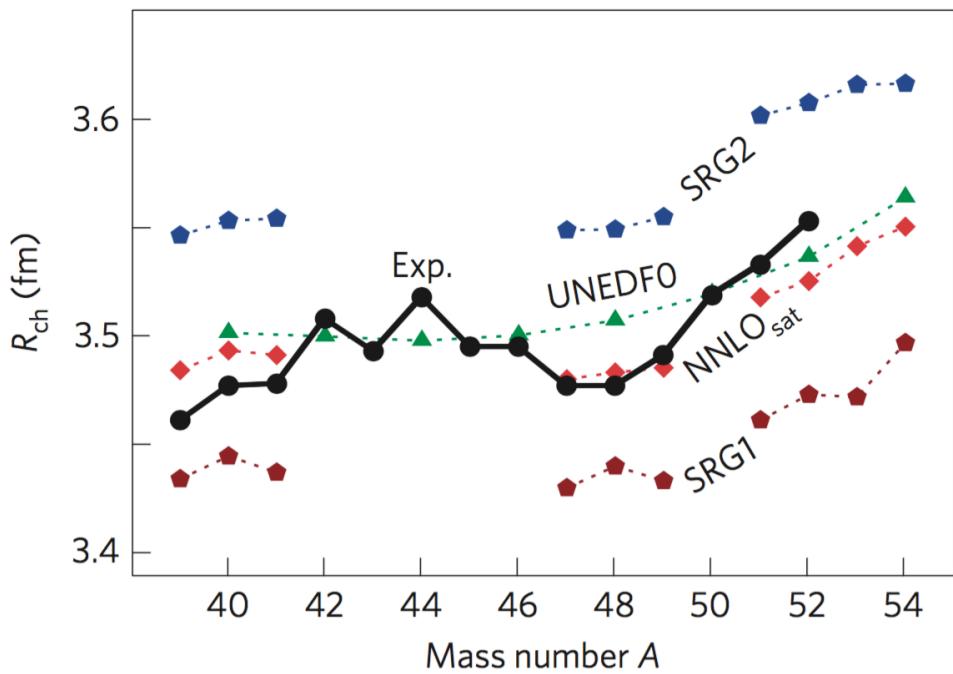
excellent agreement with theoretical NN+3N prediction

suggests $N=32$ shell closure



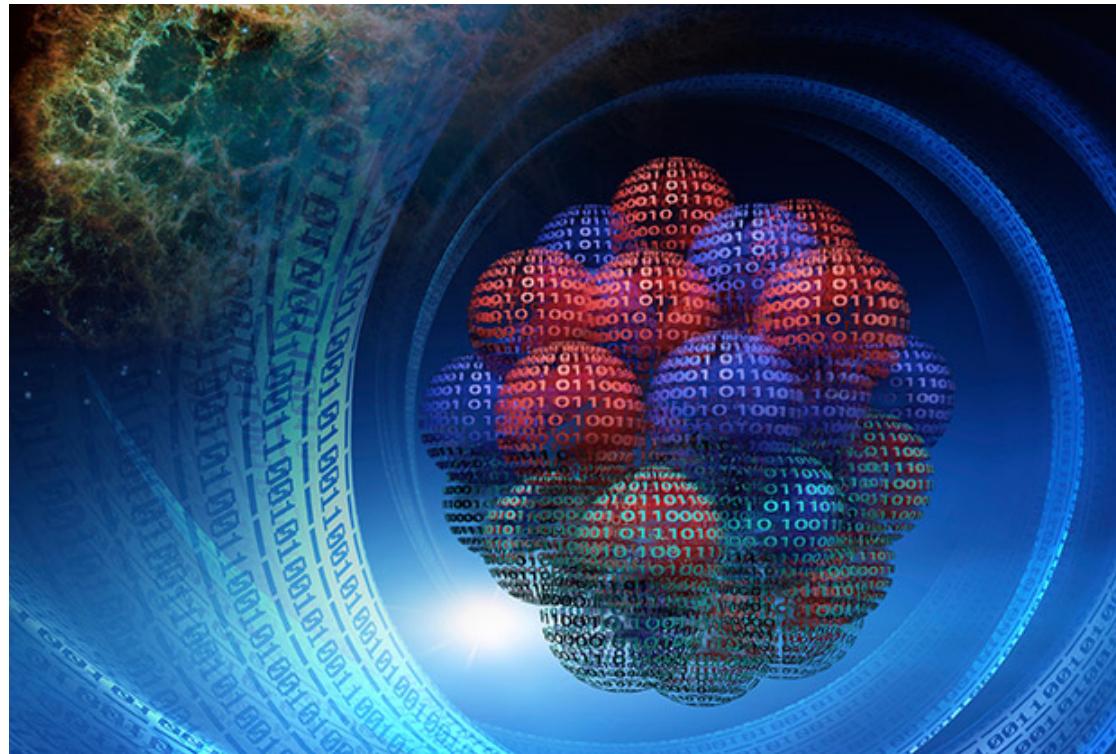
Unexpectedly large charge radii of neutron-rich calcium isotopes

R. F. Garcia Ruiz^{1*}, M. L. Bissell^{1,2}, K. Blaum³, A. Ekström^{4,5}, N. Frömmgen⁶, G. Hagen⁴, M. Hammen⁶, K. Hebeler^{7,8}, J. D. Holt⁹, G. R. Jansen^{4,5}, M. Kowalska¹⁰, K. Kreim³, W. Nazarewicz^{4,11,12}, R. Neugart^{3,6}, G. Neyens¹, W. Nörtershäuser^{6,7}, T. Papenbrock^{4,5}, J. Papuga¹, A. Schwenk^{3,7,8}, J. Simonis^{7,8}, K. A. Wendt^{4,5} and D. T. Yordanov^{3,13}



Neutron and weak-charge distributions of the ^{48}Ca nucleus

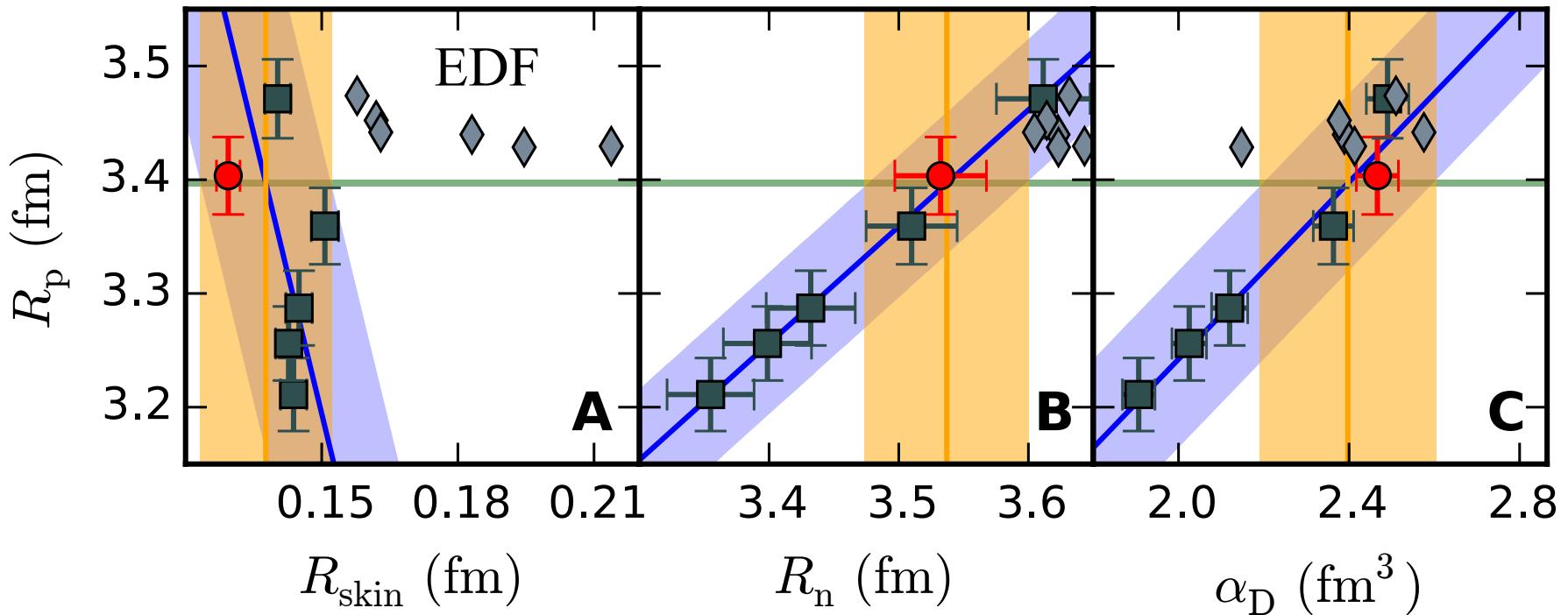
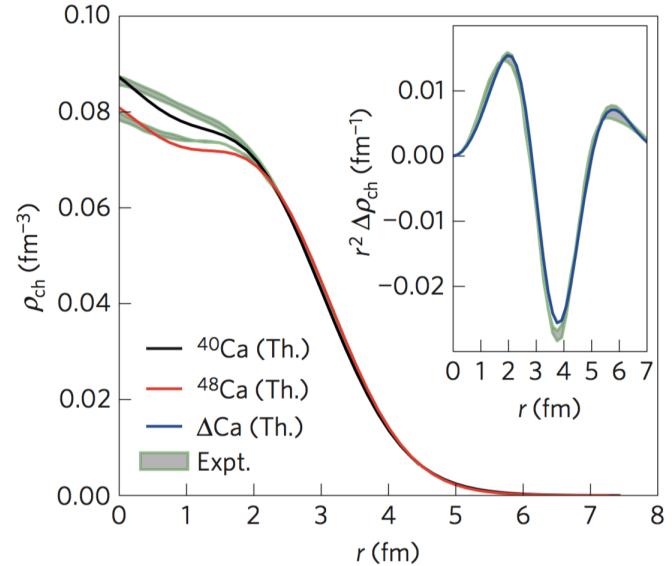
G. Hagen^{1,2*}, A. Ekström^{1,2}, C. Forssén^{1,2,3}, G. R. Jansen^{1,2}, W. Nazarewicz^{1,4,5}, T. Papenbrock^{1,2}, K. A. Wendt^{1,2}, S. Bacca^{6,7}, N. Barnea⁸, B. Carlsson³, C. Drischler^{9,10}, K. Hebeler^{9,10}, M. Hjorth-Jensen^{4,11}, M. Miorelli^{6,12}, G. Orlandini^{13,14}, A. Schwenk^{9,10} and J. Simonis^{9,10}



Neutron and weak-charge distributions of ^{48}Ca

ab initio calculations lead to charge distributions consistent with experiment

predict **small neutron skin**,
dipole polarizability, and
weak formfactor



Outline

Advances in nuclear forces and nuclear structure

- opens up ab initio calculations of $0\nu\beta\beta$ nuclei
- theoretical uncertainties from chiral EFT

Progress in nuclear matrix elements for
neutrinoless double beta decay

Perspectives

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Advances in nuclear forces and nuclear structure

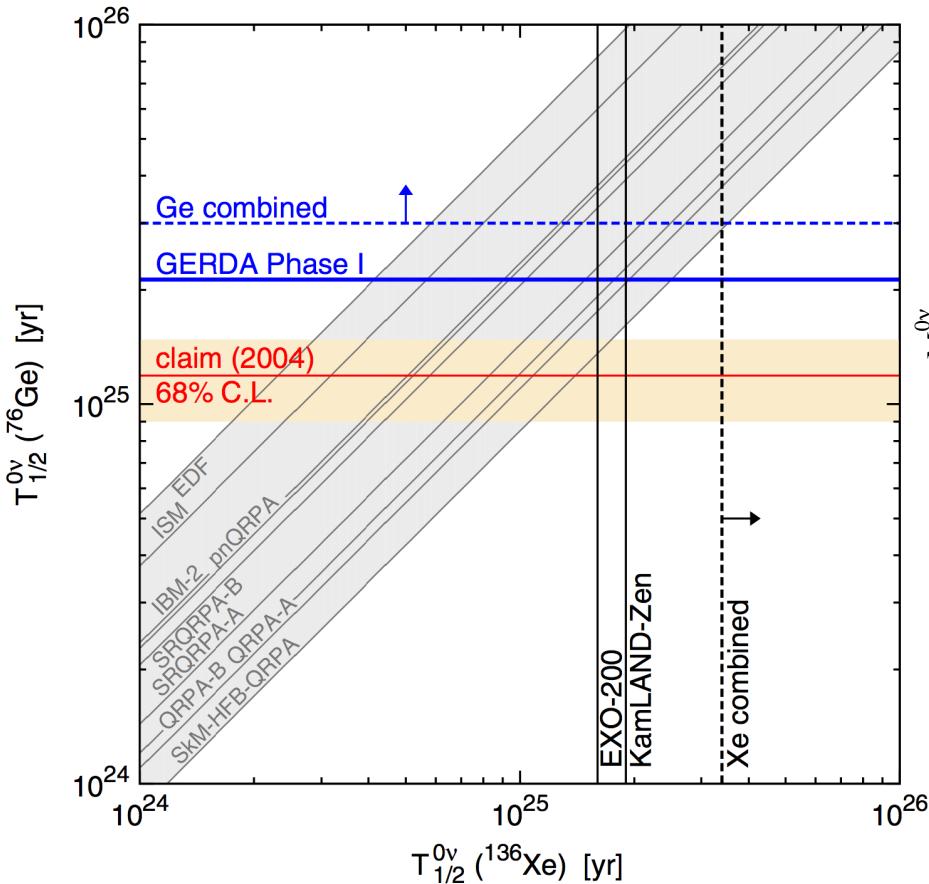
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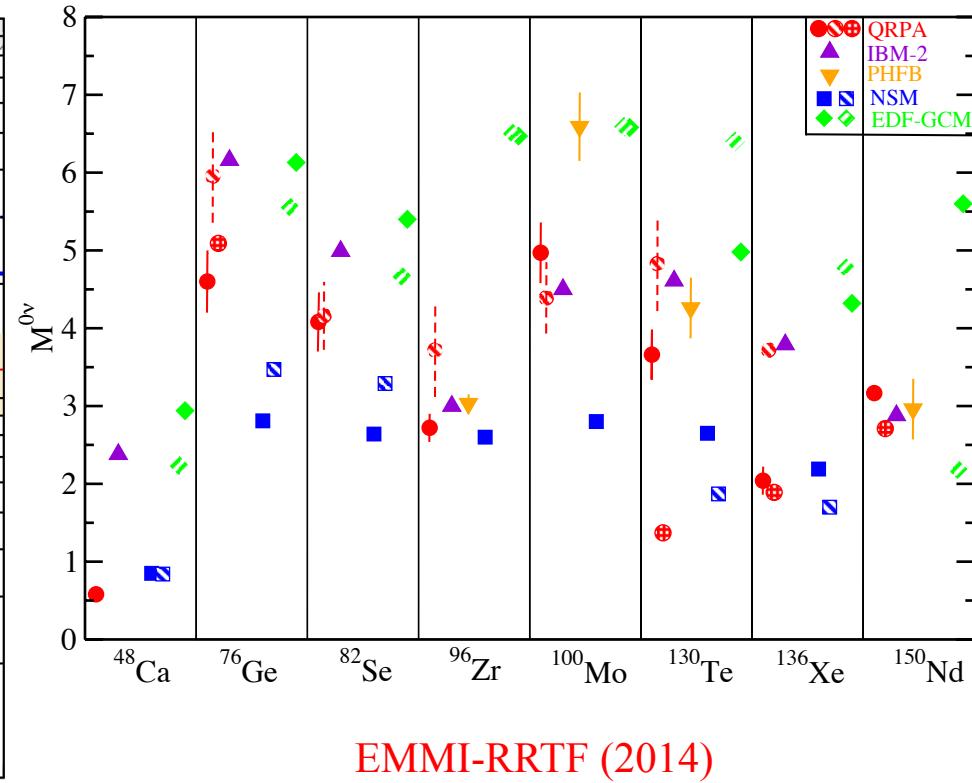
Perspectives

Range of nuclear matrix elements

different NME calculations result in large spread



GERDA Collaboration (2013)

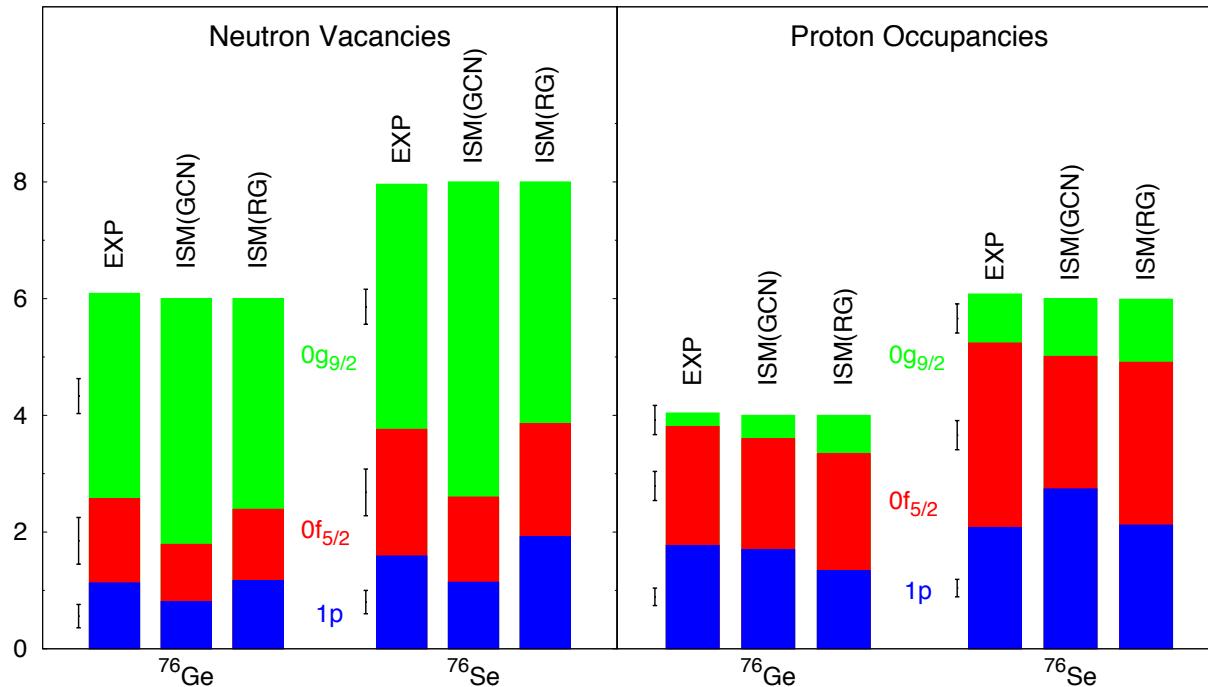


Can we understand the differences in nuclear structure?

Nuclear structure of $0\nu\beta\beta$ nuclei

shell model calculations: valence spaces and interactions successfully tested with experiment (spectra, occupancies)! largest spaces used

Experimental occupancies are well described!



Calculations using state-of-the-art ISM interactions and valence spaces (NATHAN code)

$$M^{0\nu\beta\beta} =$$

2.81 (GCN)
3.26 (RG)

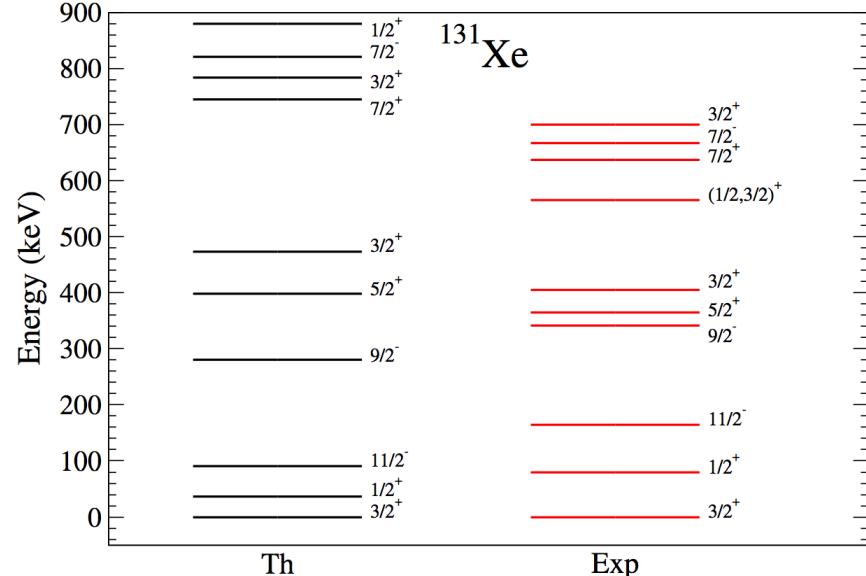
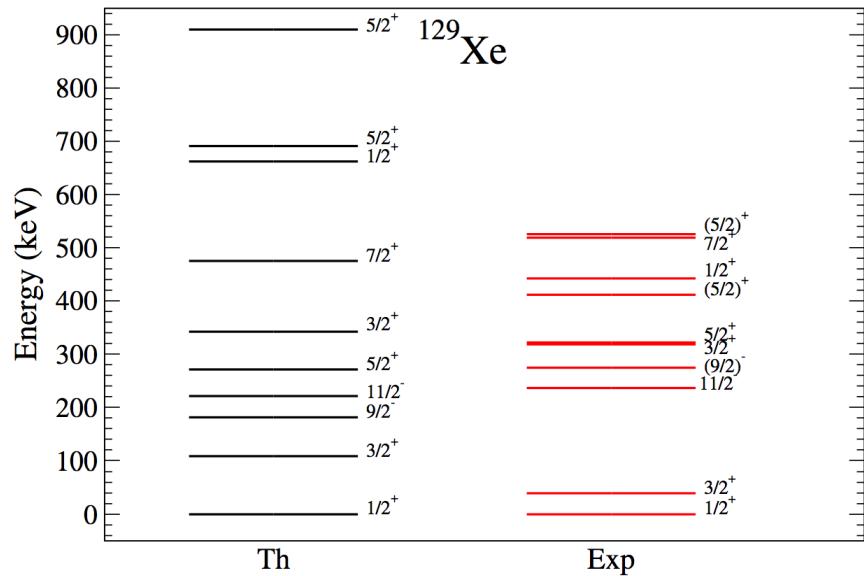
Experiment: Schiffer et al. PRL100 112501(2009), Kay et al. PRC79 021301(2009)

Theory: JM, Caurier, Nowacki, Poves PRC80 048501 (2009)

Other methods do not demonstrate that structure is well reproduced

An example from WIMP-nucleus scattering

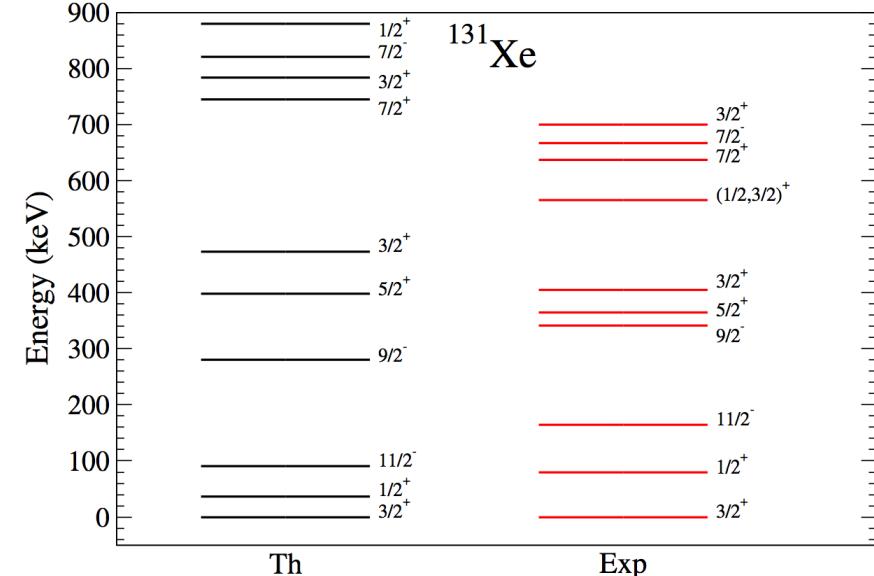
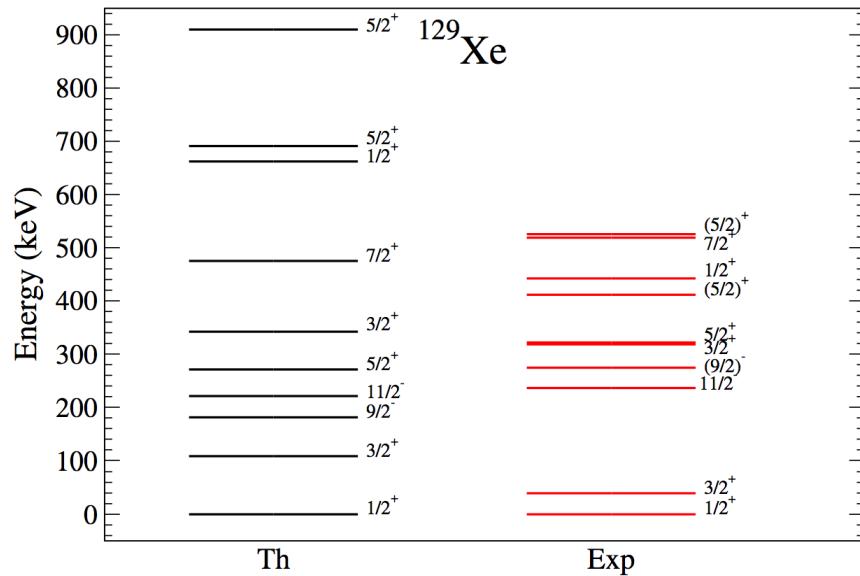
shell model calculation based on the same interactions as for $0\nu\beta\beta$
very good agreement for spectra, ordering and grouping well reproduced



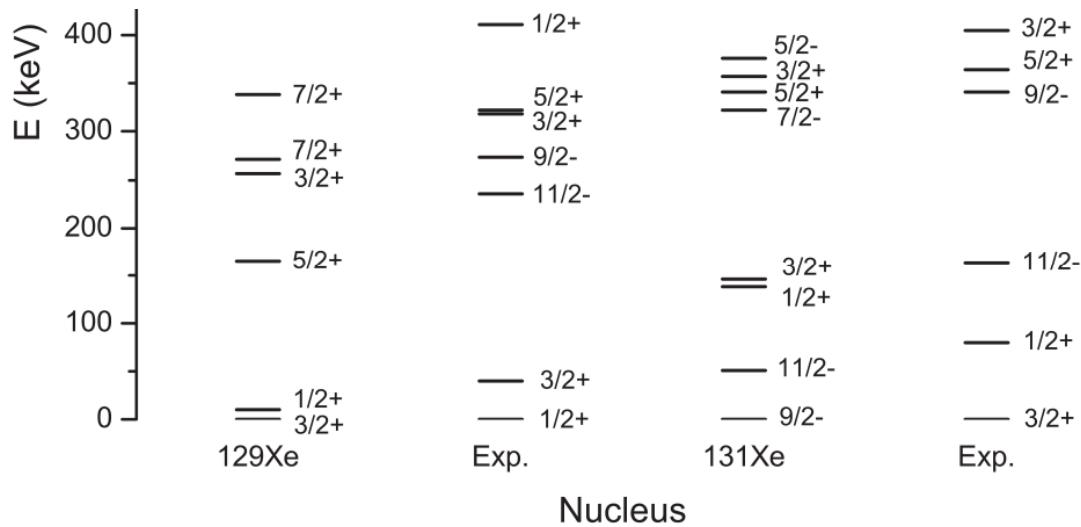
Menendez, Gazit, AS (2012)

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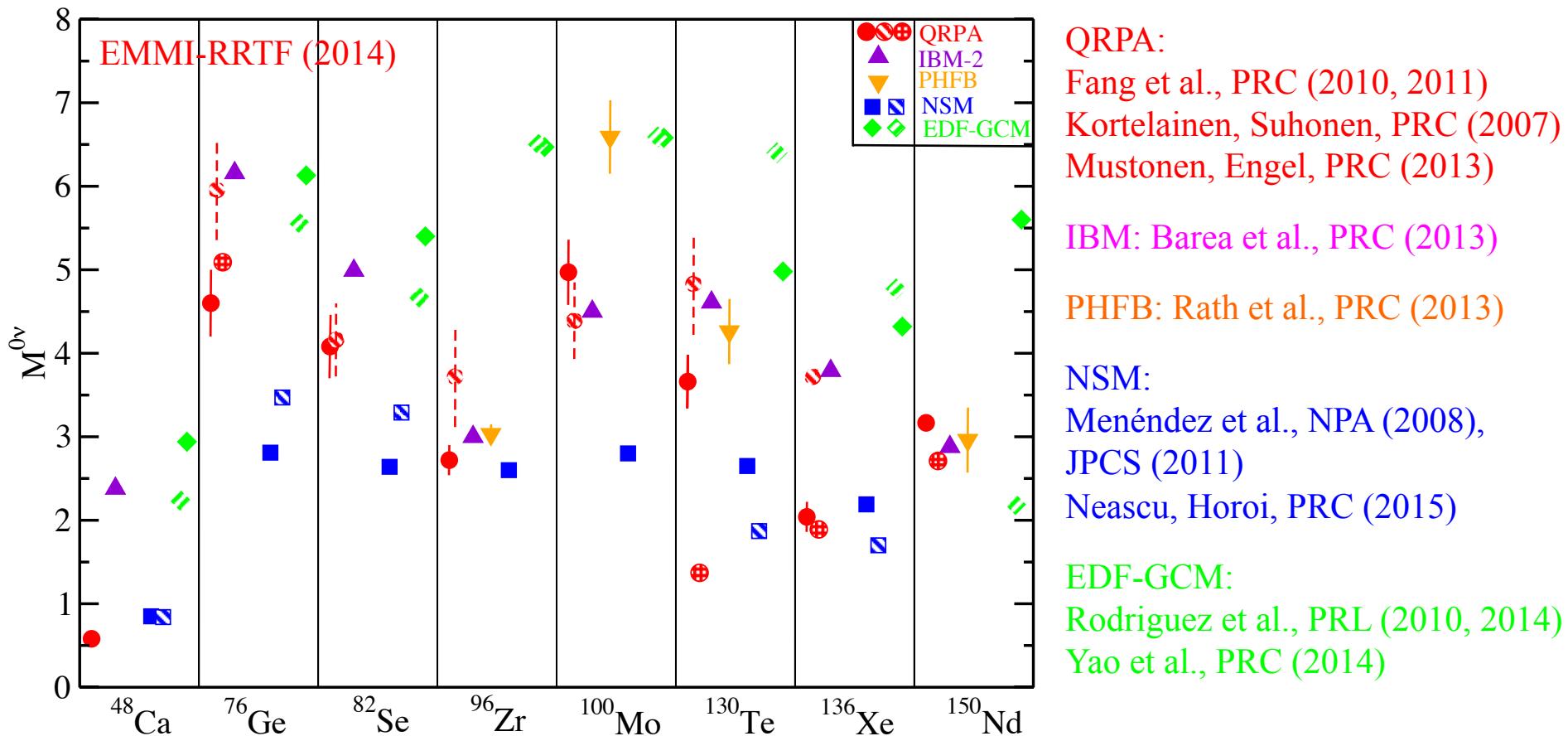
Menéndez, Gazit, AS (2012)



Previous “state-of-the-art”
 calculation... Toivanen et al. (2009)

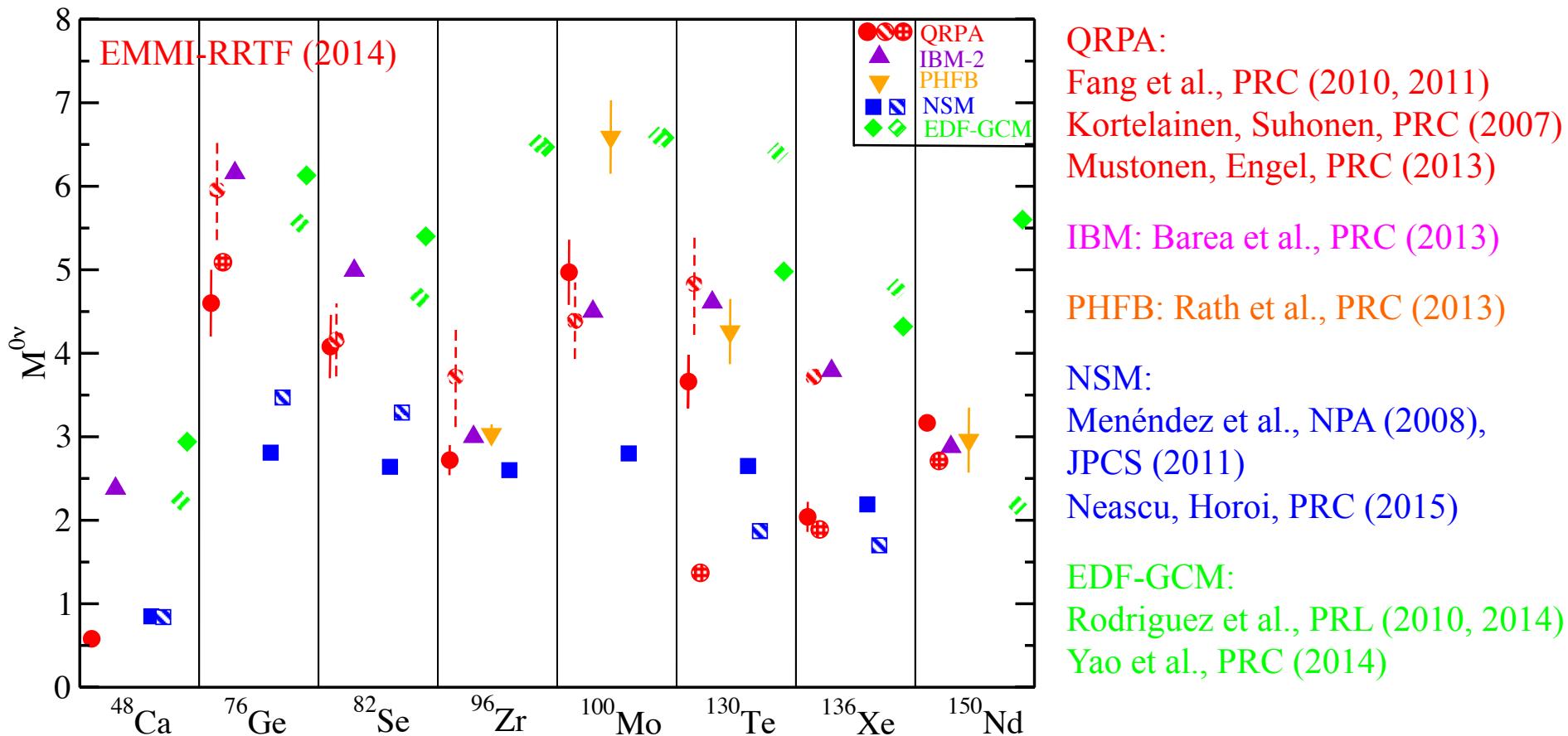
Differences of nuclear matrix elements

NME differ by factor ~ 2 in Ge (range is not good measure of theo. uncertainty)



Differences of nuclear matrix elements

NME differ by factor ~ 2 in Ge (range is not good measure of theo. uncertainty)



largest difference: shell model (NSM) smaller than “mean-field results”
QRPA, PHFB, EDF-GCM

Can we trace this to missing correlations?

Pairing part of nuclear interactions

$T=1, J=0$ interactions included in all calculations,

$T=0, J=1$ channel only in shell model

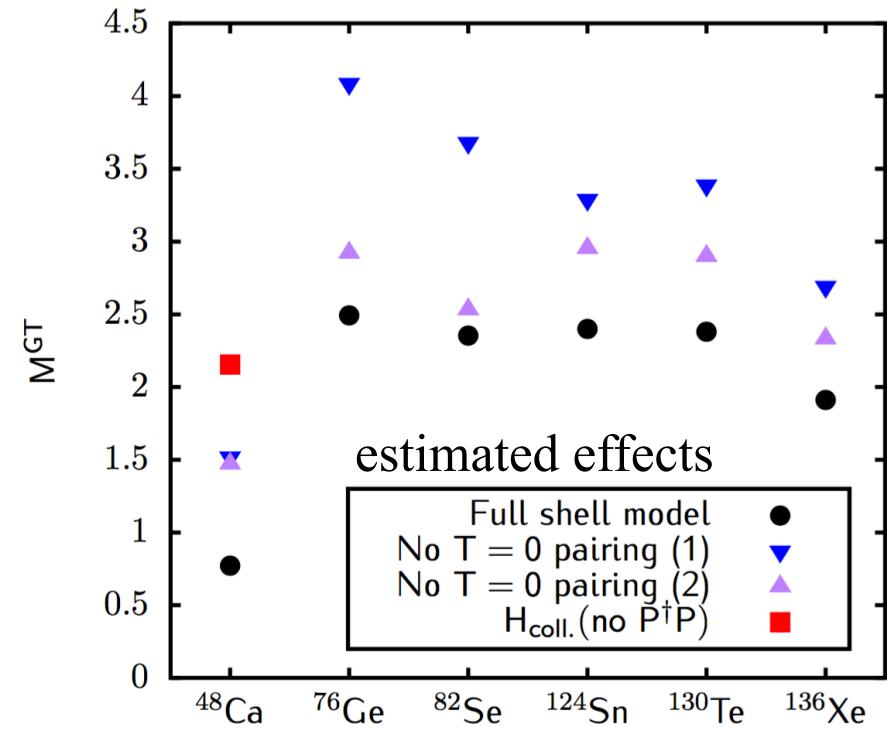
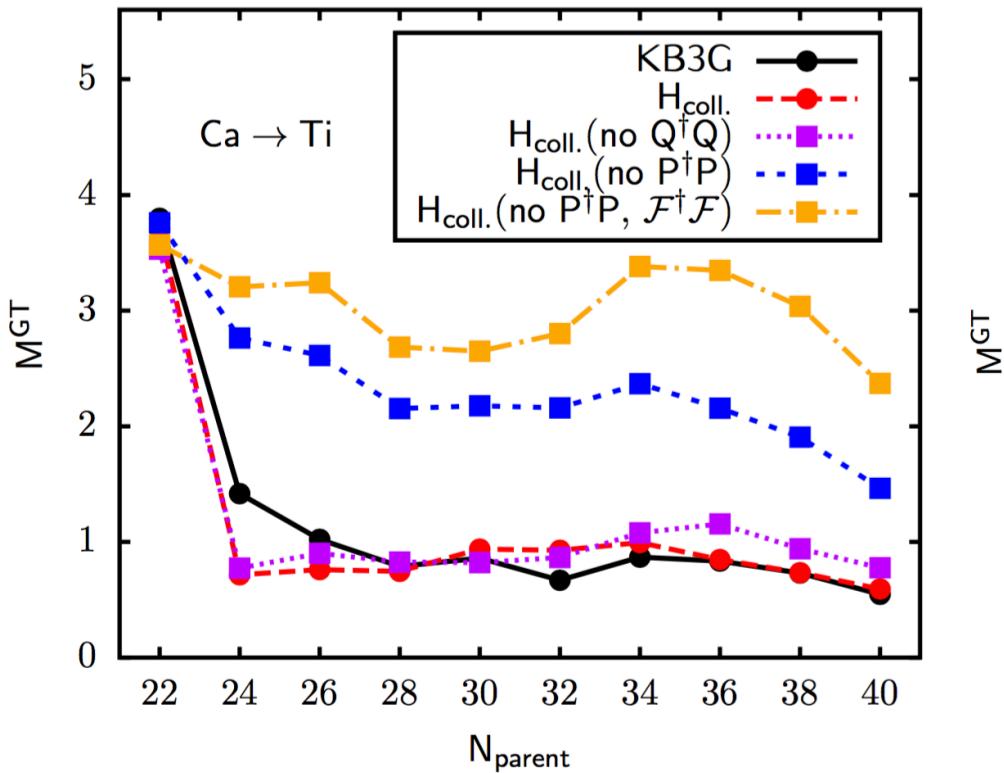
What about correlations due to $T=0, J=1$ interactions (np pairs)?

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What about correlations due to $T=0, J=1$ interactions (np pairs)?

consider Ca chain, without $T=0$ channel NMEs are factor ~ 2 smaller

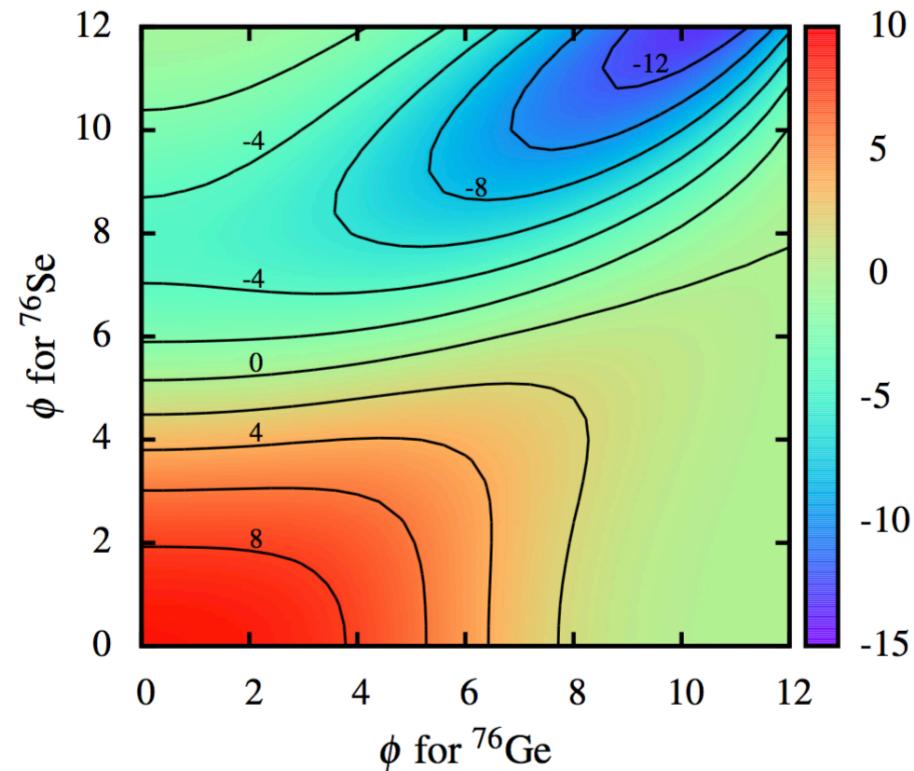
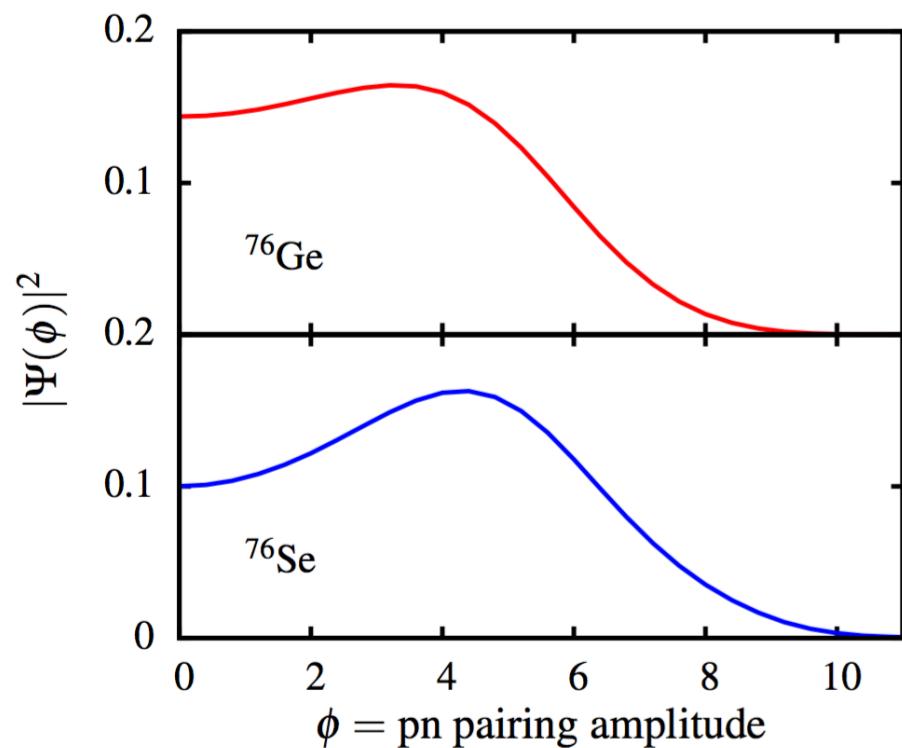


Pairing part of nuclear interactions

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What about correlations due to $T=0, J=1$ interactions (np pairs)?

np pairing correlations in EDF-GCM also reduce NMEs



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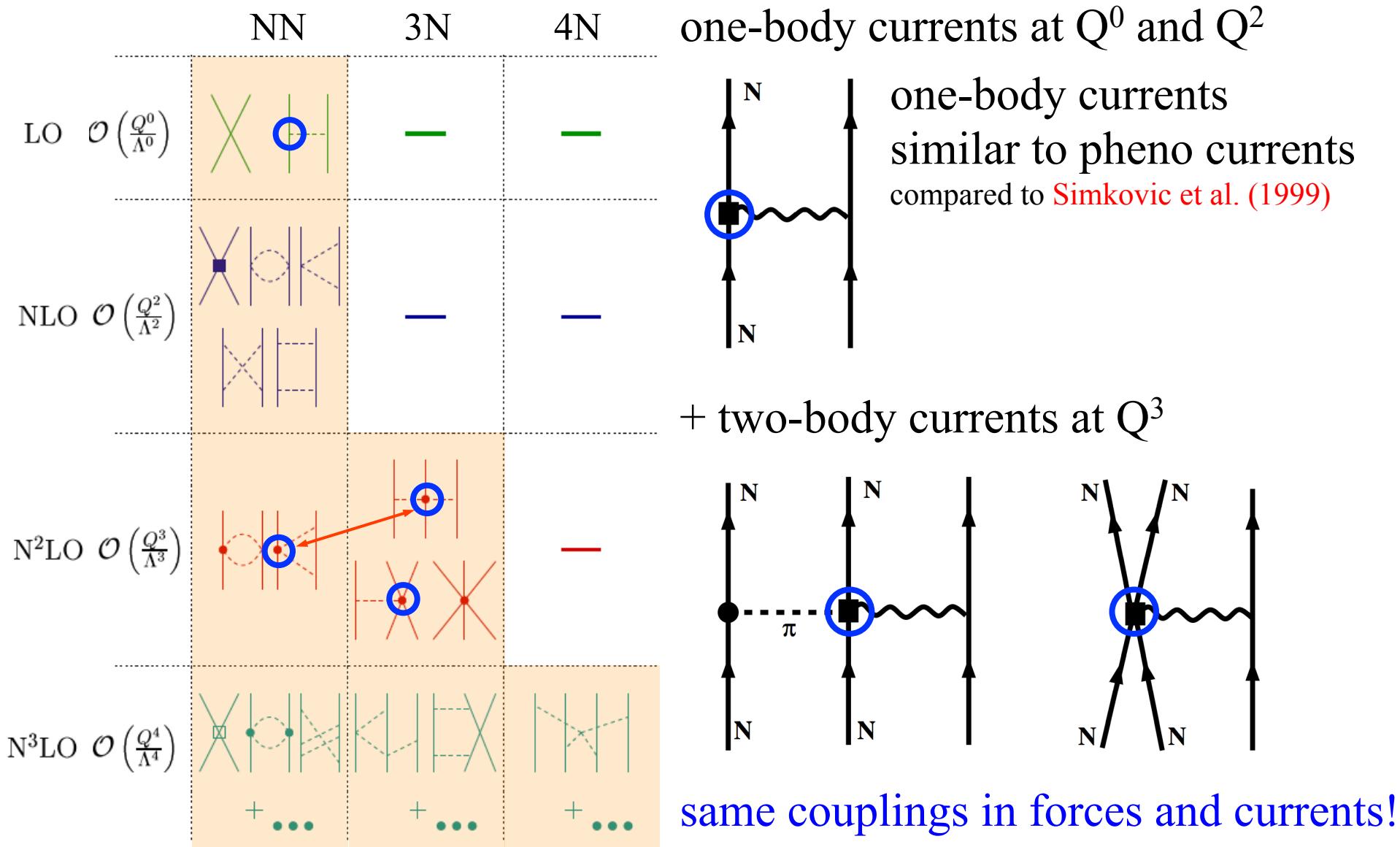
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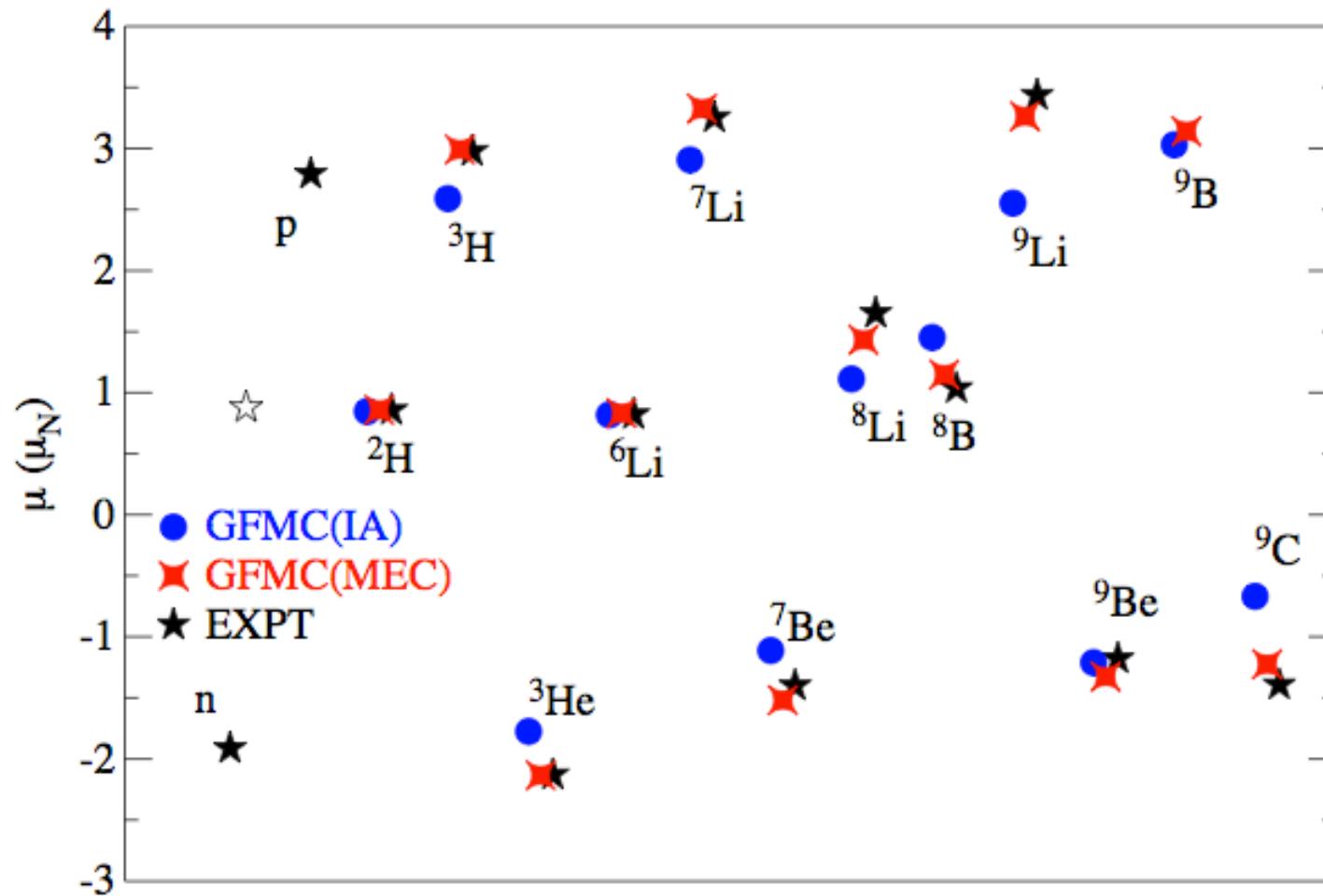
Chiral EFT for weak currents in nuclei



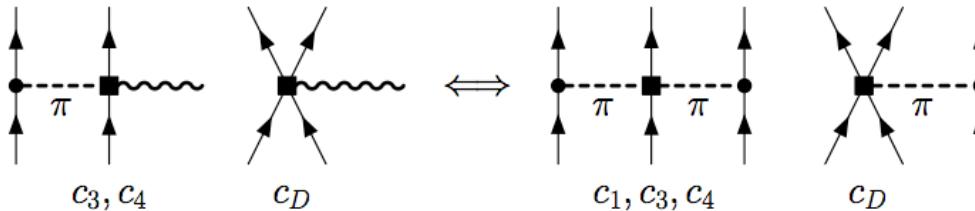
Chiral EFT currents and electromagnetic interactions

predicts consistent electromagnetic 1+2-body currents

GFMC calculations of magnetic moments in light nuclei [Pastore et al. \(2012\)](#)
2-body currents (meson-exchange currents=MEC) are key!



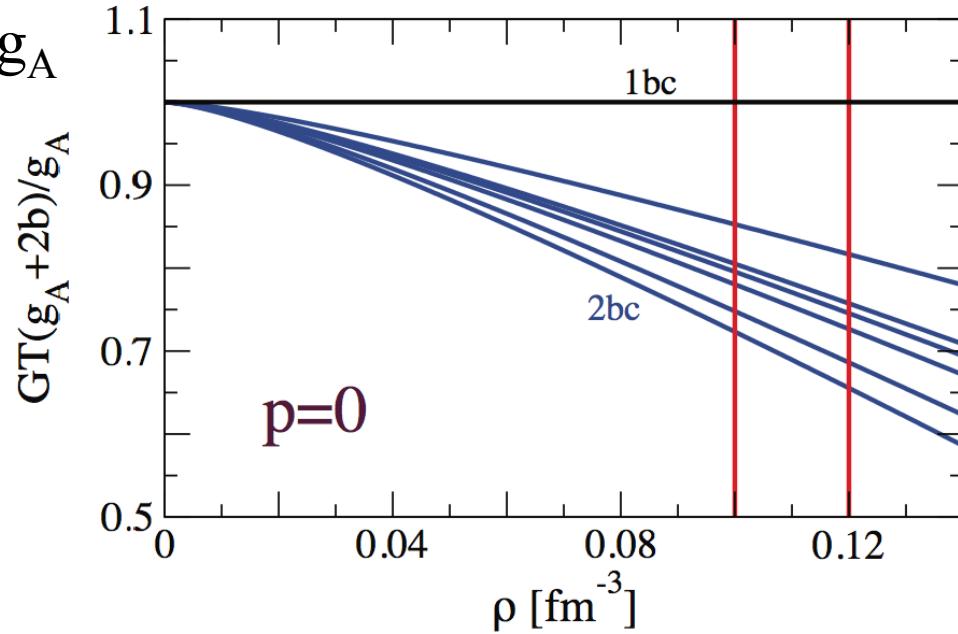
Electroweak interactions and 3N forces



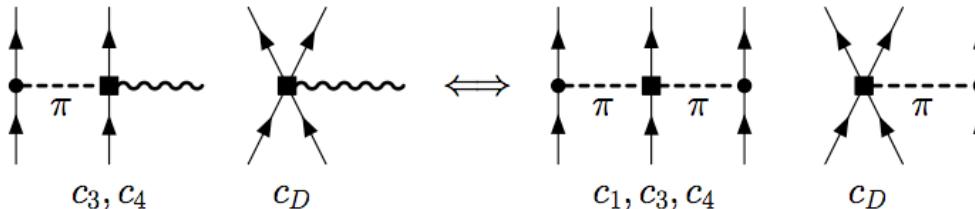
3N couplings predict quenching of g_A
needed in beta decay calculations

Menendez, Gazit, AS (2011)

comparable to empirical $q \sim 0.75$
for β , $2\nu\beta\beta$ decays



Electroweak interactions and 3N forces



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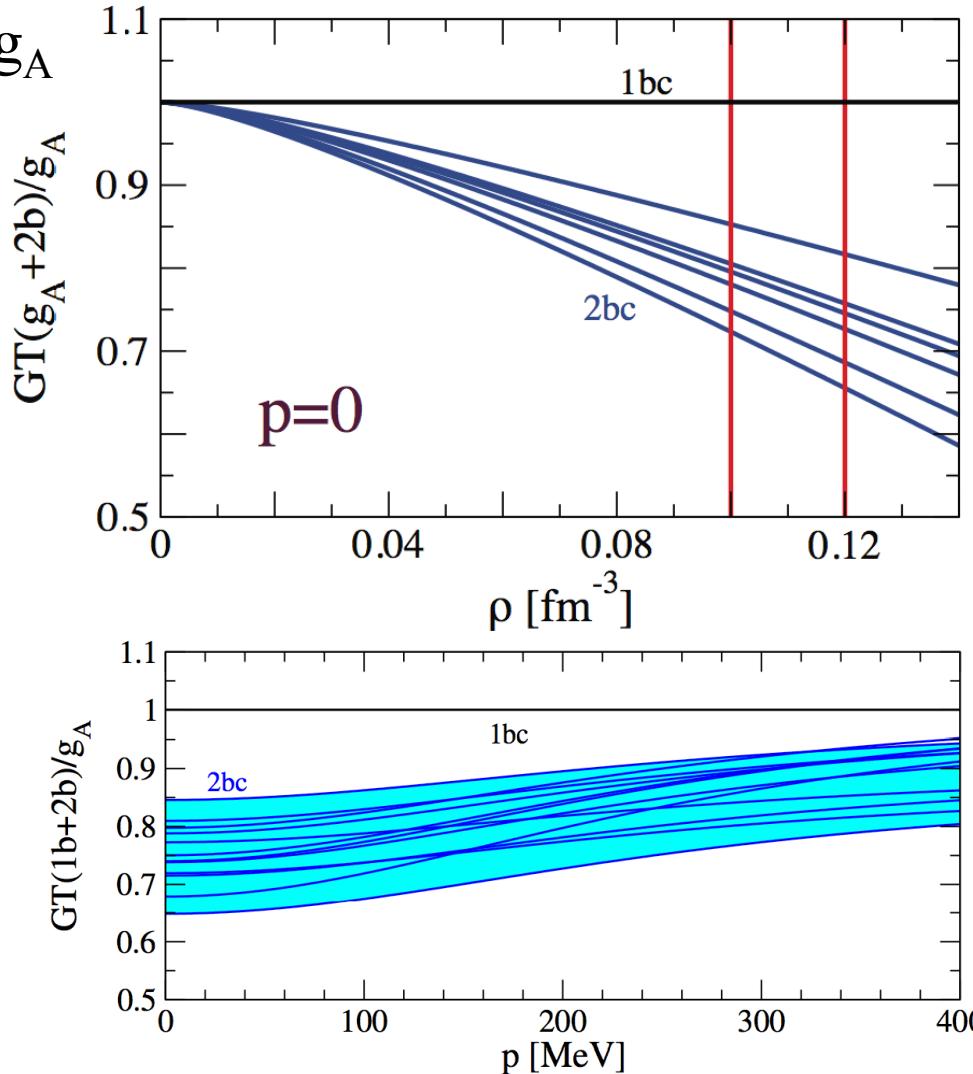
comparable to empirical $q \sim 0.75$
for β , $2\nu\beta\beta$ decays

predicts momentum dependence,
weaker quenching for larger p

Menendez, Gazit, AS (2011)

less quenching for $0\nu\beta\beta$ for $p \sim m_\pi$

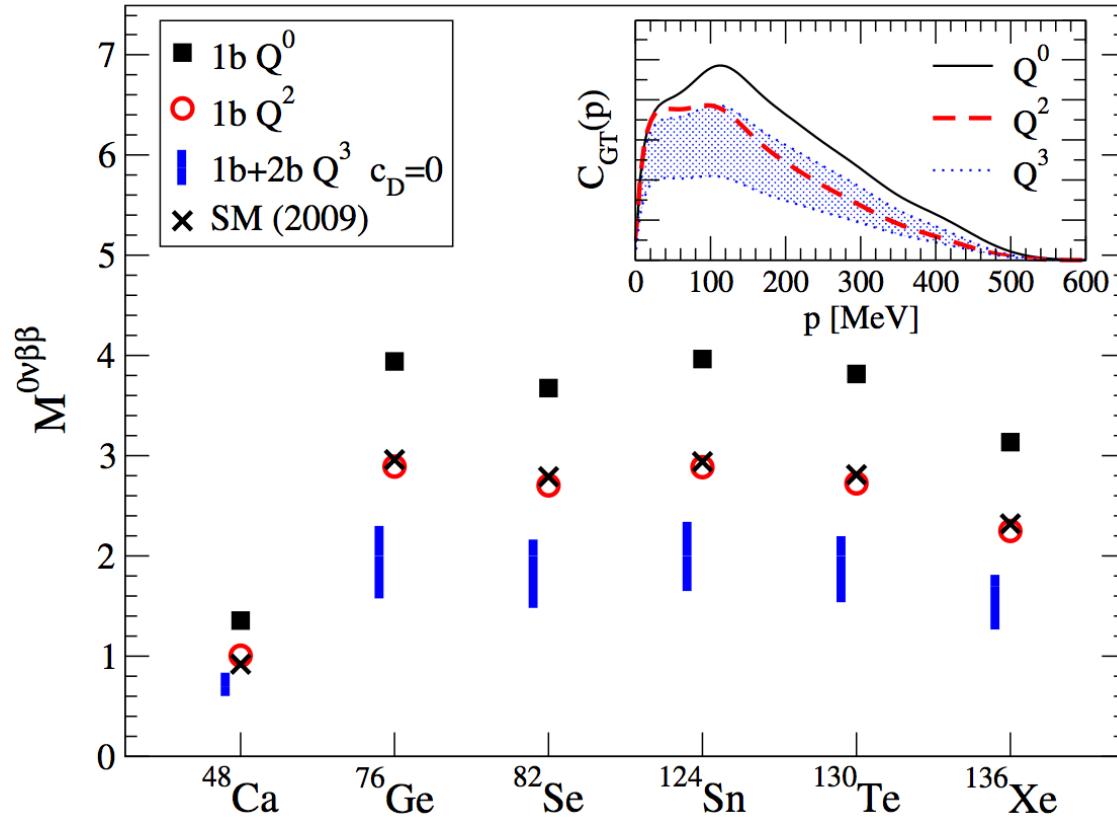
$$M^{0\nu\beta\beta} \propto g_A^2 \Rightarrow \left(T_{1/2}^{0\nu\beta\beta} \right)^{-1} \propto g_A^4$$



Chiral EFT and $0\nu\beta\beta$ decay

NMEs for $0\nu\beta\beta$ decay based on chiral EFT operator

Menendez, Gazit, AS (2011)



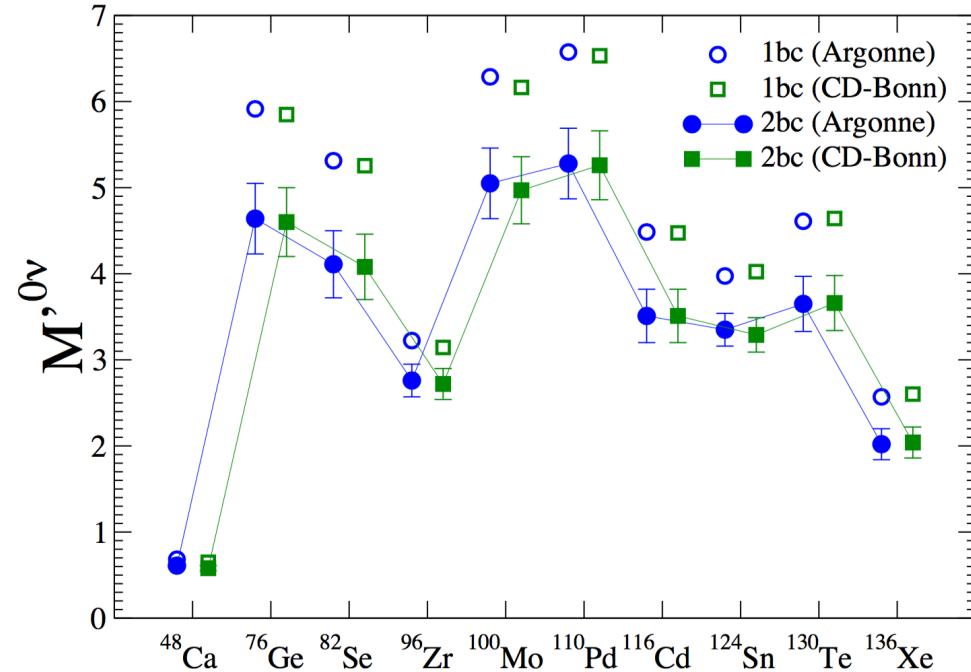
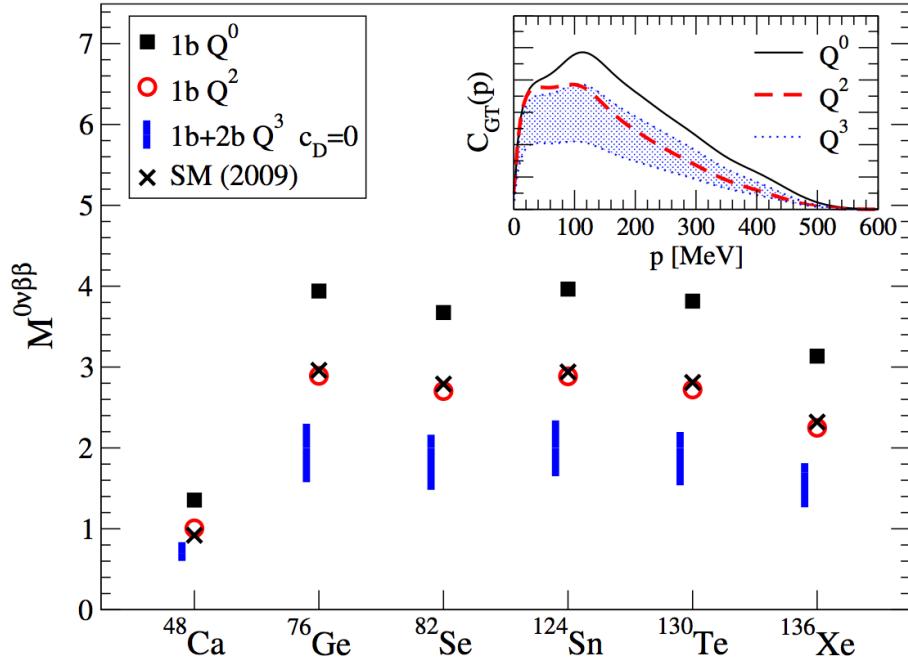
modest quenching because $0\nu\beta\beta$ decay probes higher momentum transfer

two-body currents reduce NME by $\sim 15\% - 40\%$,
need to be included in all calculations of $0\nu\beta\beta$ decay

Chiral EFT and $0\nu\beta\beta$ decay

NMEs for $0\nu\beta\beta$ decay based on chiral EFT operator

Menendez, Gazit, AS (2011)



Engel, Simkovic ,Vogel (2014)

modest quenching because $0\nu\beta\beta$ decay probes higher momentum transfer
 two-body currents reduce NME by $\sim 15\% - 40\%$,
 need to be included in all calculations of $0\nu\beta\beta$ decay

Summary and perspectives

chiral effective field theory

nuclear forces and electroweak/WIMP/... interactions,
systematic for energies below ~ 300 MeV

advances in ab initio calculations up to $A \sim 50$

progress in understanding **nuclear structure differences in NMEs**

first calculation of **$0\nu\beta\beta$ operator in chiral EFT**
less quenching compared to β , $2\nu\beta\beta$ decays

Future perspectives: ^{48}Ca NME ab initio calculation as benchmark

ab initio calculation of NMEs with consistent electroweak ints.

US-DOE topical collaboration lead by J. Engel (UNC), see CIPANP 1511.00074